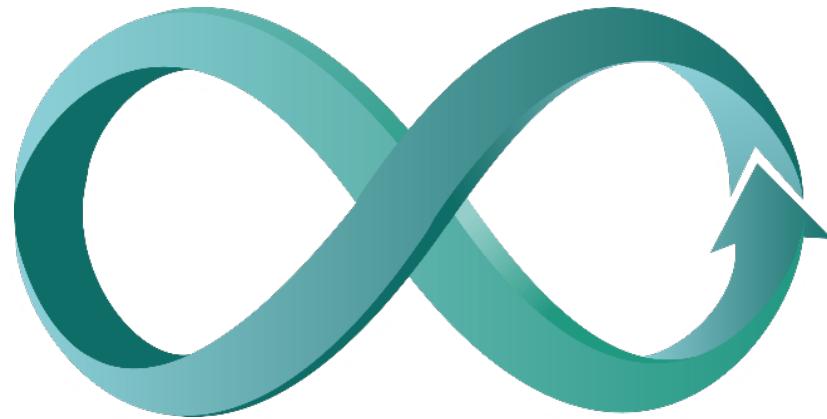




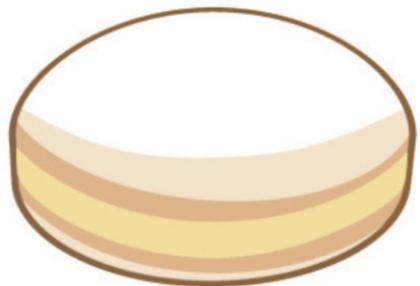
MAX-PLANCK-INSTITUT

Topology, Magnetism und Chirality



Claudia Felser, MPI CPfS
felser@cpfs.mpg.de

topology

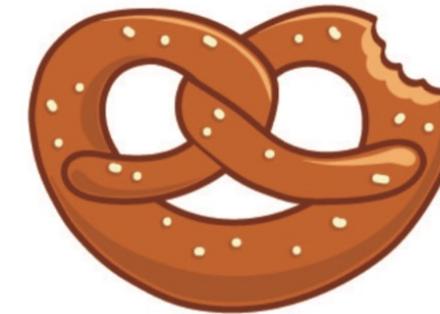


Genius $g = 0$

Gauss–Bonnet
Theorem: $\chi = 2(1 - g) = \frac{1}{2\pi} \int K \cdot dS$



$g = 1$



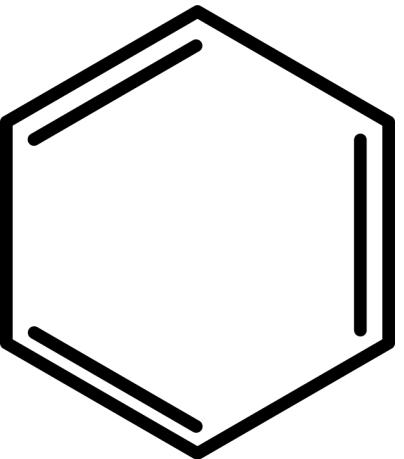
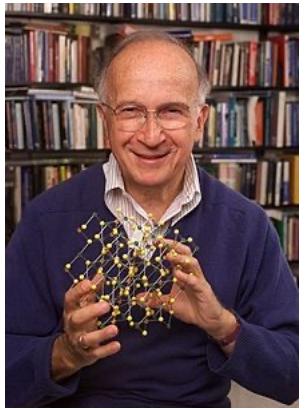
$g = 3$

C. F. Gauss, , Werke , 8 , K. Gesellschaft Wissenschaft. Göttingen (1900)
O. Bonnet, J. École Polytechnique , 19 (1848) pp. 1–146
J. Kroder, et al., Chemie in unserer Zeit 56 (2022) 12–20

topology

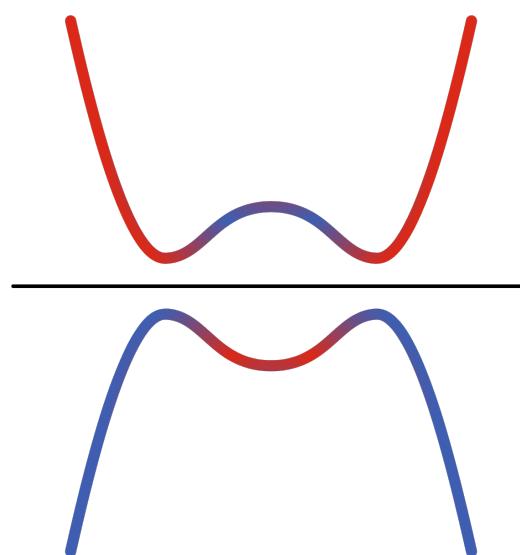
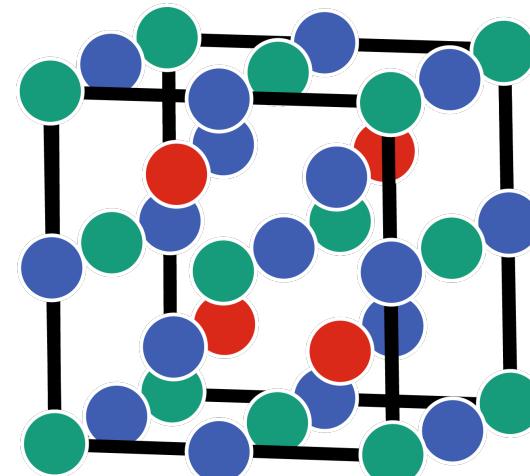
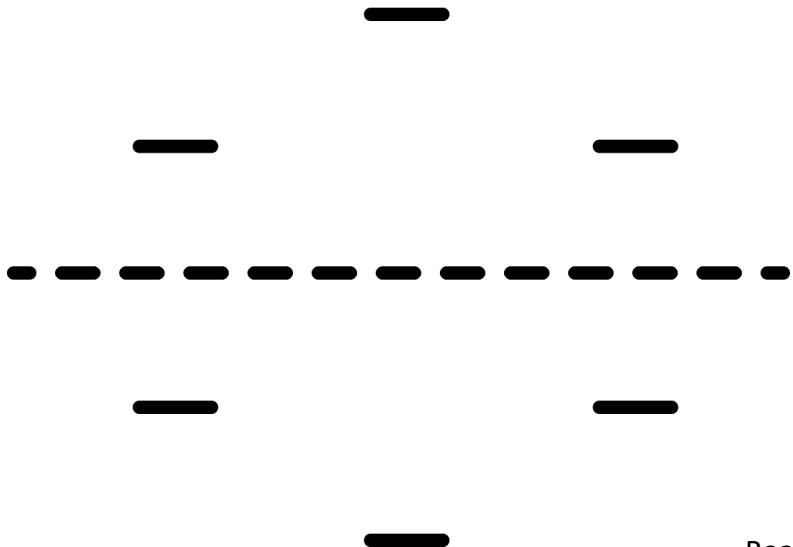


from molecule to solid

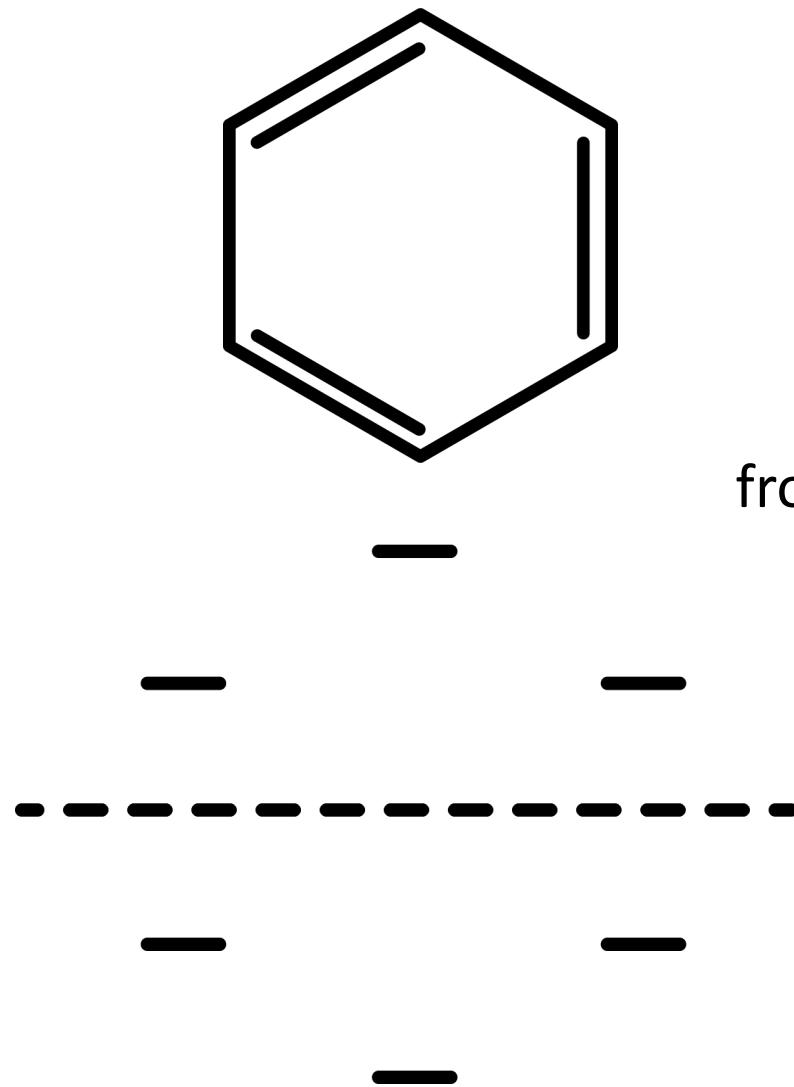


$$n = 0 \quad 1 \quad 2 \quad 3 \quad 4 \dots$$
$$\chi_0 \quad \chi_1 \quad \chi_2 \quad \chi_3 \quad \chi_4$$
$$\psi_k = \sum_n e^{ikna} \chi_n$$

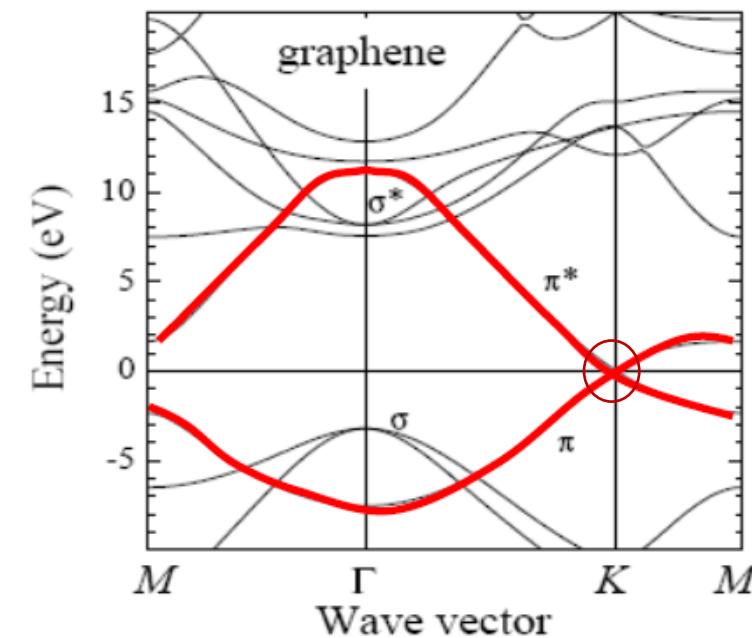
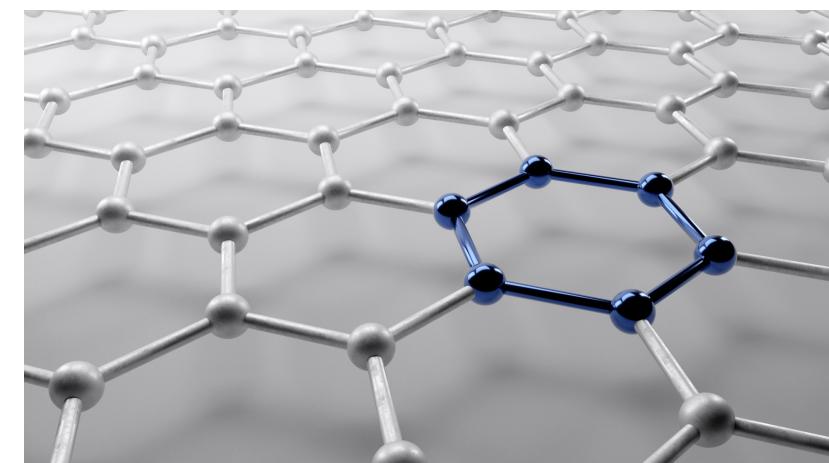
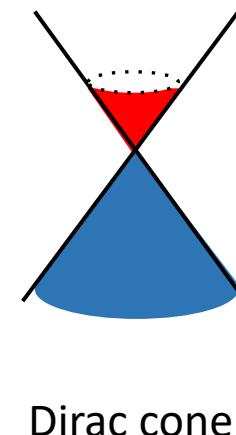
from orbitals to bands



from benzen to graphene

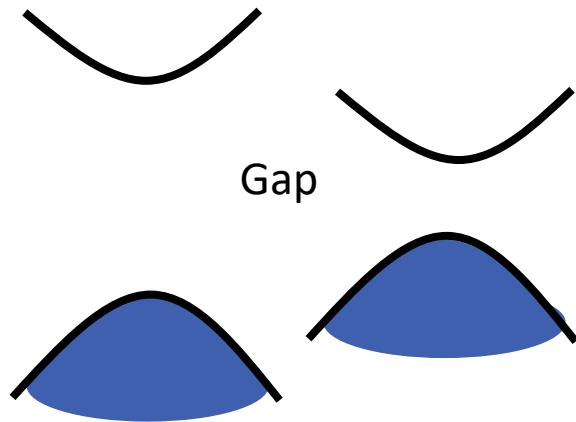


from bonds to bands

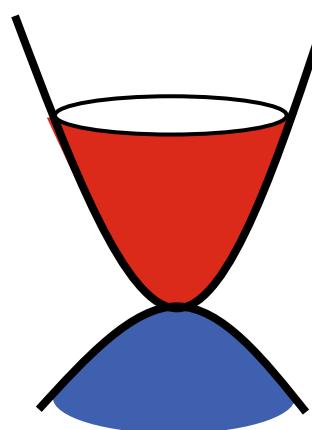


electronic structure

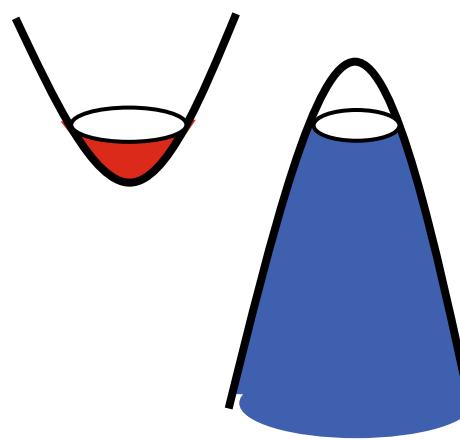
Insulator Semiconductor



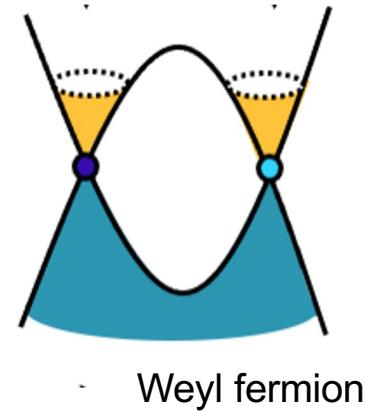
Metal



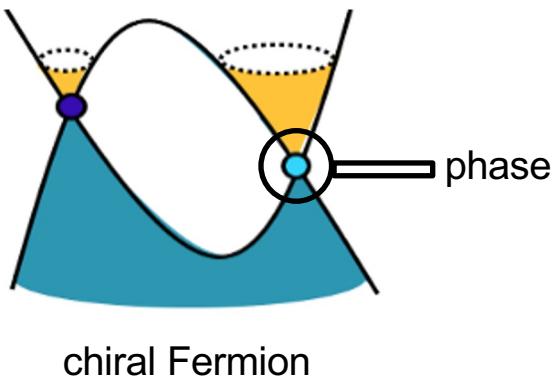
Semi metal



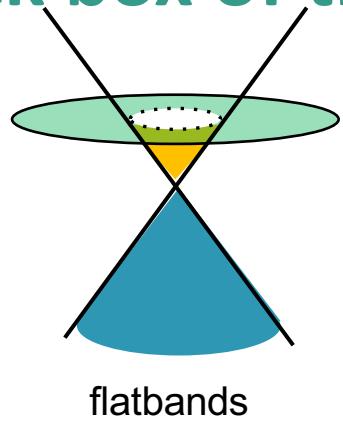
black box of the electronic structure



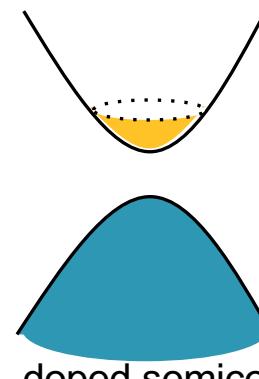
Weyl fermion



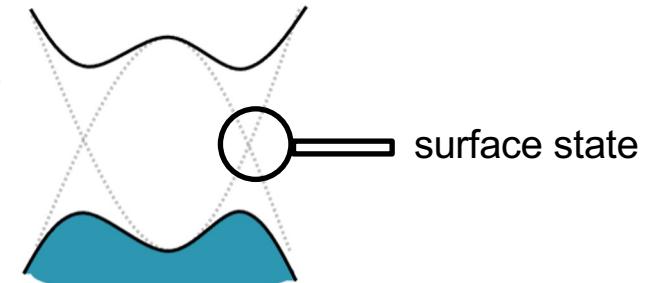
chiral Fermion



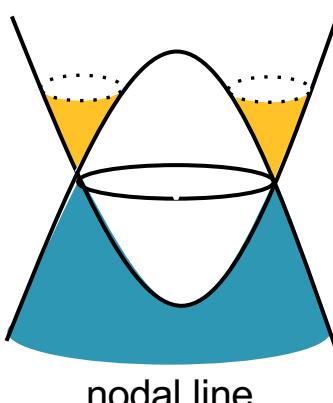
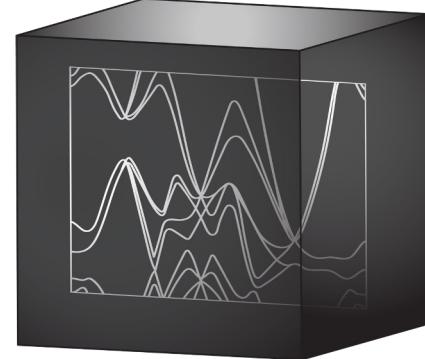
flatbands



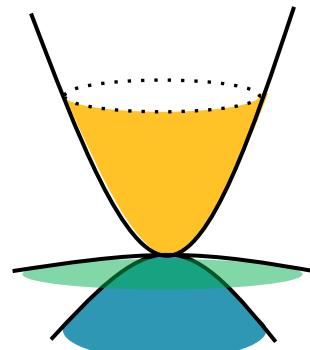
doped semiconductors



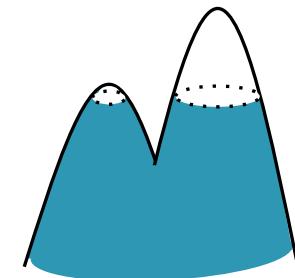
surface state



nodal line

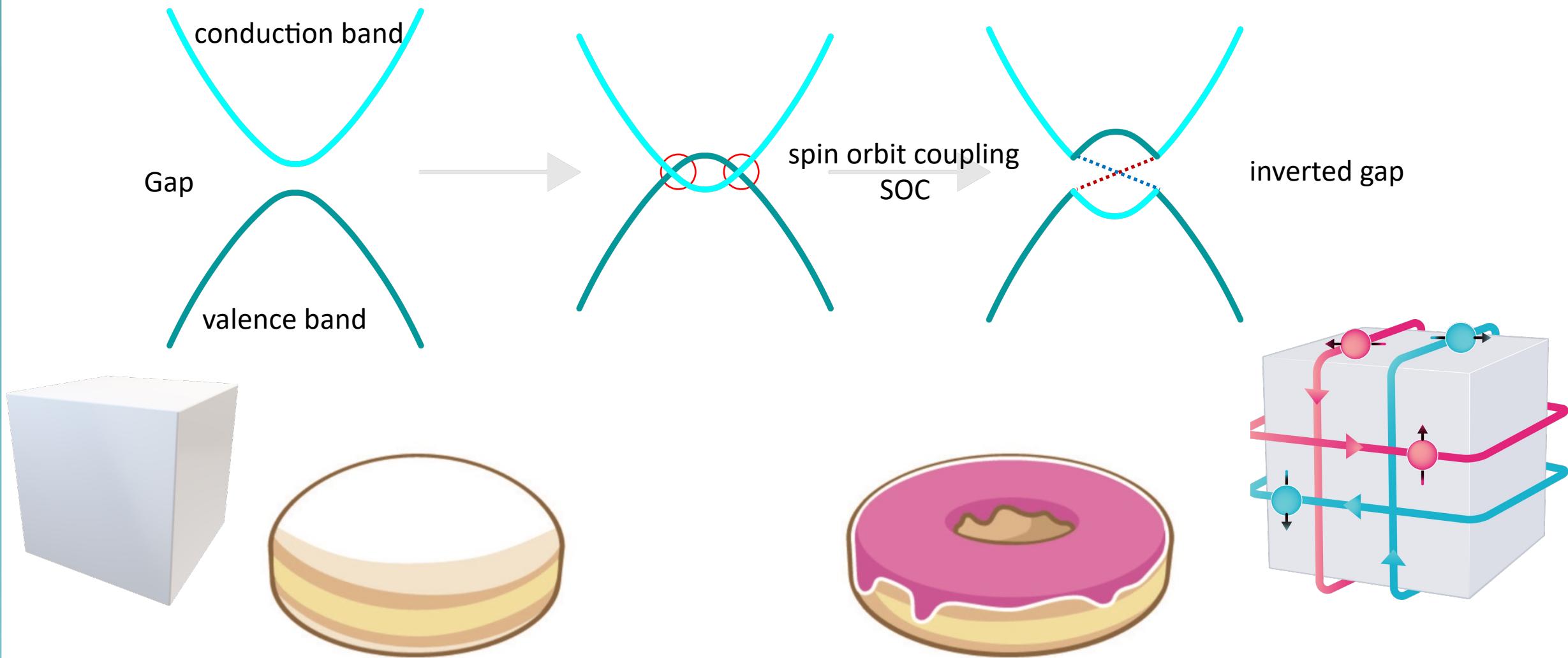


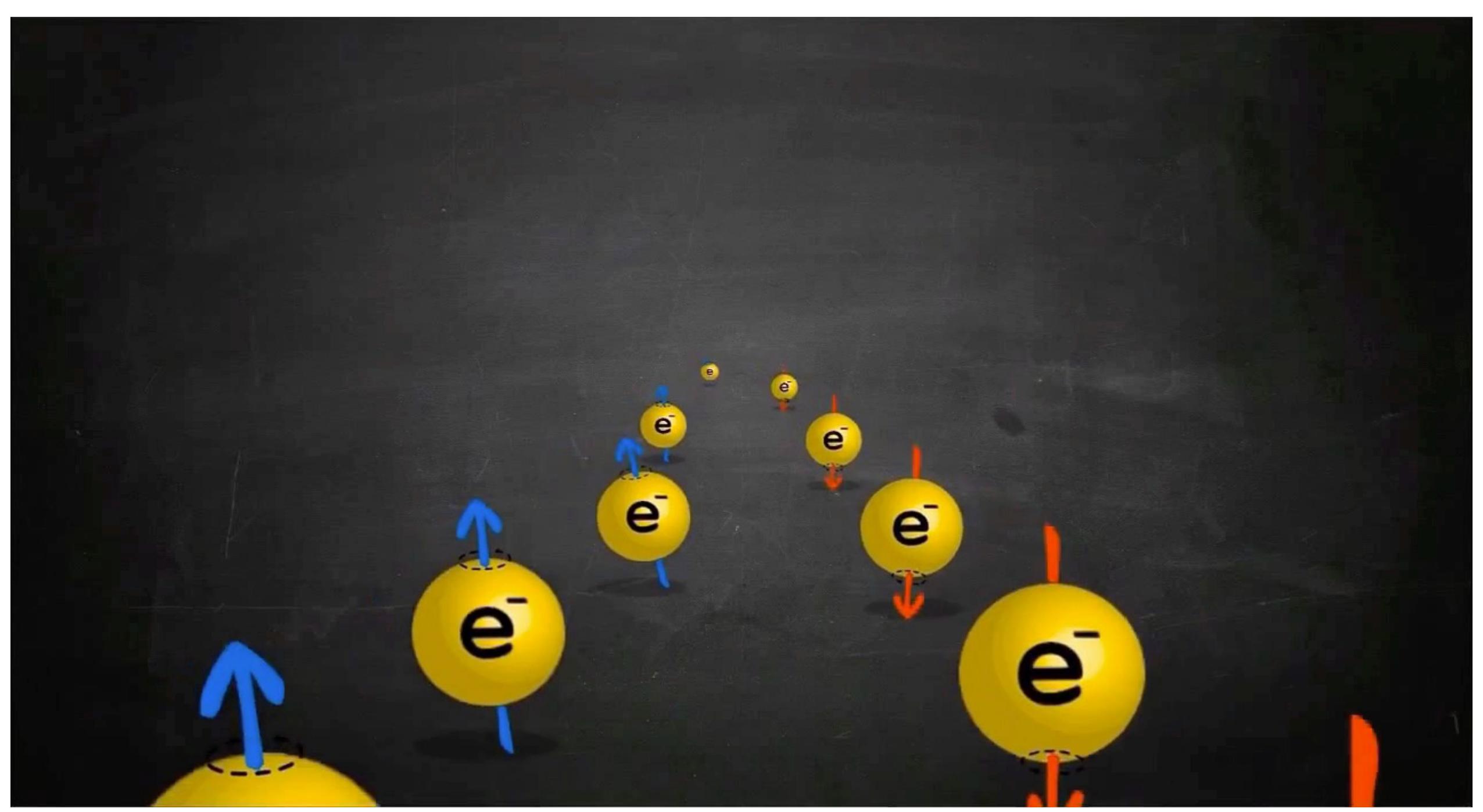
new Fermion



van Hove Singularity

topological insulators







from all 250 000 known
inorganic compounds
25% are topological

TICLE

RESEARCH

Topological

Barry Bradlyn^{1*}, L. Elcoro^{2*}, Jennifer



TOPOLOGICAL MATTER

All topological bands of all nonmagnetic stoichiometric materials

Maia G. Vergniory^{*†}, Benjamin J. Wieder^{*†}, Luis Elcoro, Stuart S. P. Parkin, Claudia Felser, B. Andrei Bernevig, Nicolas Regnault^{*}

topological materials

<https://doi.org/10.1038/s41586-020-2837-0>

Received: 27 January 2020 | Accepted: 24 August 2020 | <https://doi.org/10.1038/s41586-020-2837-0>

Yuanfeng Xu¹, Luis Elcoro², Zhida Song³, Benjamin J. Wieder^{3,4,5}, M. G. Vergniory^{6,7}, Nicolas Regnault^{3,8}, Yulin Chen^{9,10,11,12}, Claudia Felser^{13,14} & B. Andrei Bernevig^{1,3,15}



^{8,9*} & Zhijun Wang^{7,10}

<https://doi.org/10.1038/s41586-019-0937-5>

topological indicators

<https://doi.org/10.1038/s41586-019-0944-6>

electronic

ing^{1,4}, Hongming Weng^{1,5,6,7,8*} &

s of magnetic



topologicalquantumchemistry.org

37 Rb	38 Sr	39 Y	40 Zr	41 Nb	42 Mo	43 Tc	44 Ru	45 Rh	46 Pd	47 Ag	48 Cd	49 In	50 Sn	51 Sb	52 Te	53 I	54 Xe
55 Cs	56 Ba	57 La	72 Hf	73 Ta	74 W	75 Re	76 Os	77 Ir	78 Pt	79 Au	80 Hg	81 Tl	82 Pb	83 Bi	84 Po	85 At	86 Rn
87 Fr	88 Ra	89 Ac	104 Rf	105 Db	106 Sg	107 Bh	108 Hs	109 Mt	110 Ds	111 Rg	112 Cn	113 Nh	114 Fl	115 Mc	116 Lv	117 Ts	118 Og
			58 Ce	59 Pr	60 Nd	61 Pm	62 Sm	63 Eu	64 Gd	65 Tb	66 Dy	67 Ho	68 Er	69 Tm	70 Yb	71 Lu	
			90 Th	91 Pa	92 U	93 Np	94 Pu	95 Am	96 Cm	97 Bk	98 Cf	99 Es	100 Fm	101 Md	102 No	103 Lr	

12 Entries found for **Hf, Te**, showing:

ALL (12)

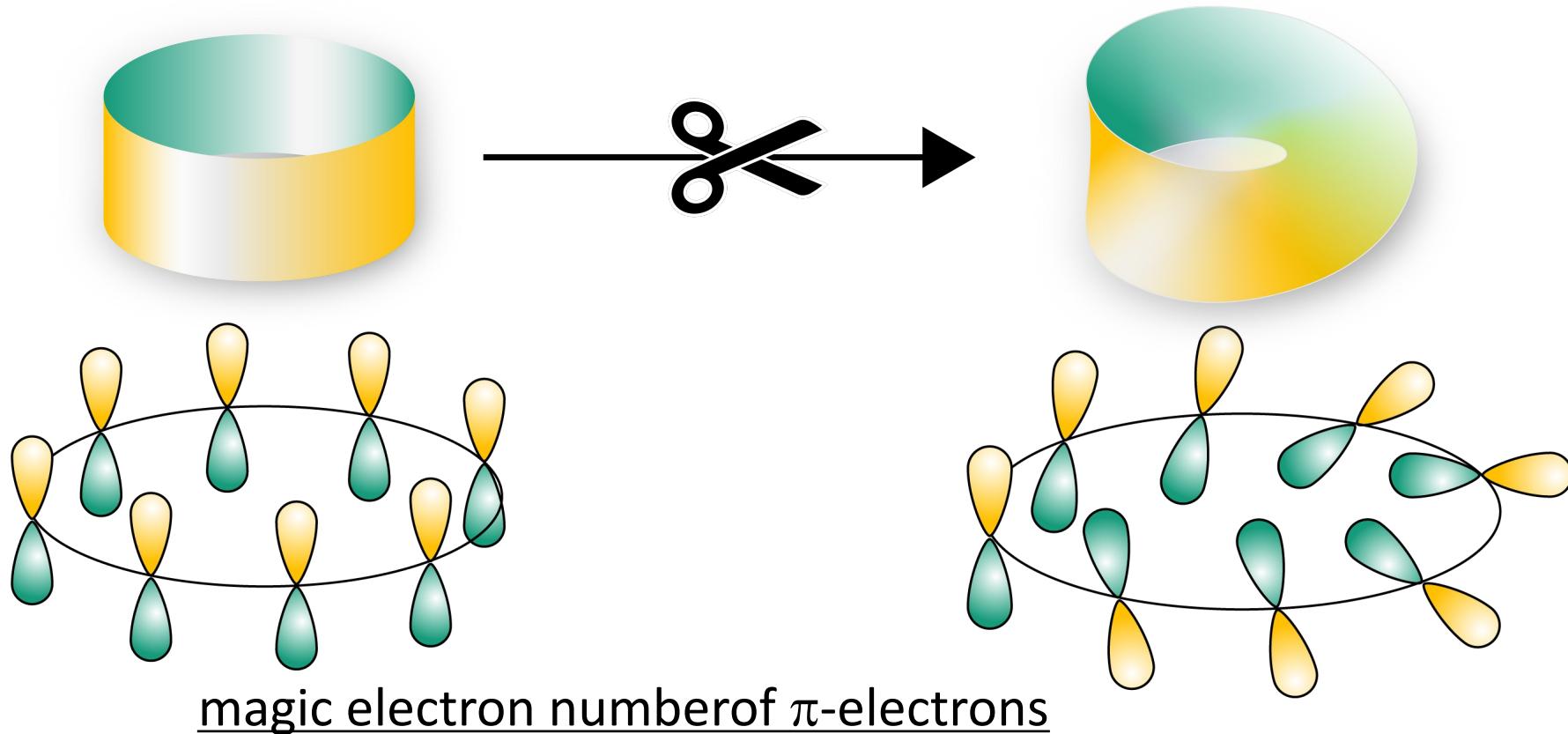
TI (5)

SM (1)

Trivial (6)

Compound	Symmetry Group	Topological Indices	Crossing Type	Type
Hf ₁ Te ₂	164 (<i>P-3m1</i>)		Line	ES
Hf ₁ Te ₃	11 (<i>P2₁/m</i>)			LCEBR
Hf ₁ Te ₅	63 (<i>Cmcm</i>)	$Z_4=3, Z_{2w,1}=1, Z_{2w,2}=1$		NLC

chemistry



Hückel:

$4n+2$ aromatic

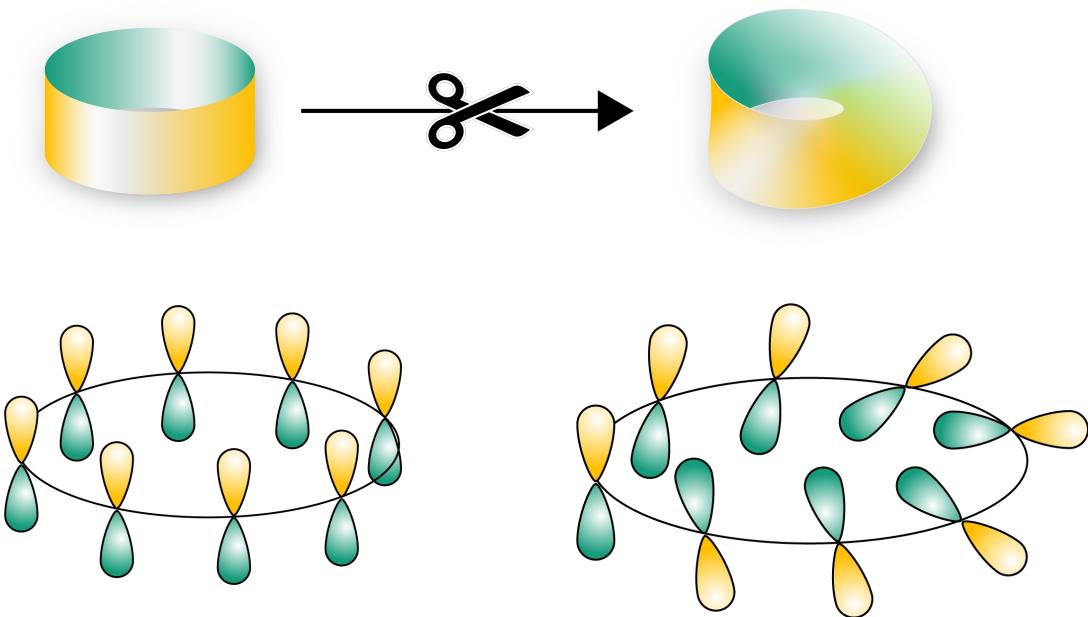
$4n$ antiaromatic

Möbius:

$4n$ aromatic

$4n+2$ antiaromatic

topology in chemistry



Magic electron numbers

Hückel:

$4n+2$	aromatic
$4n$	antiaromatic

Möbius:

$4n$	aromatic
$4n+2$	antiaromatic

ORGANIC CHEMISTRY

Aromatics with a twist

Rainer Herges

The properties of flat aromatic molecules are well known to chemists, but some non-planar aromatics remain a mystery. A molecule that can twist into a Möbius band on command might shed light on their features.

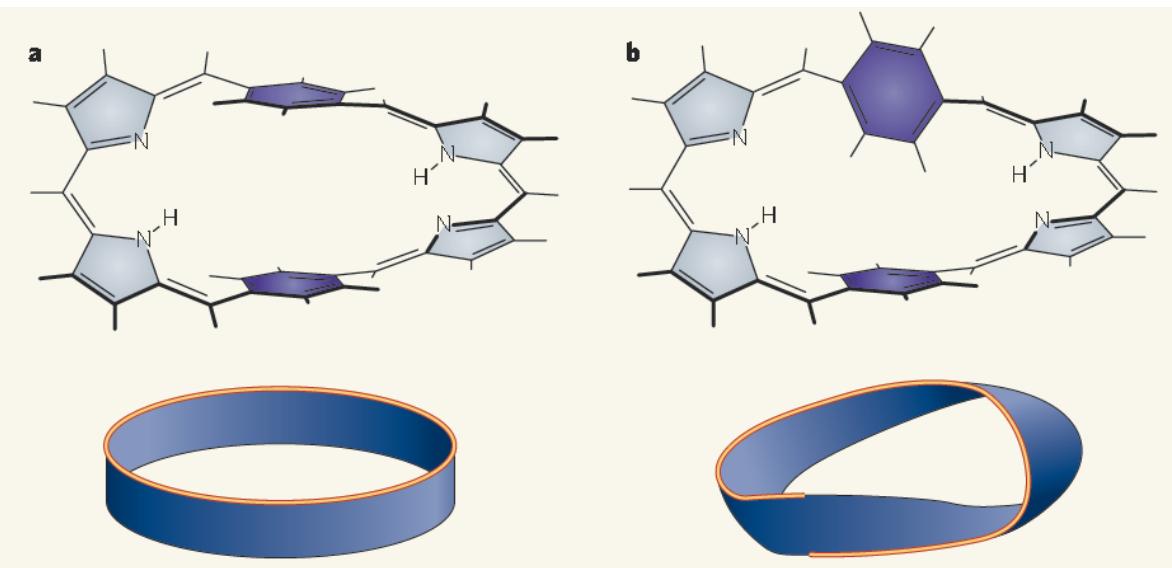


Figure 2 | A molecular topological switch. Latos-Grażyński and colleagues¹ have made a compound that is antiaromatic in nonpolar solvents, but not in polar solvents. **a**, In nonpolar solvents, the two benzene rings (purple) in the molecule are parallel, and the molecule is a two-sided, non-twisted band. **b**, In polar solvents, the upper benzene ring twists by 90°, so that the molecule becomes a one-sided, Möbius structure. This conformational change alters the aromaticity of the molecule.



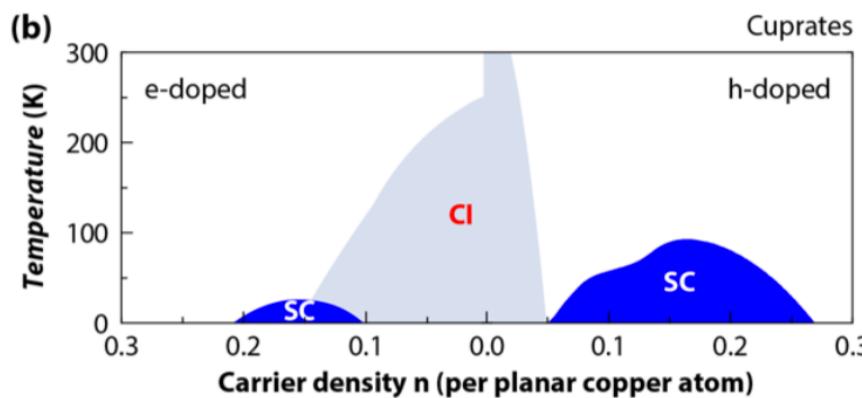
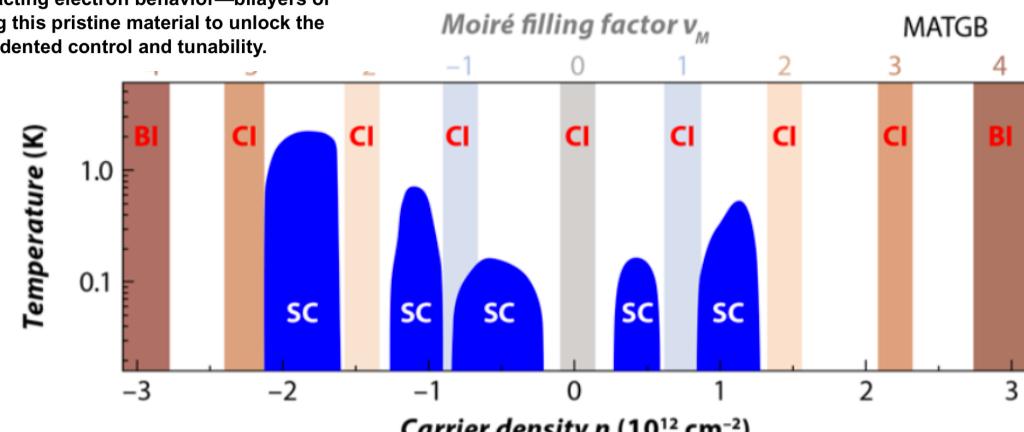
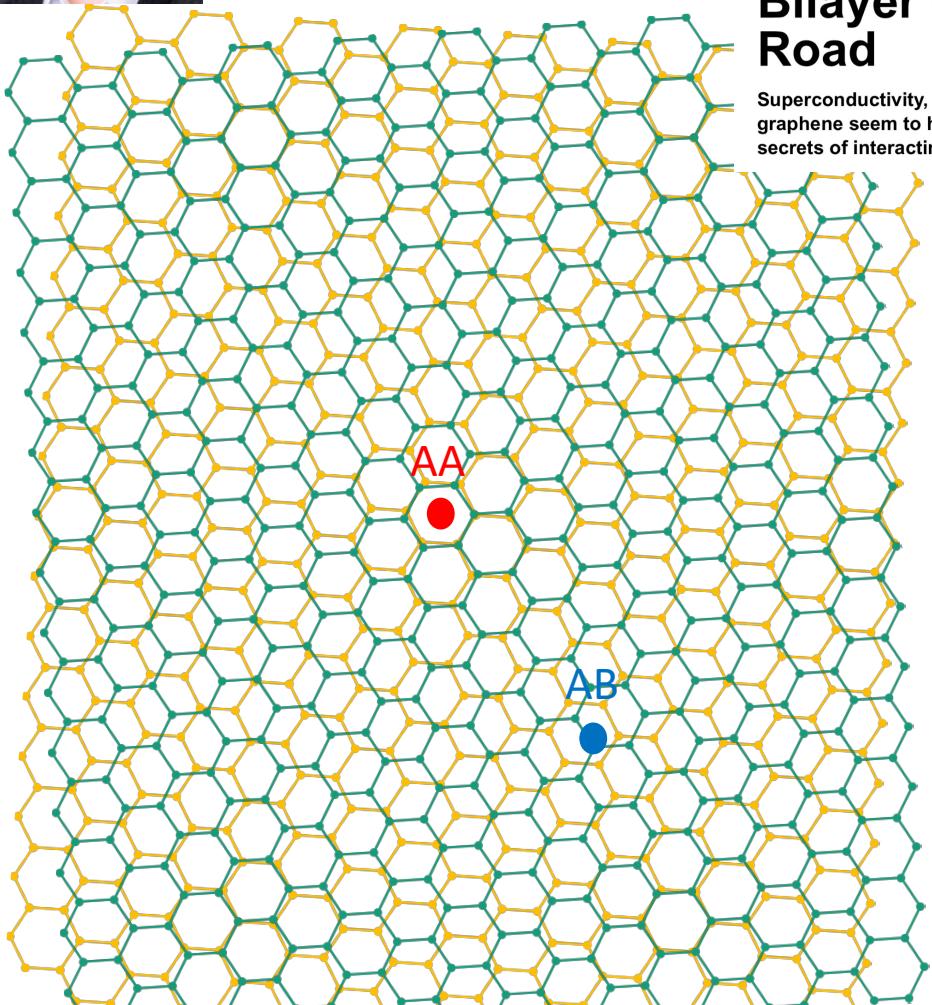
graphene – the twist

Physics

TREND

Bilayer Graphene's Wicked, Twisted Road

Superconductivity, magnetism, and other forms of interacting electron behavior—bilayers of graphene seem to have it all. Researchers are now using this pristine material to unlock the secrets of interacting-electron phenomena with unprecedented control and tunability.





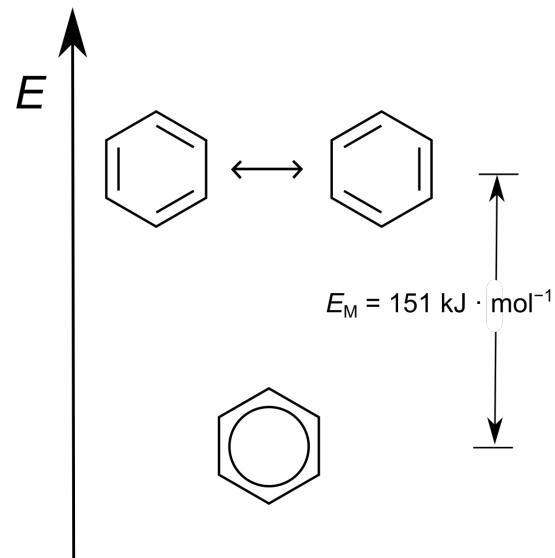
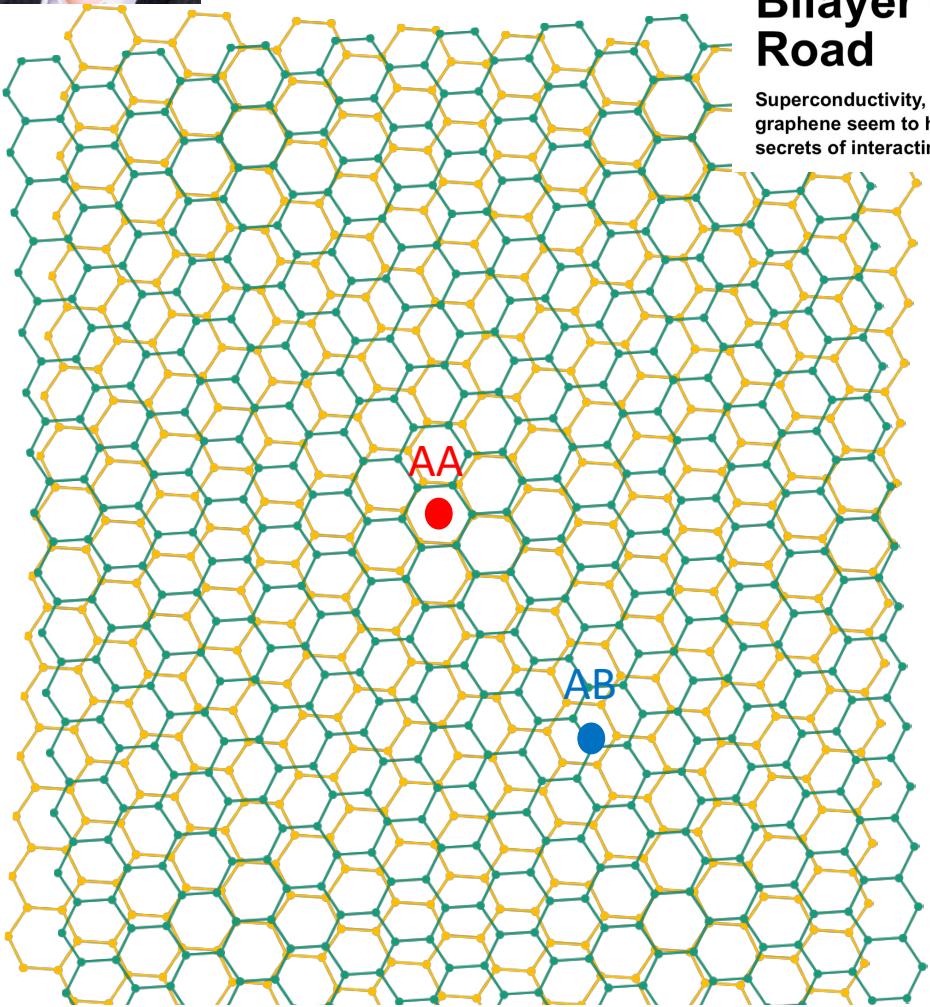
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Physics

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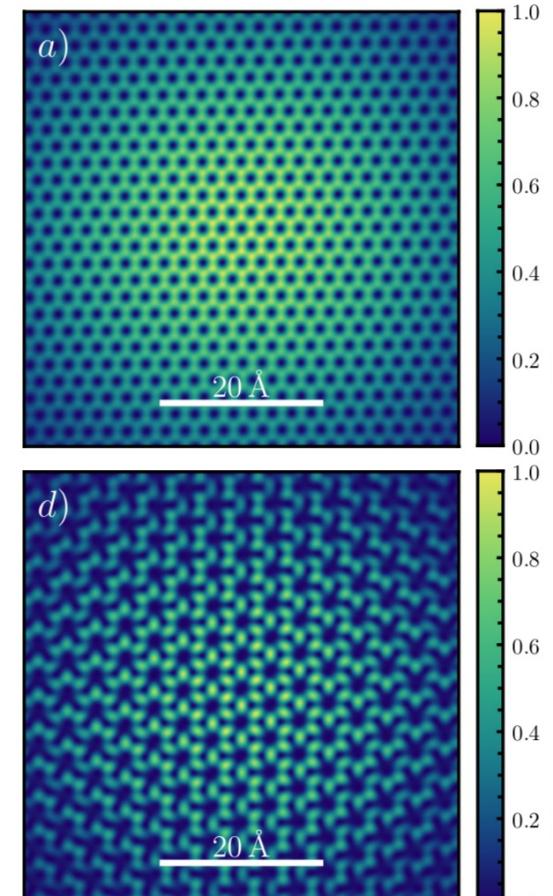


AA π -elektrons localised
AB π -electrons delocalised

Unconventional superconductivity in magic-angle ..., Y Cao, , et al., Nature 556 (2018), 43-50, Spectroscopy of Twisted Bilayer Graphene ..., Dumitru Călugăru et al., arXiv:2110.15300



Spectroscopy with STM





graphene – the twist

VIP Nanographenes Very Important Paper

International Edition: DOI: 10.1002/anie.201808178
German Edition: DOI: 10.1002/ange.201808178

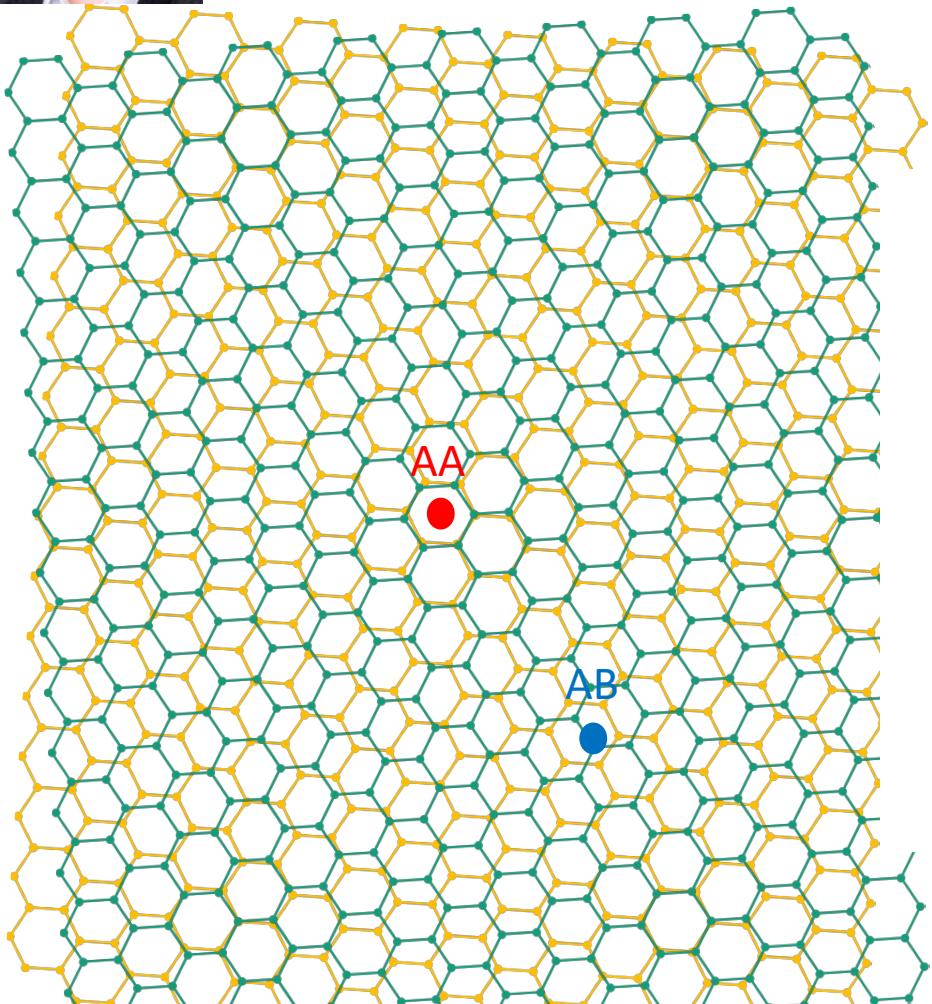
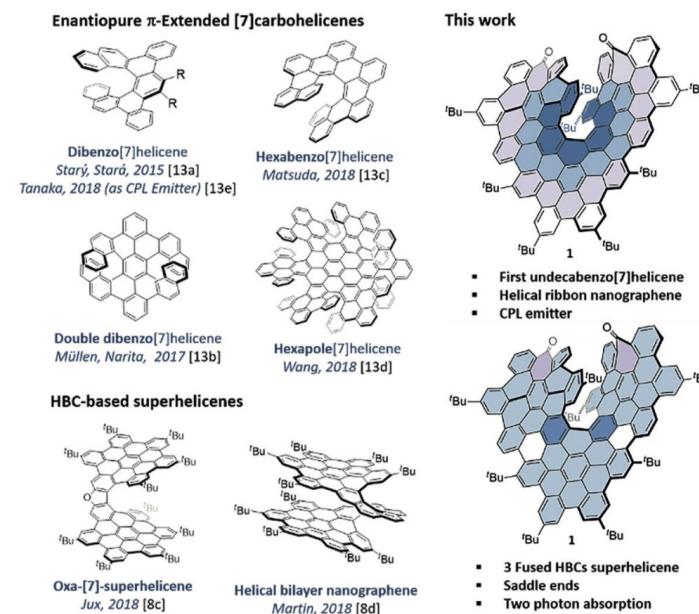


Figure 1. Background and novel structural features of compound 1.

twisted graphene – aromatic – antiaromatic ???



Unconventional superconductivity in magic-angle ..., Y Cao, , et al., Nature 556 (2018), 43-50, Spectroscopy of Twisted Bilayer Graphene ..., Dumitru Călugăru et al., arXiv:2110.15300

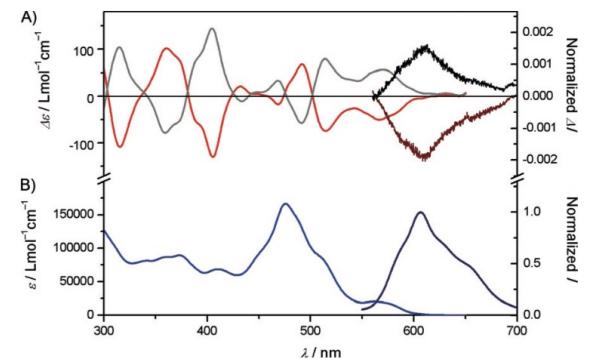
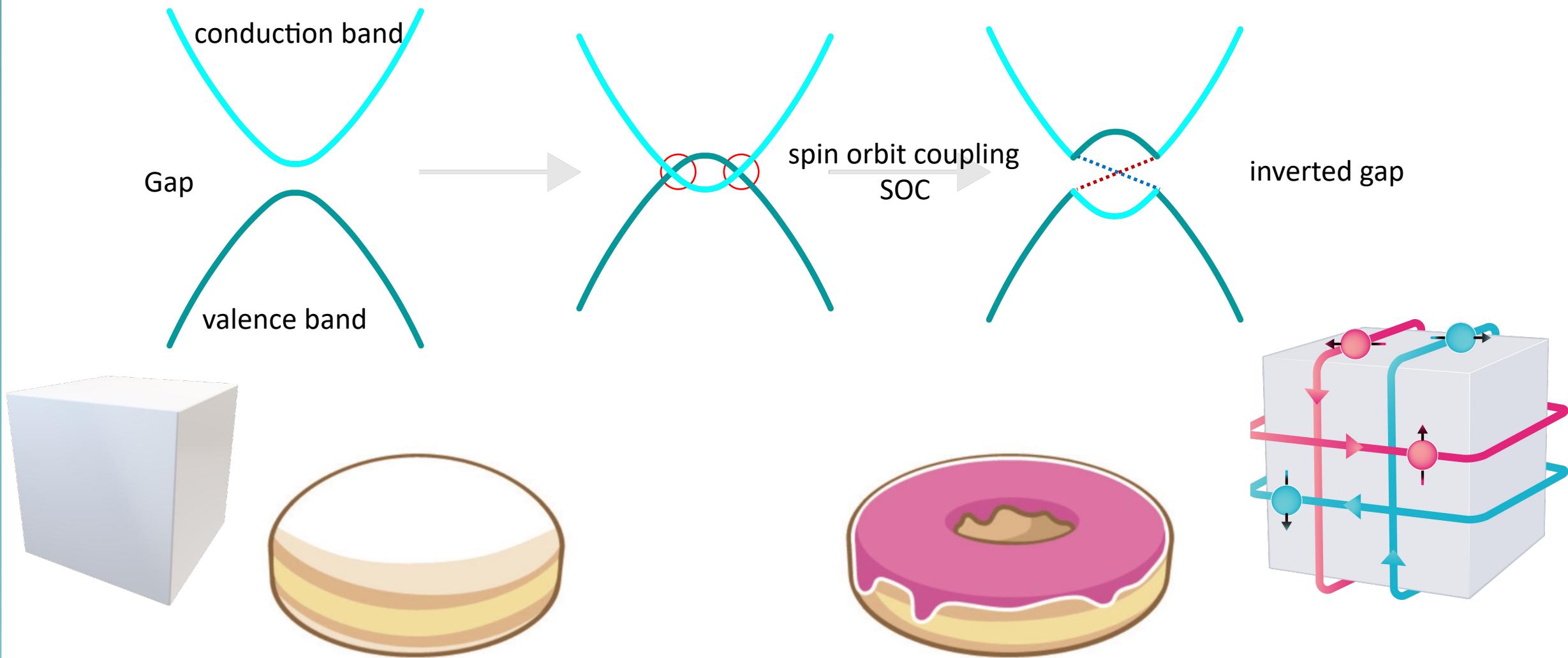
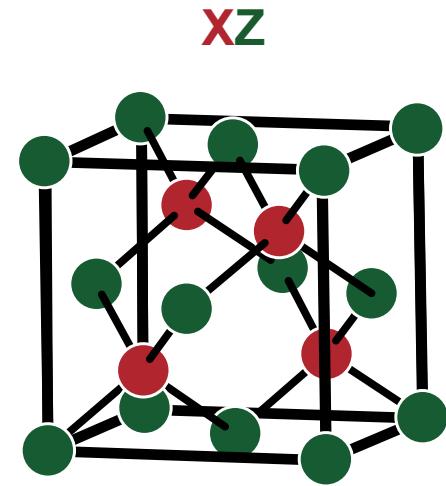
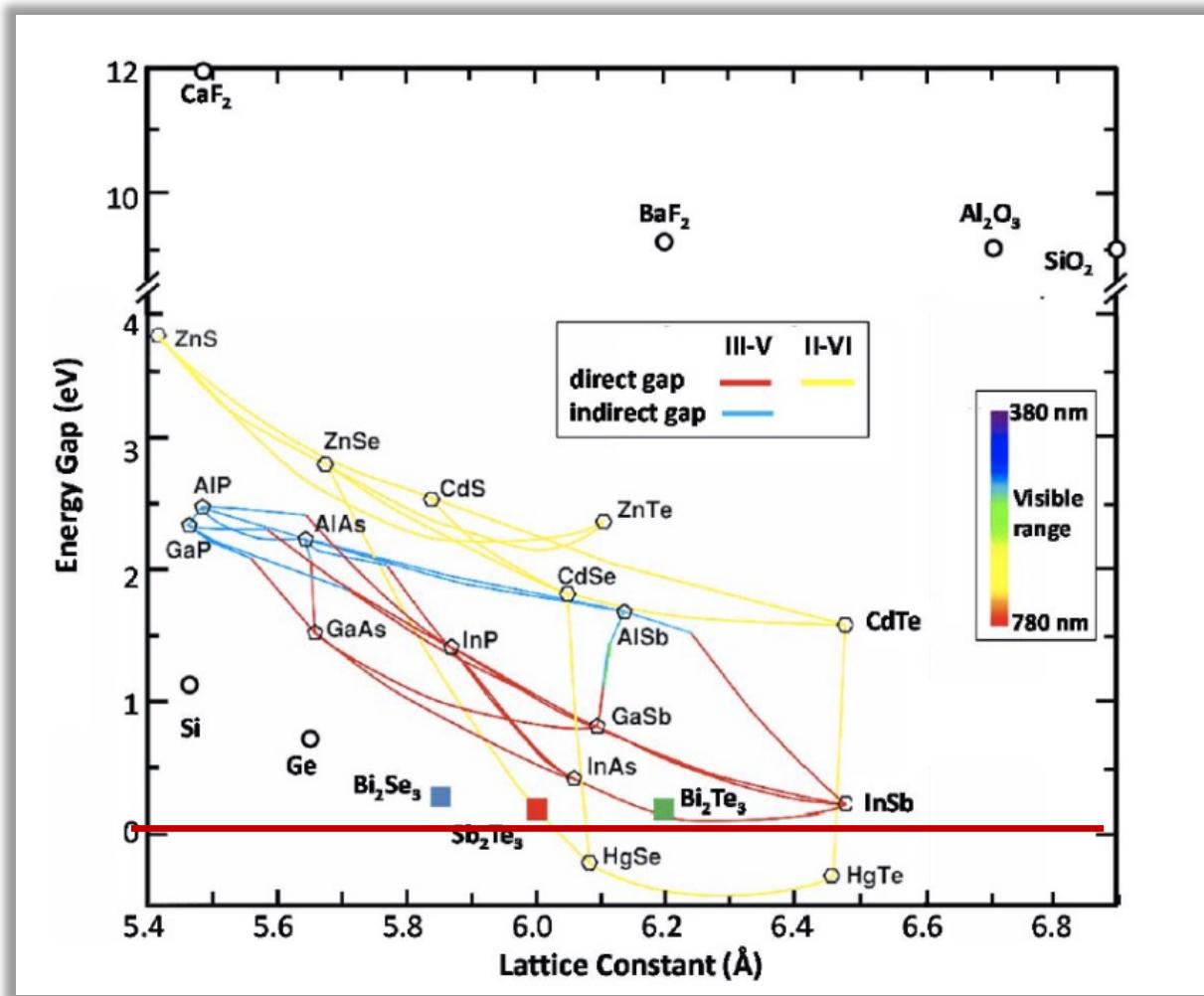
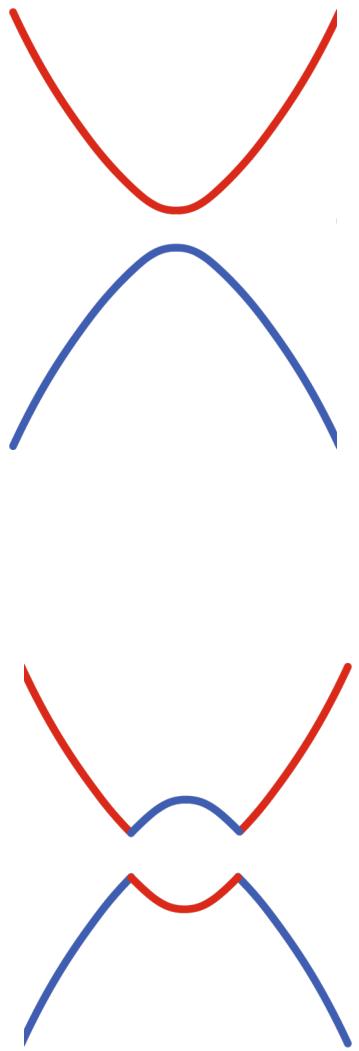


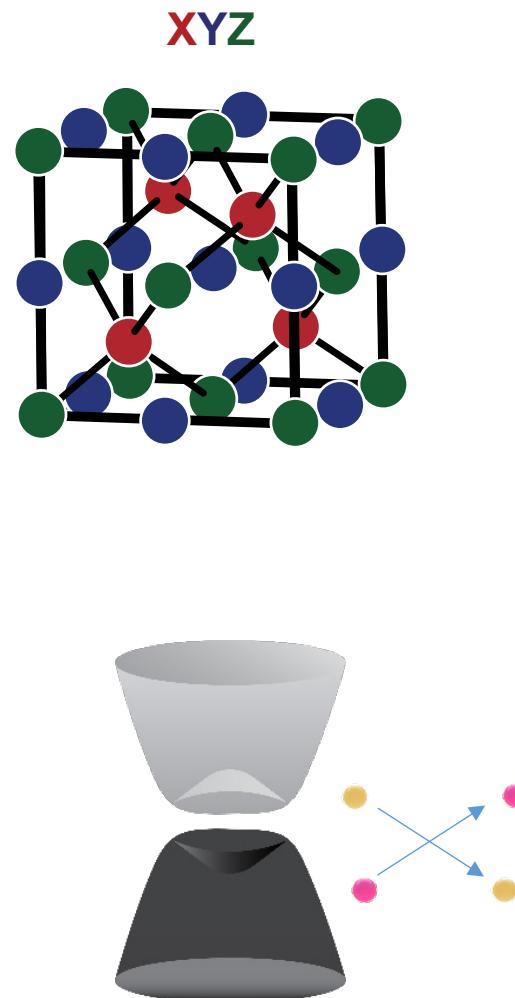
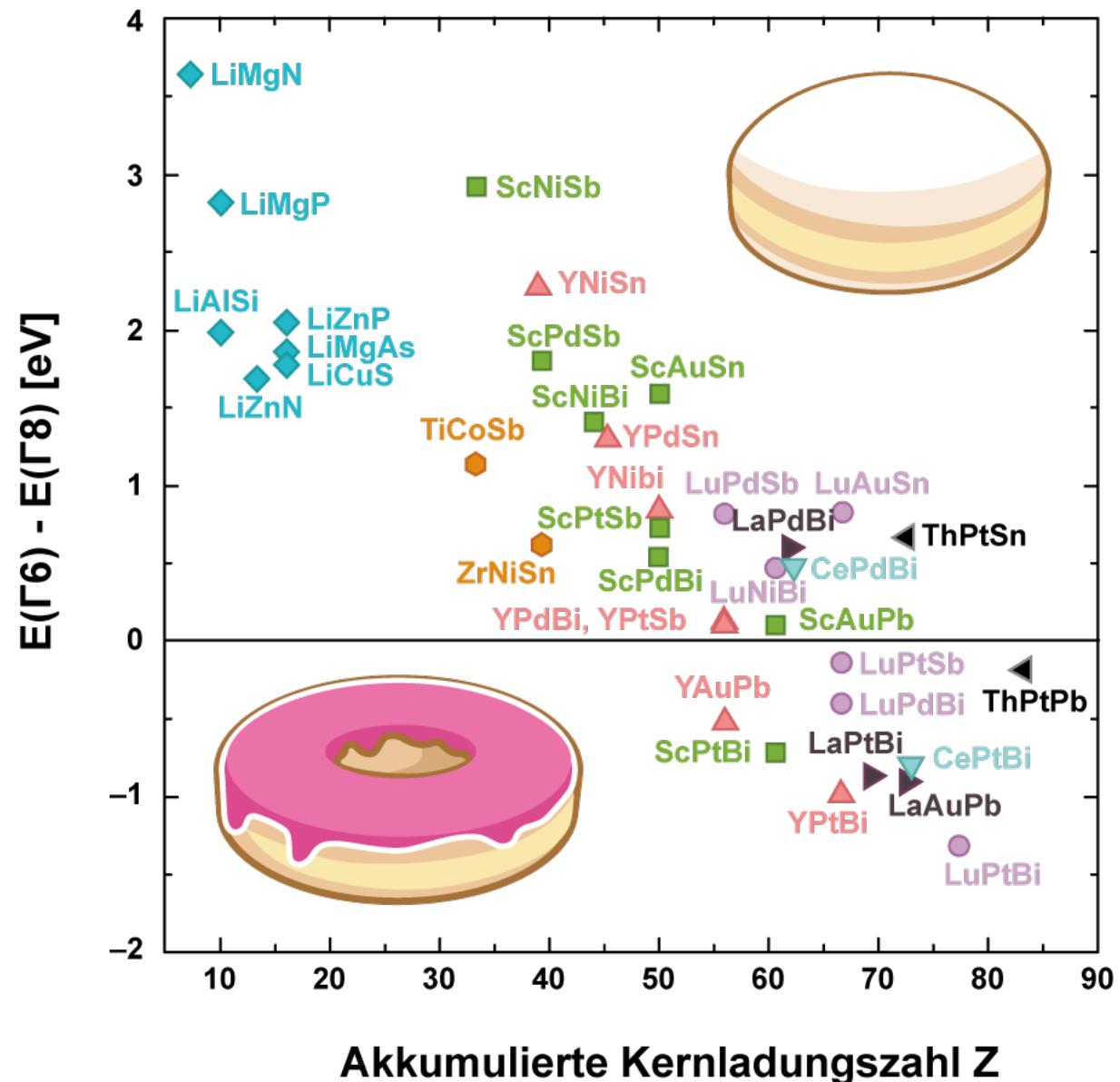
Figure 3. a) Experimental ECD (left) and CPL (right, $\lambda_{exc}=490$ nm) of M) (red and scarlet) and (P) (gray and black) enantiomers of 1 in CH_2Cl_2 at ca. 5×10^{-6} M. b) Experimental UV/Vis (blue) and fluorescence (navy, $\lambda_{exc}=490$ nm) spectra of 1 in CH_2Cl_2 at ca. 5×10^{-6} M.

topological insulators

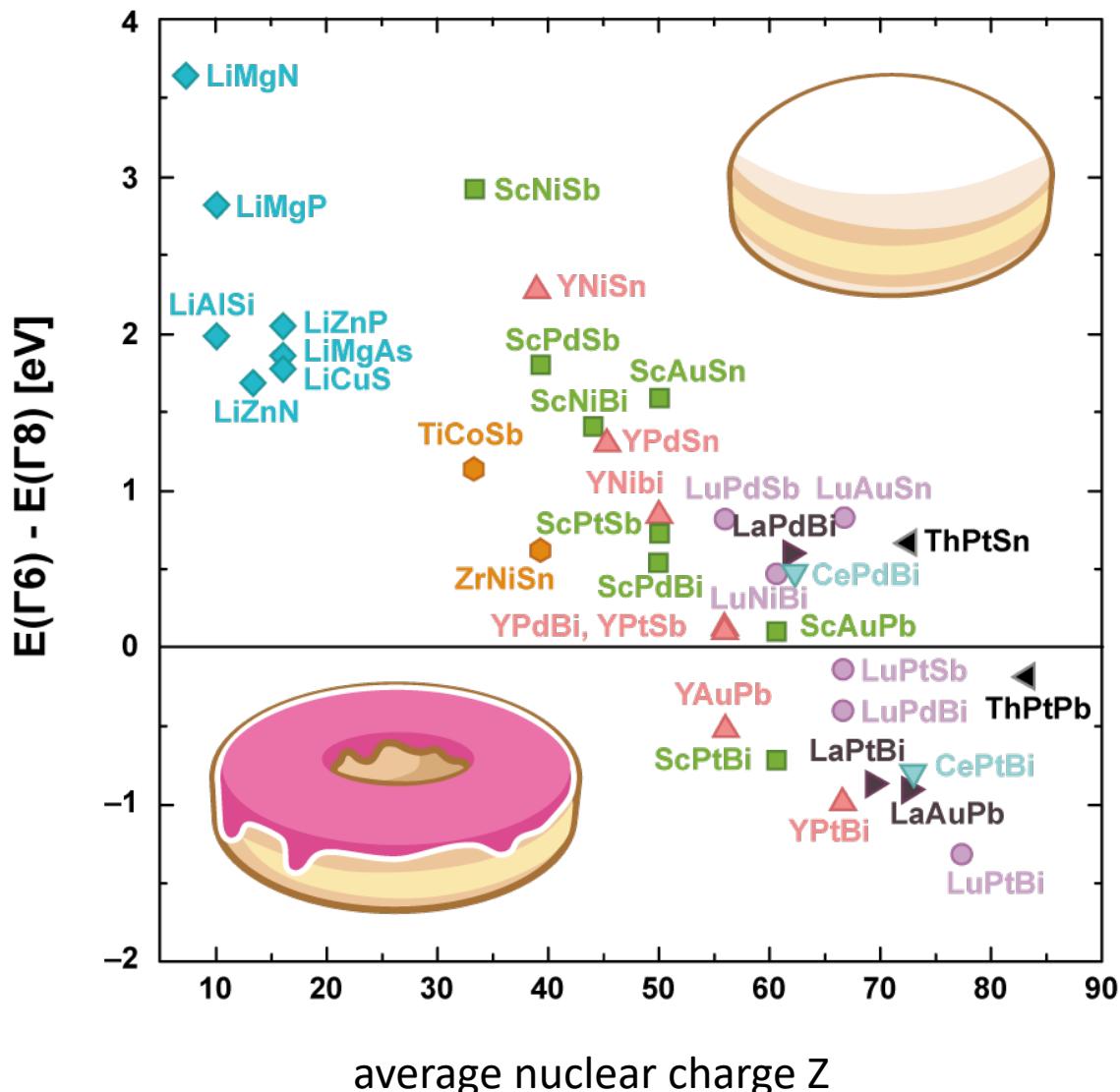


semiconductor

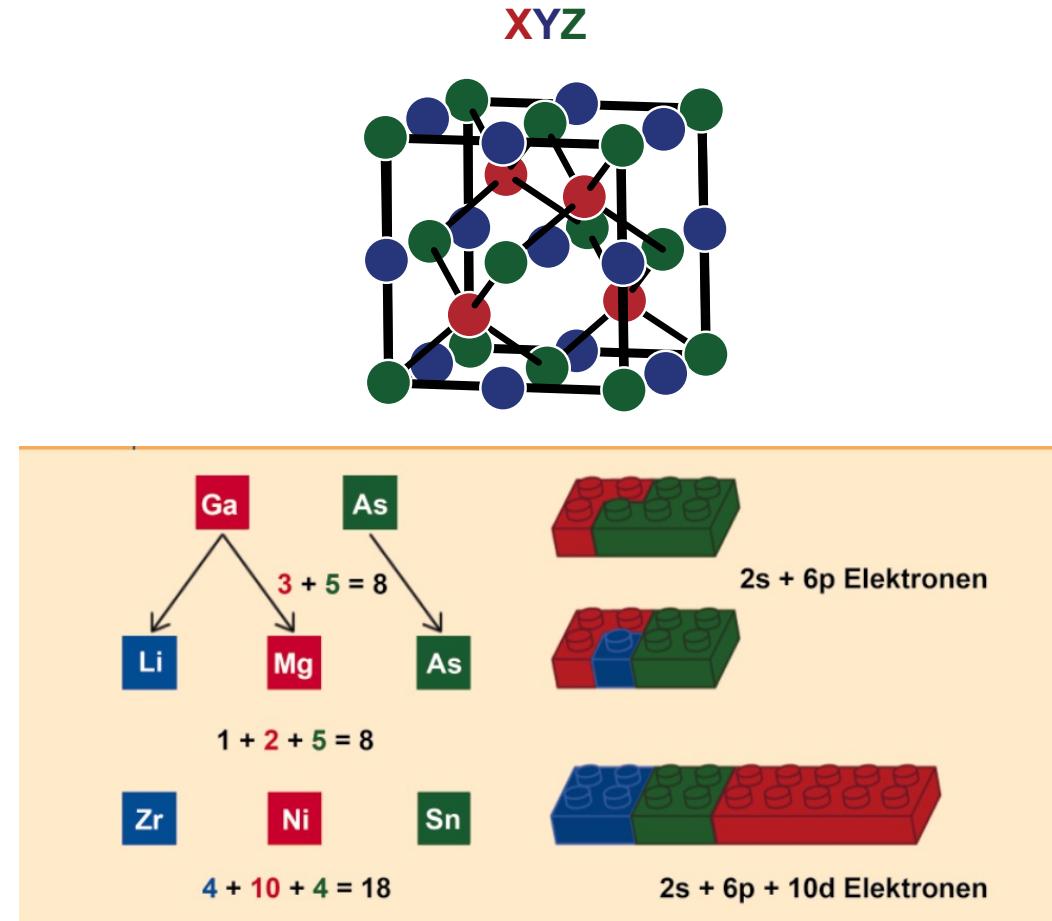




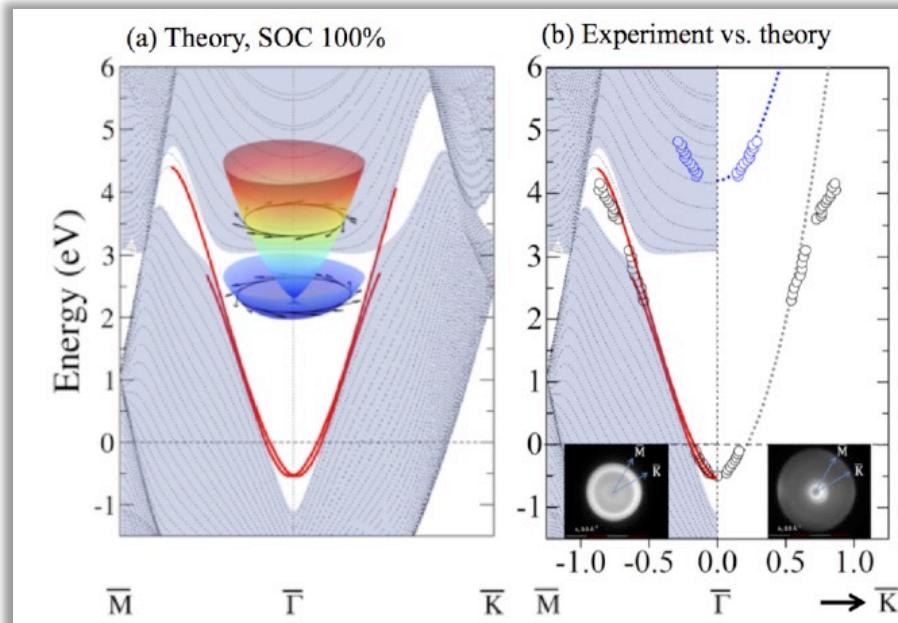
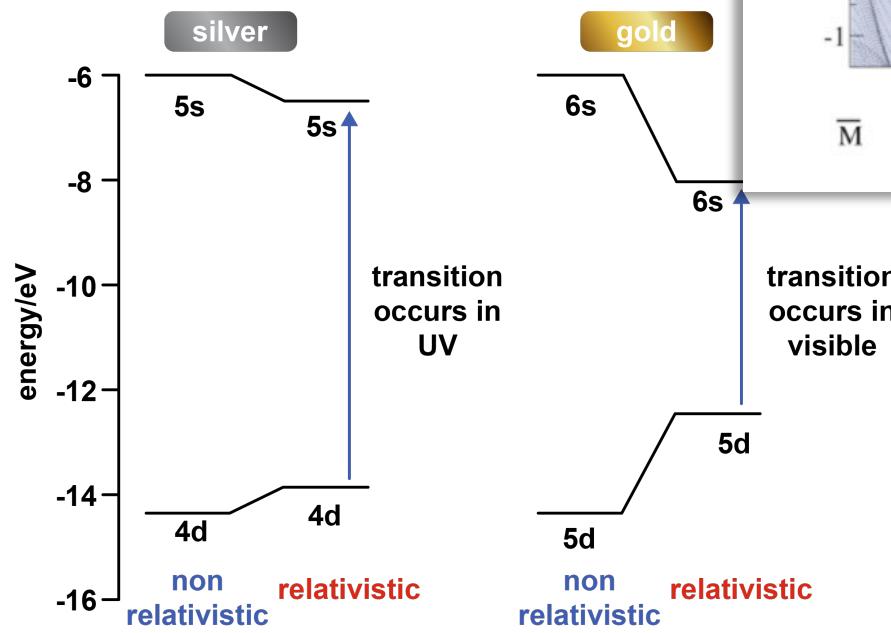
ternary semiconductor



half Heusler compounds are ternary semiconductors

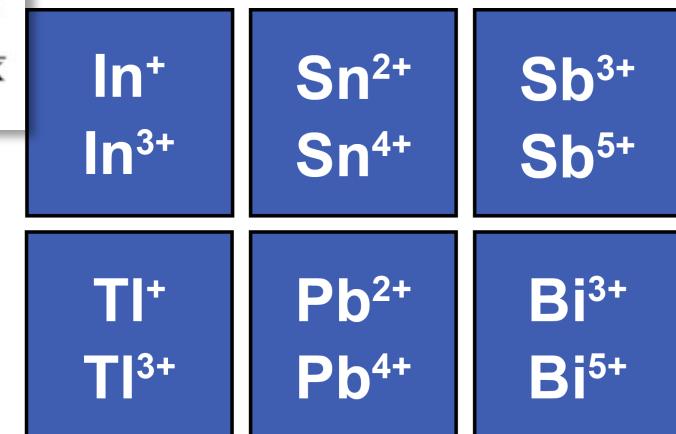


gold: a topological metal

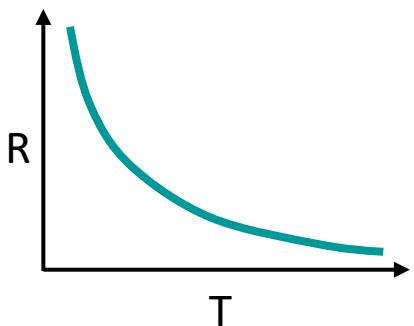


Inert pair effect

contraction of the 1s-orbital =
higher s orbitals will be „core“ like



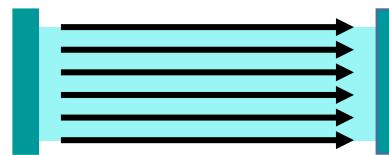
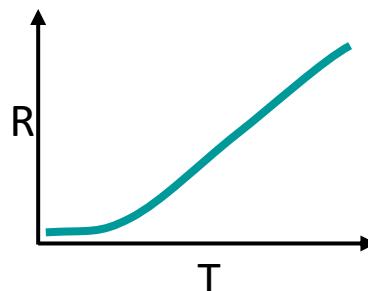
measurements



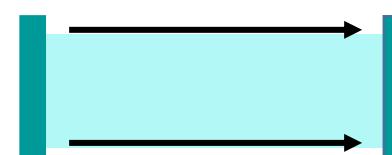
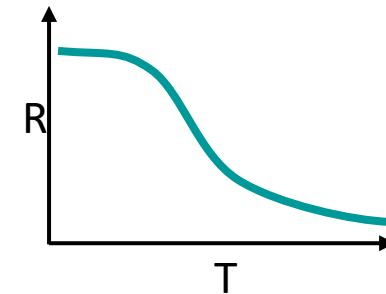
copyright Shekhar Chandra et al.



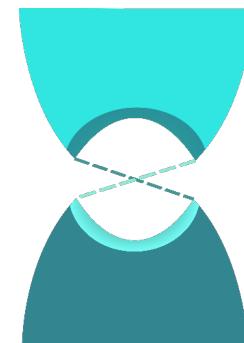
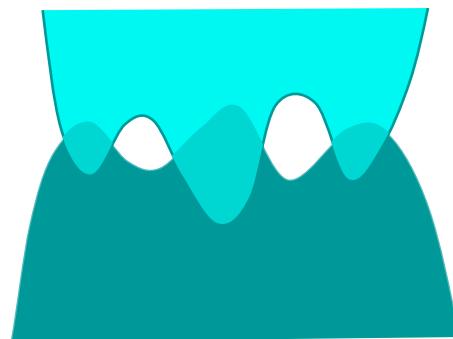
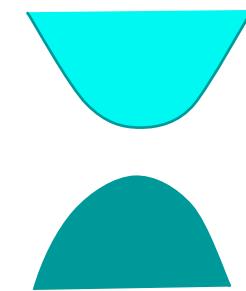
insulator



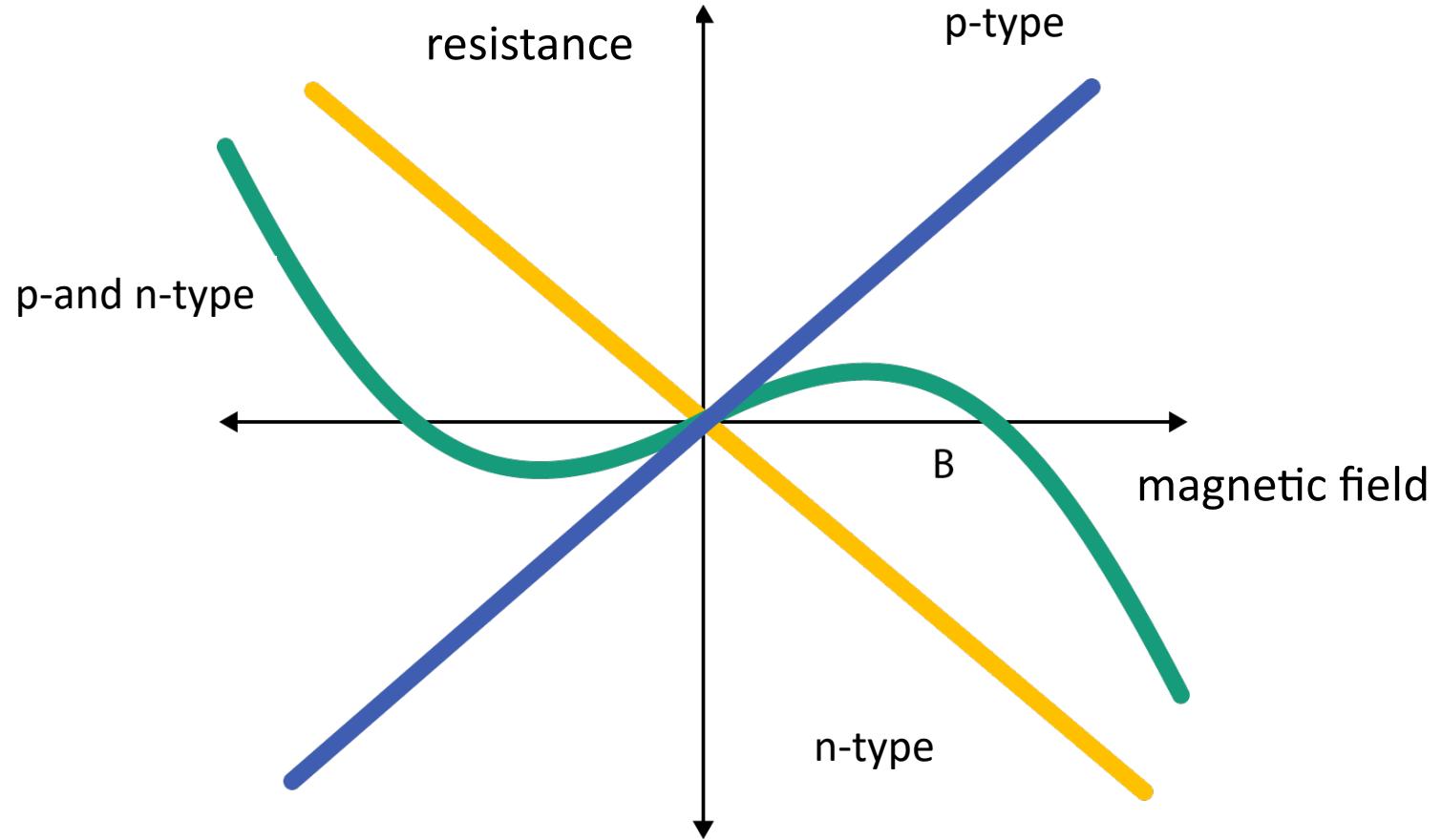
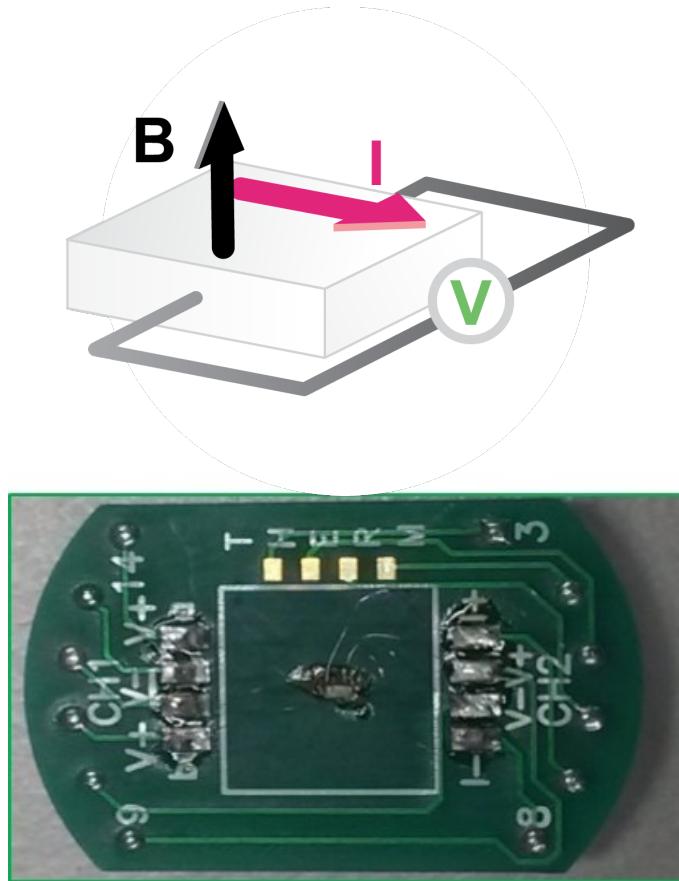
metal



topological insulator



Hall-measurements of semiconductors



quantum Hall effect

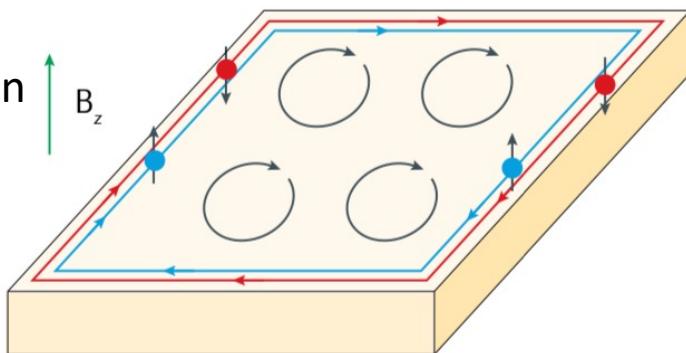
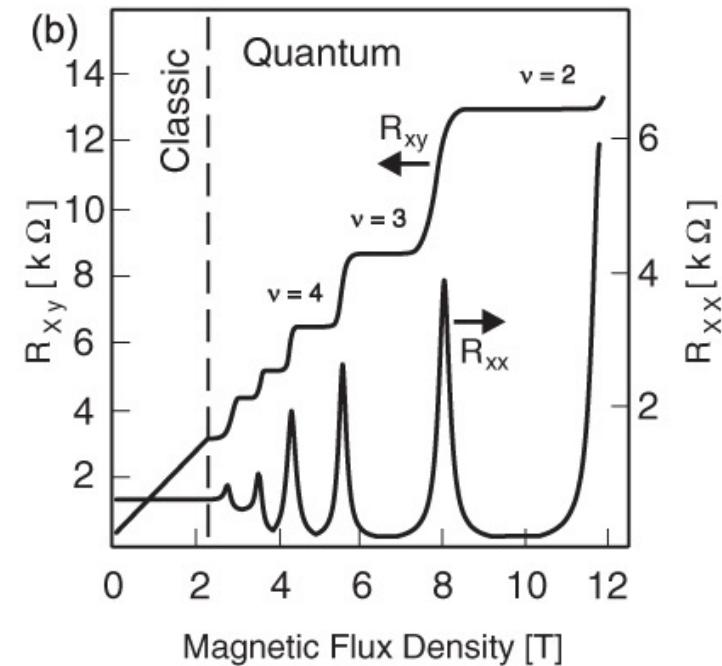
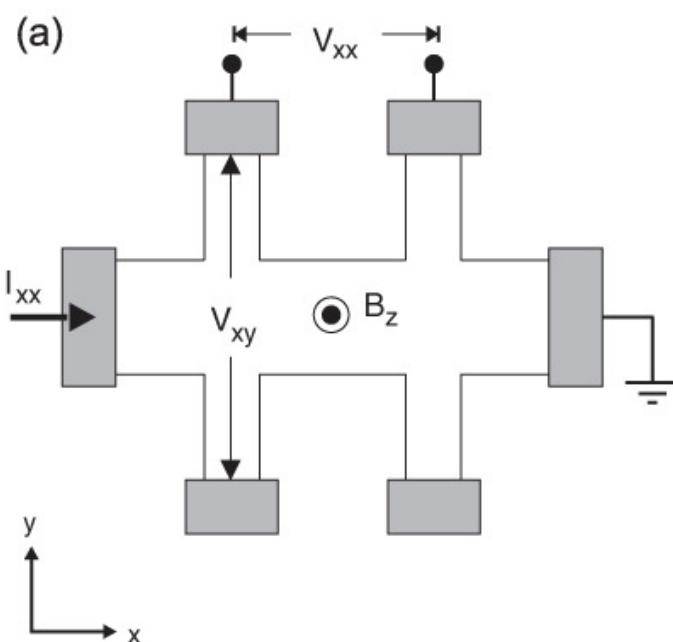
or the beginning



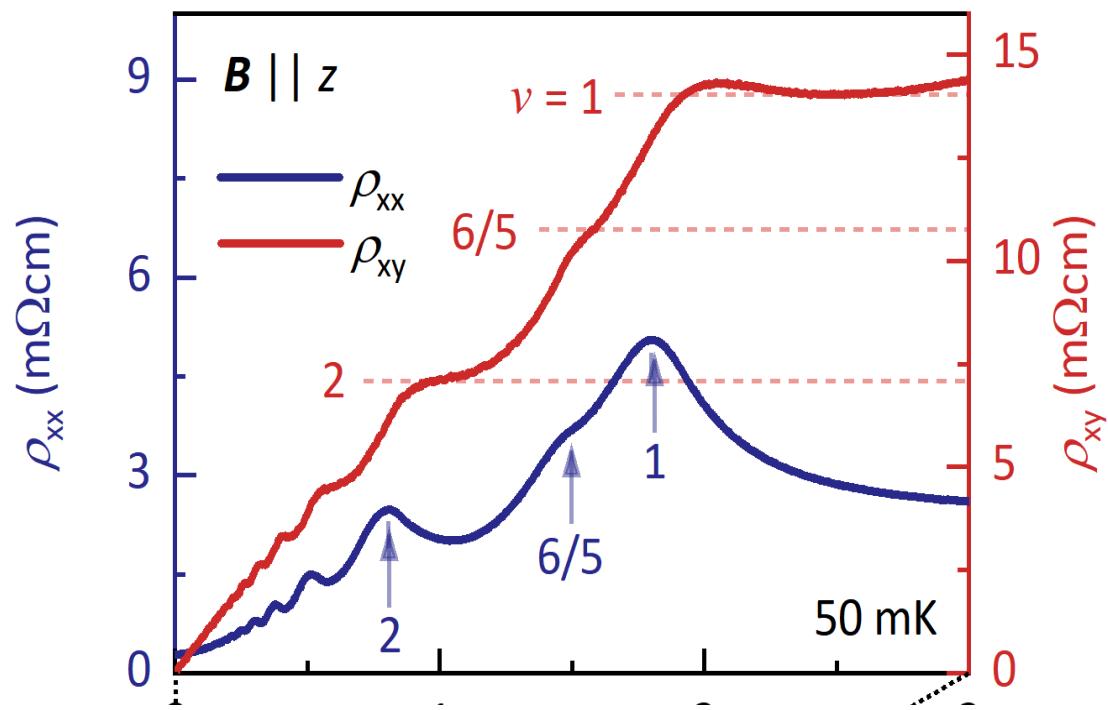
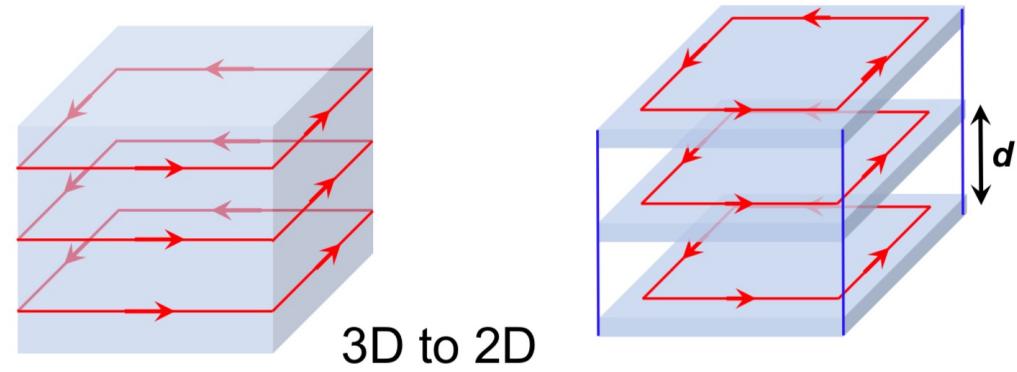
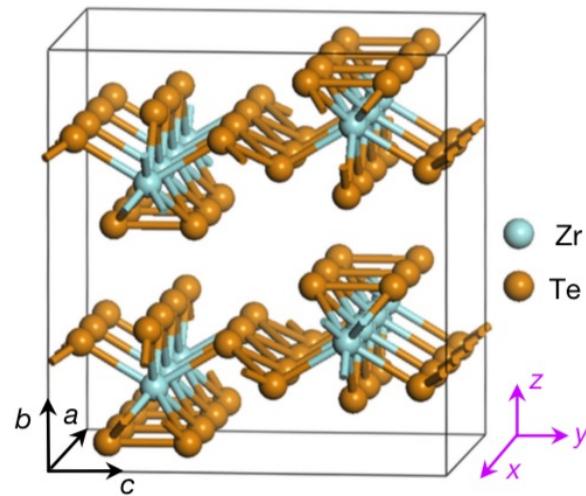
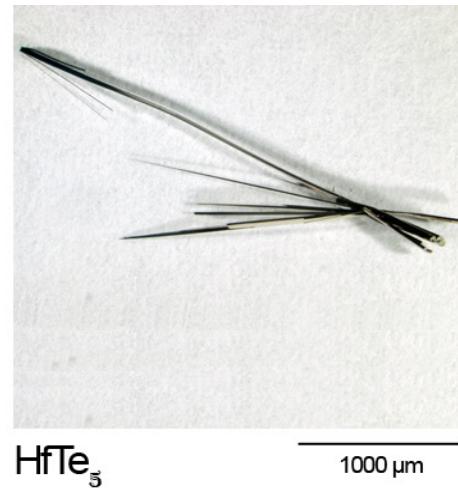
Klaus von Klitzing
©mpg

quantum Hall effect:

- exotic effect
- in complex layer systems of semiconductors
- precisely adjusted charge carrier concentration
- high magnetic fields
- low temperature



quantum Hall effect in a crystal





VOLUME 61, NUMBER 18

PHYSICAL REVIEW LETTERS

31 OCTOBER 1988

**Model for a Quantum Hall Effect without Landau Levels:
Condensed-Matter Realization of the “Parity Anomaly”**

F. D. M. Haldane

*Department of Physics, University of California, San Diego, La Jolla, California 92093
(Received 16 September 1987)*



Topological Quantum Matter

Nobel Lecture, December 8, 2016

Prediction of the quantum anomalous Hall effect without magnetic field
Solution: Magnetic materials

Haldane, PRL 61, 2015 (1988)

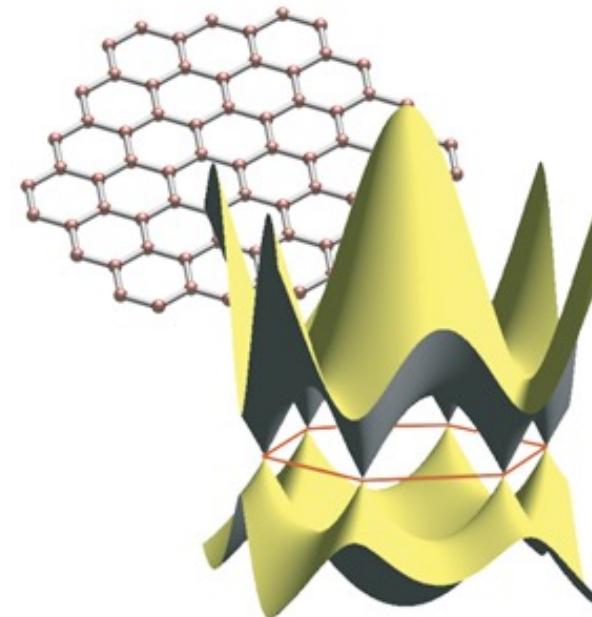
quantum anomalous Hall – quantum spin Hall



Z_2 Topological Order and the Quantum Spin Hall Effect

C. L. Kane and E. J. Mele

*Department of Physics and Astronomy, University of Pennsylvania, Philadelphia, Pennsylvania 19104, USA
(Received 22 June 2005; published 28 September 2005)*



Prediction of the quantum spin Hall effect
Solution: materials with large spin orbit coupling

Kane and Mele, PRL 95, 146802 (2005)

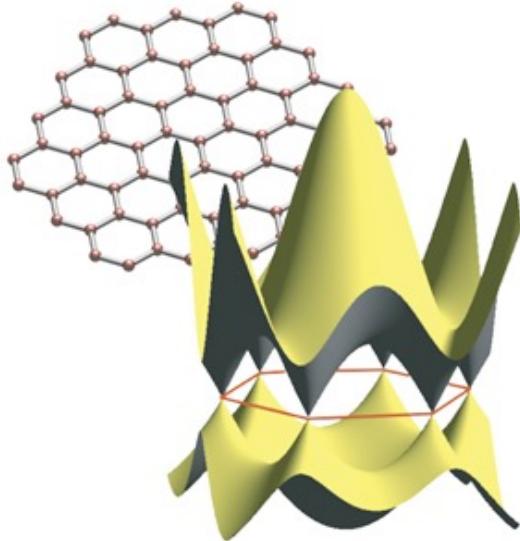


topological insulators – quantum spin Hall

Z₂ Topological Order and the Quantum Spin Hall Effect

C. L. Kane and E. J. Mele

Department of Physics and Astronomy, University of Pennsylvania, Philadelphia, Pennsylvania 19104, USA



Prediction of the quantum spin Hall effect

Solution: materials with large spin orbit coupling

Kane and Mele, PRL 95, 146802 (2005)

Heavy insulating elements?

Strained α -Sn and Bi-bilayer

$\lambda_{\text{soc}} \sim Z^2$ for valence shells



topological insulators

Element Bismuth
Bi-Sb alloys
 Bi_2Se_3 and related structures



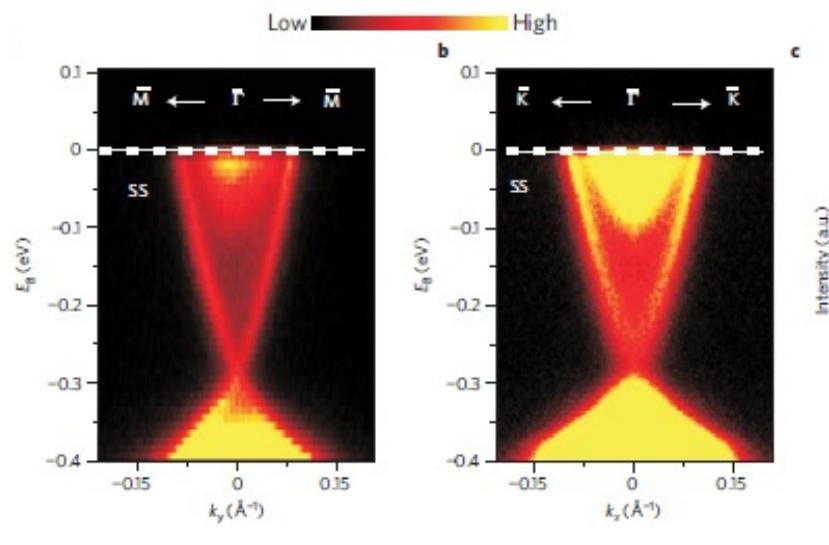
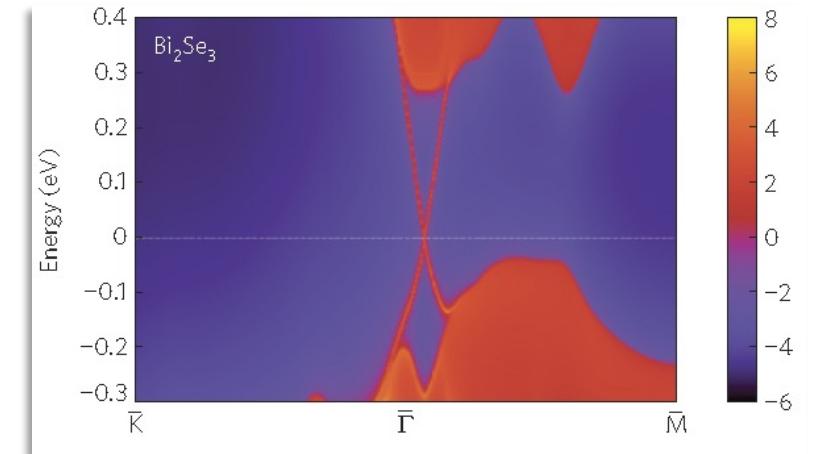
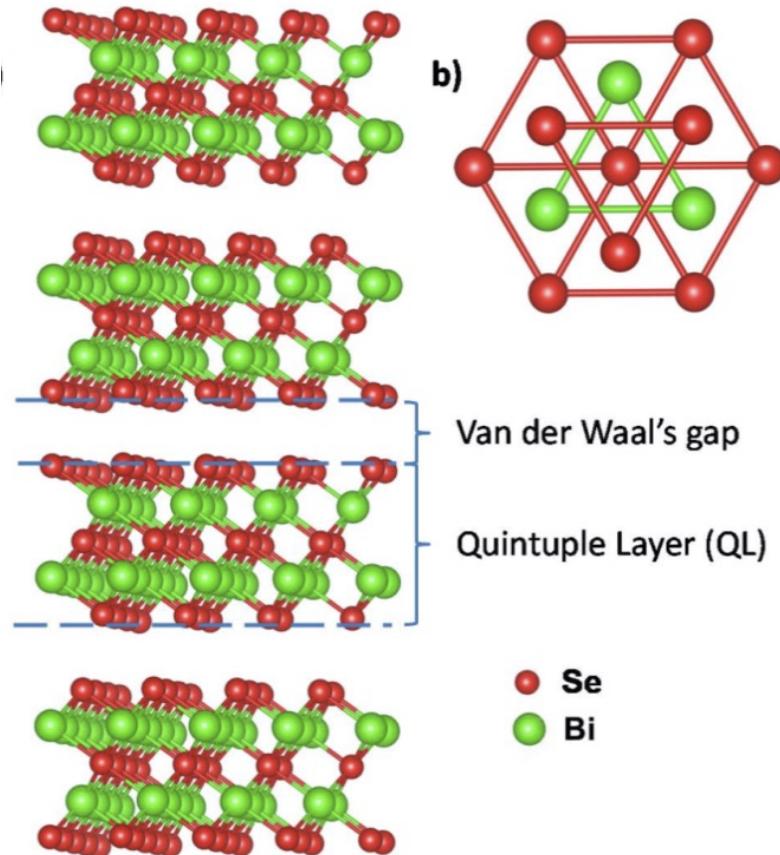
Moore and Balents, PRB 75, 121306(R) (2007)

Fu and Kane, PRB 76, 045302 (2007)

Murakami, New J. Phys. 9, 356 (2007)

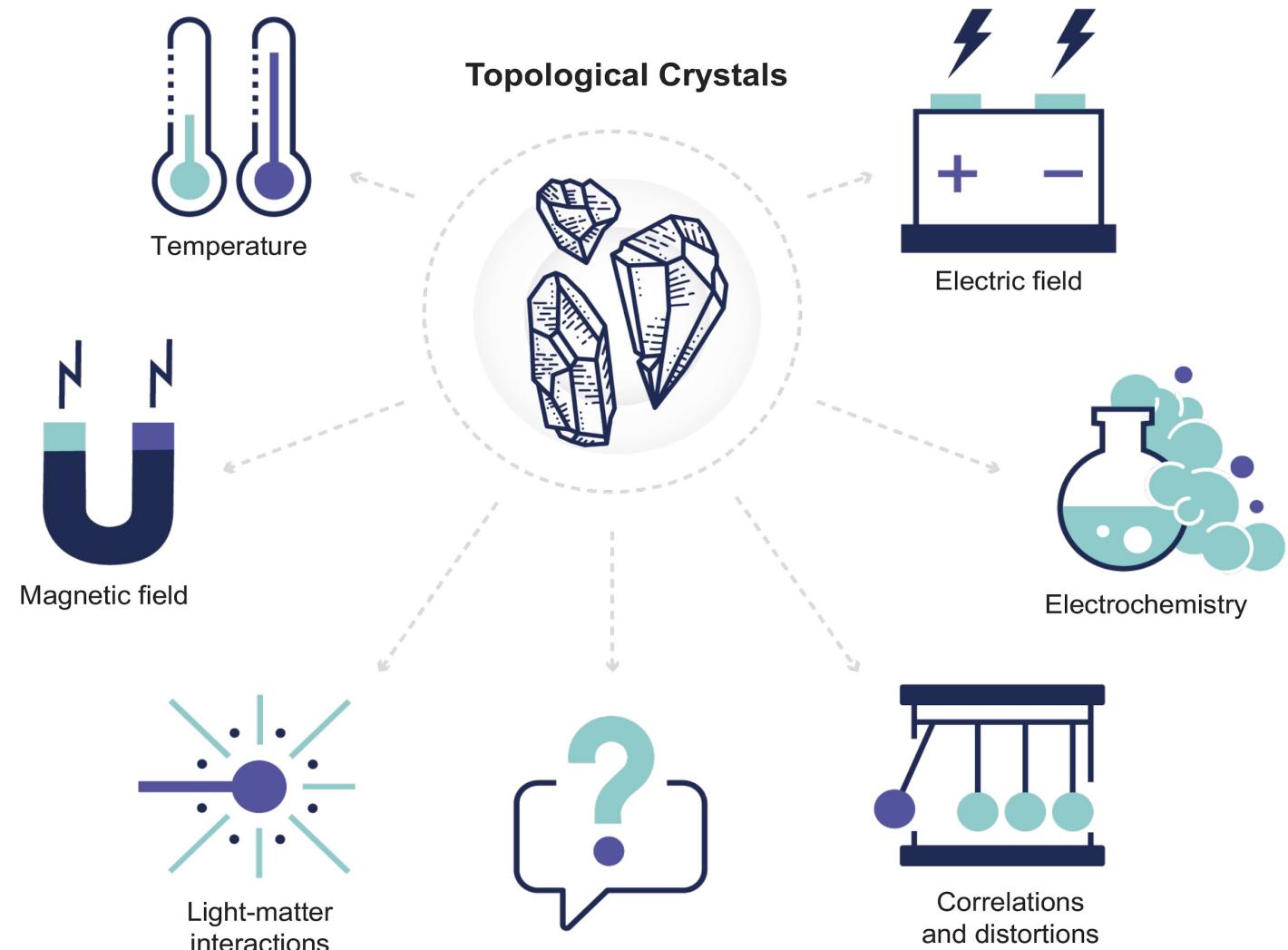
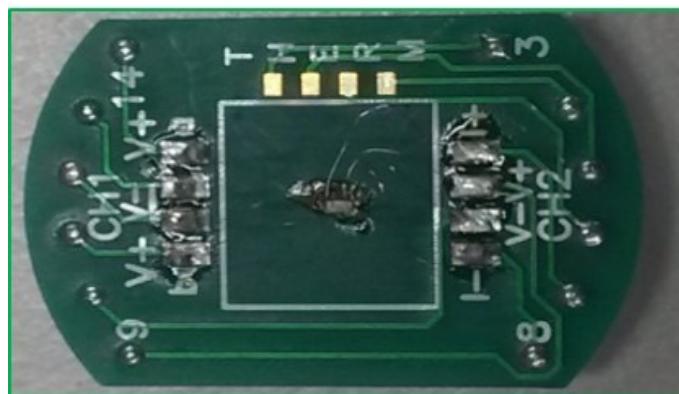
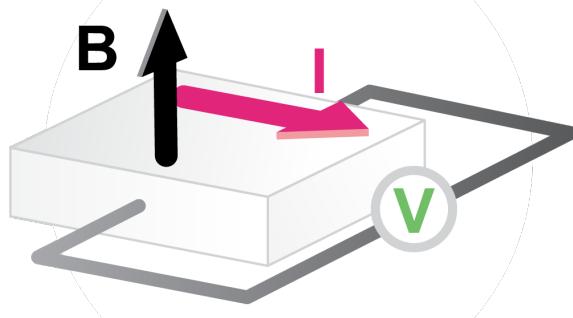
Hsieh, et al., Science 323, 919 (2009)

Xia, et al., Nature Phys. 5, 398 (2009); Zhang, et al., Nature Phys. 5, 438 (2009)



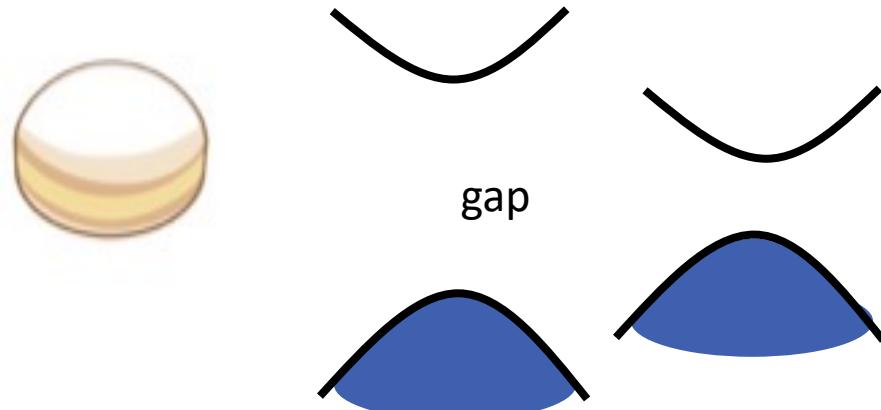
measurements

- protected surface states
- quantized effects
- large effects in response to
 - the magnetic field, ...



topologic or trivial

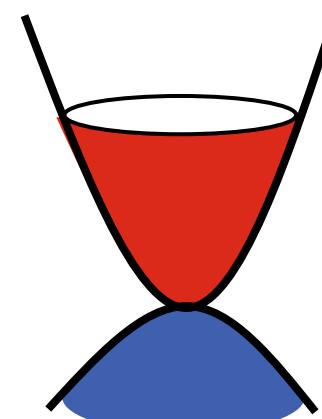
Insulator semiconductor



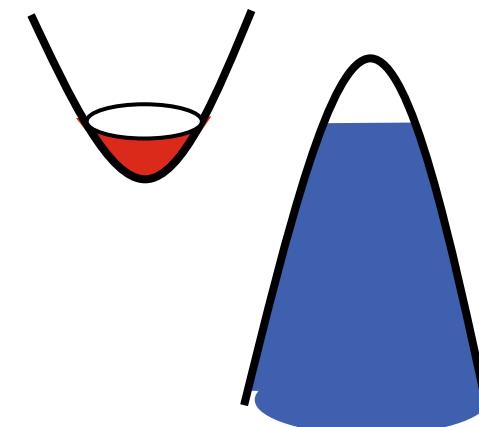
topological Insulator



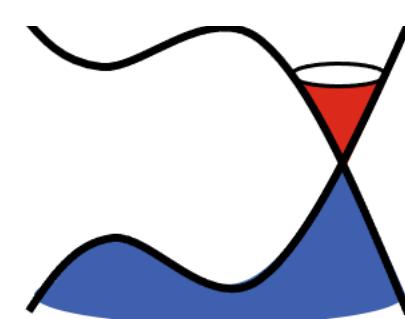
metal



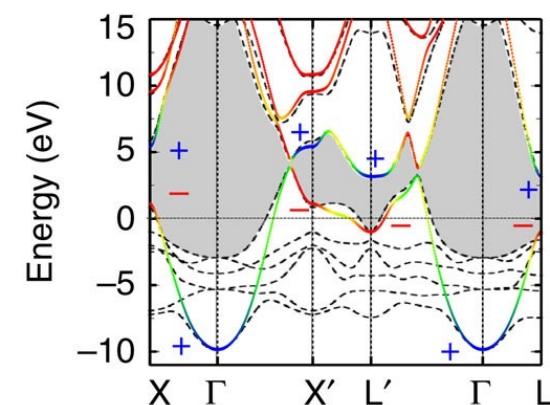
semimetal



Dirac semimetal

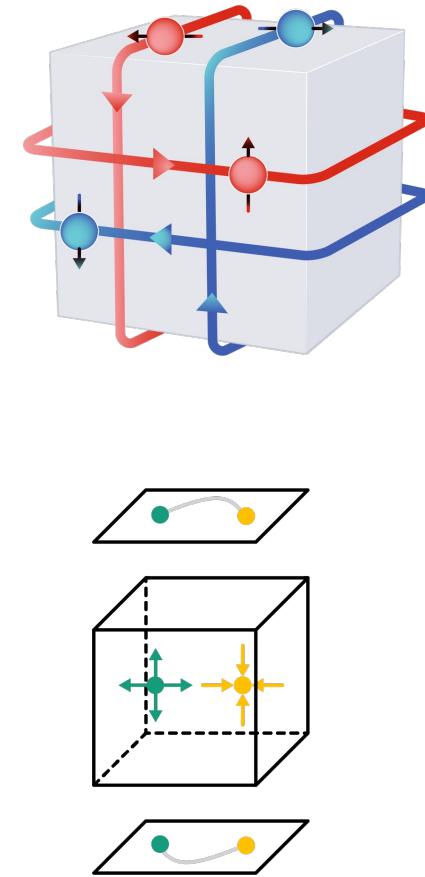
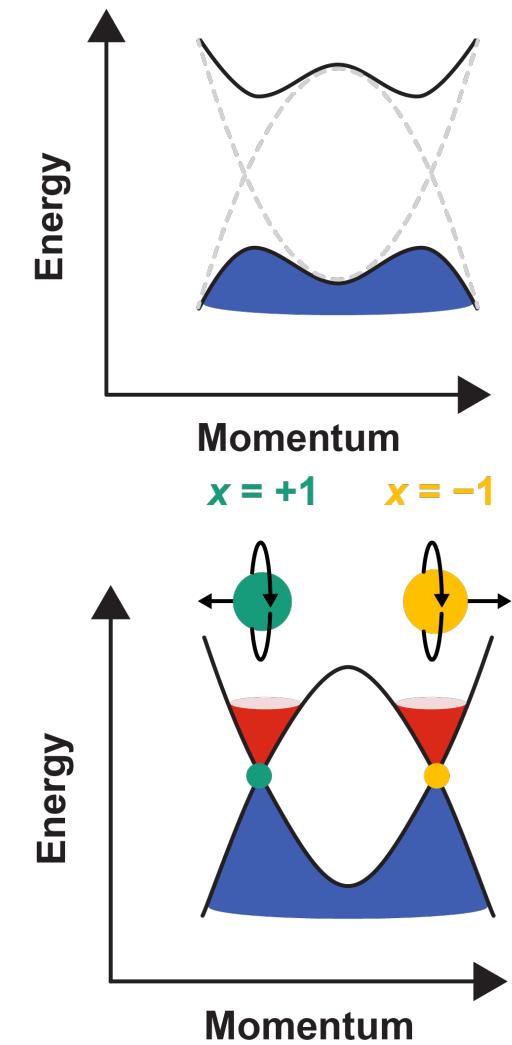
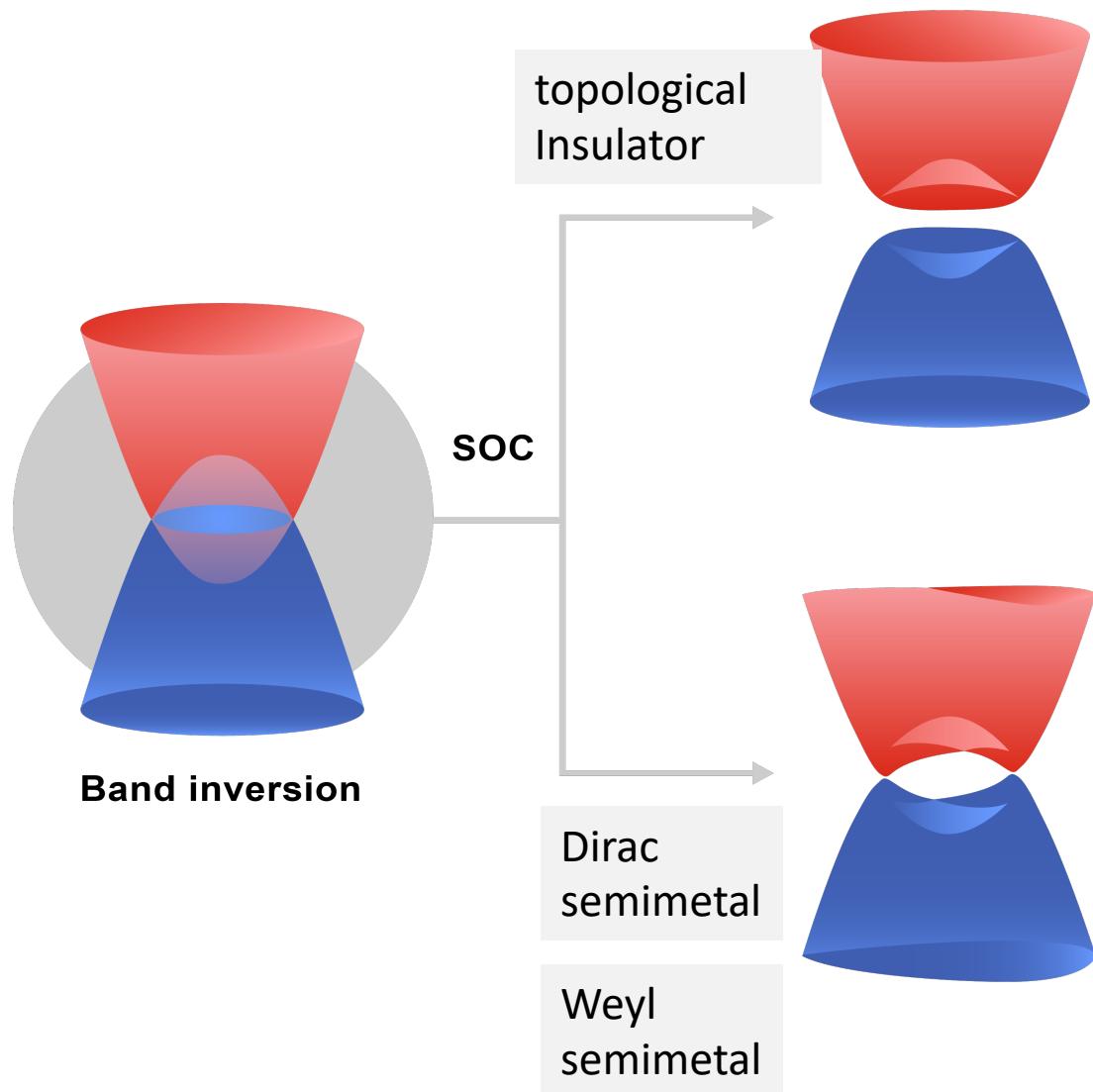


gold

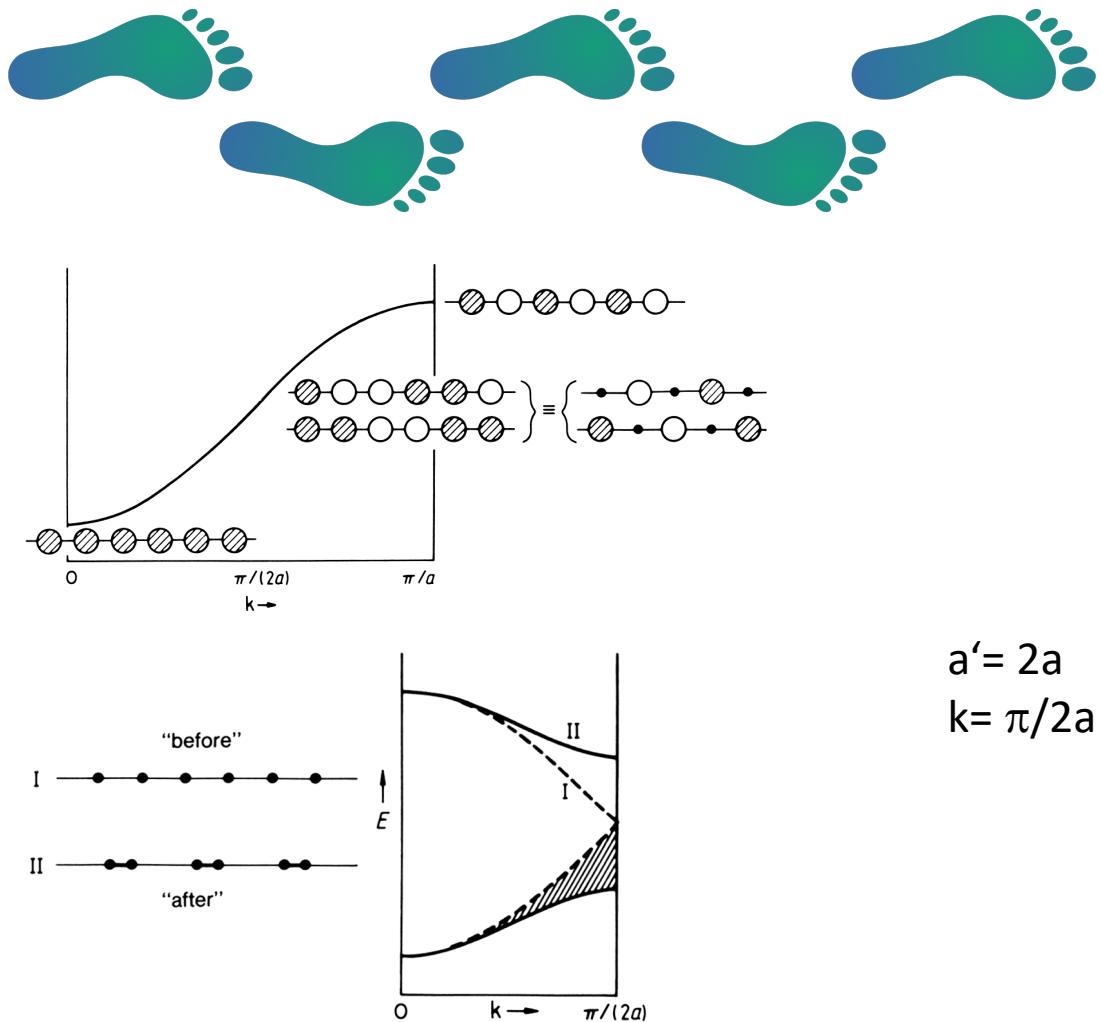


Graphene

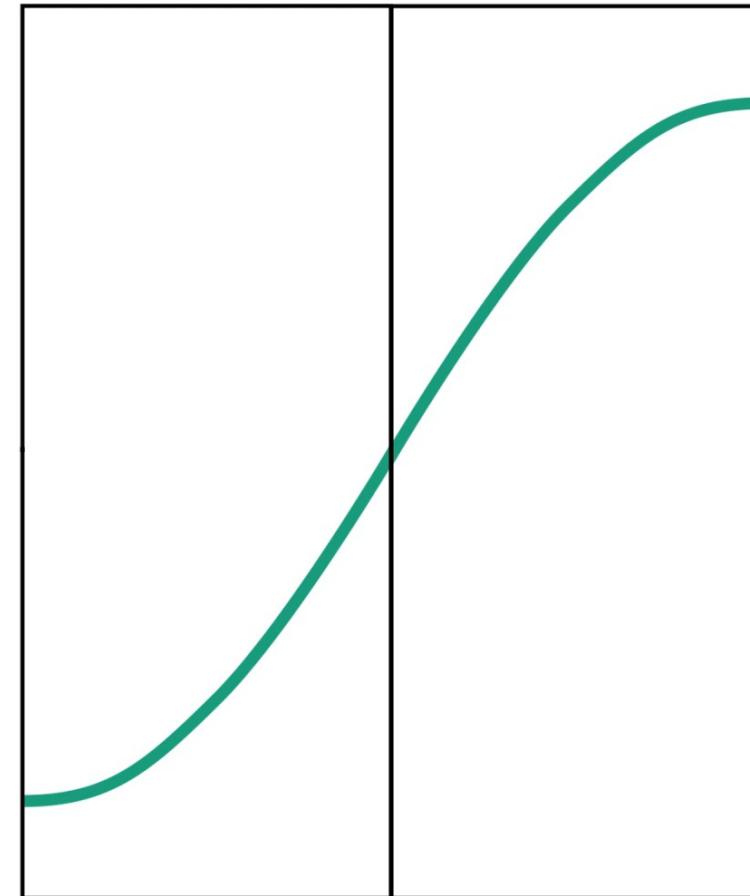
from insulator to semimetal



semimetals are common



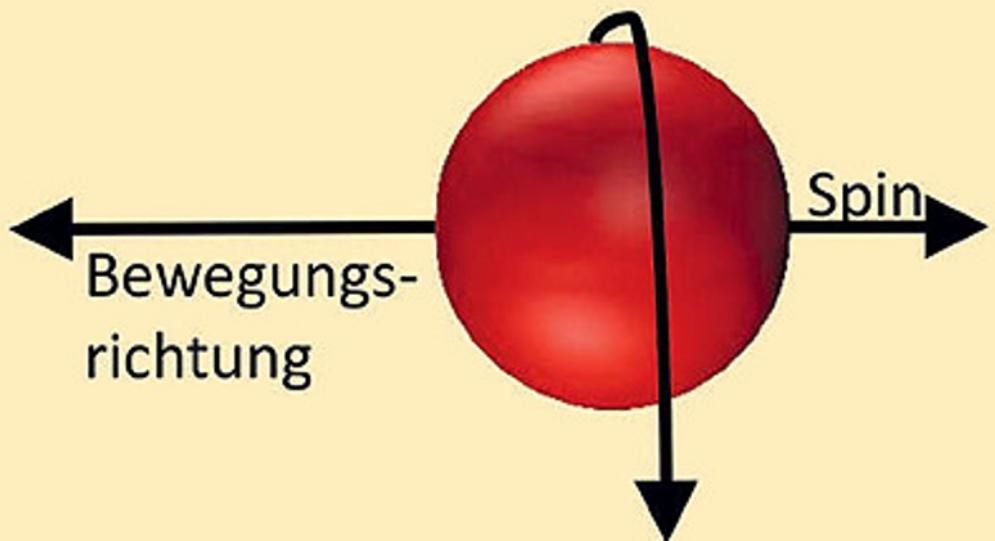
$$a' = 2a$$
$$k = \pi/2a$$



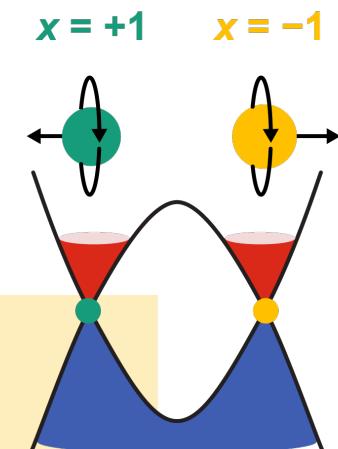
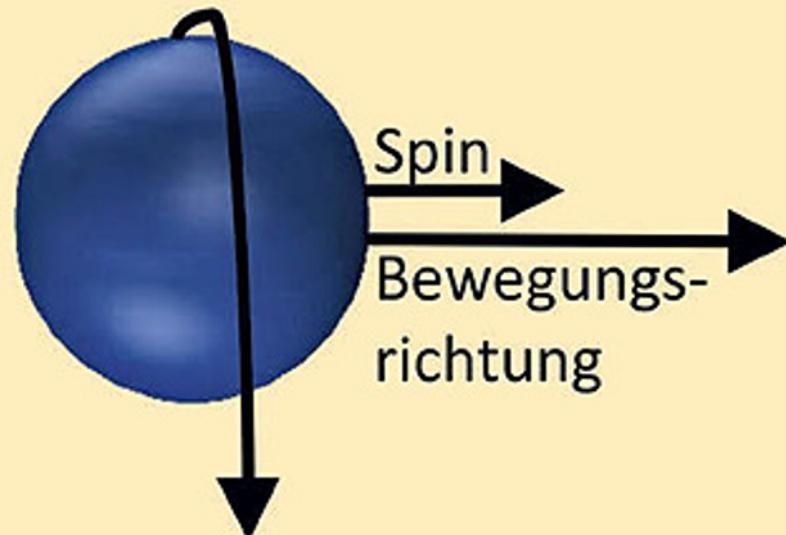
chiral electrons

in bulk

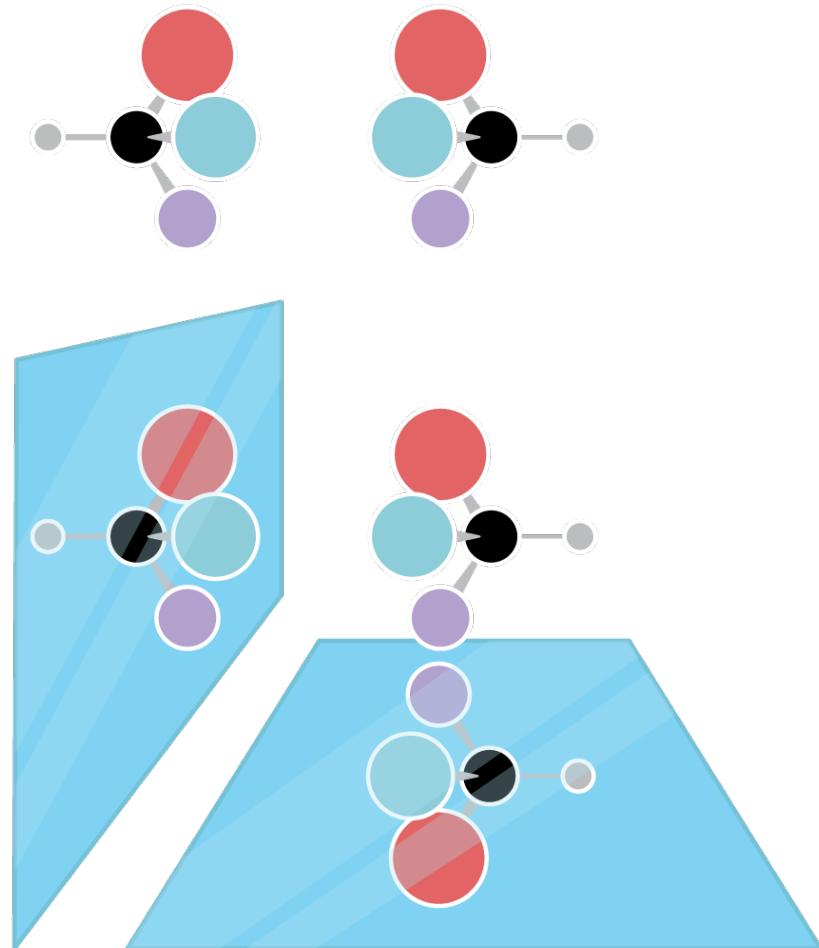
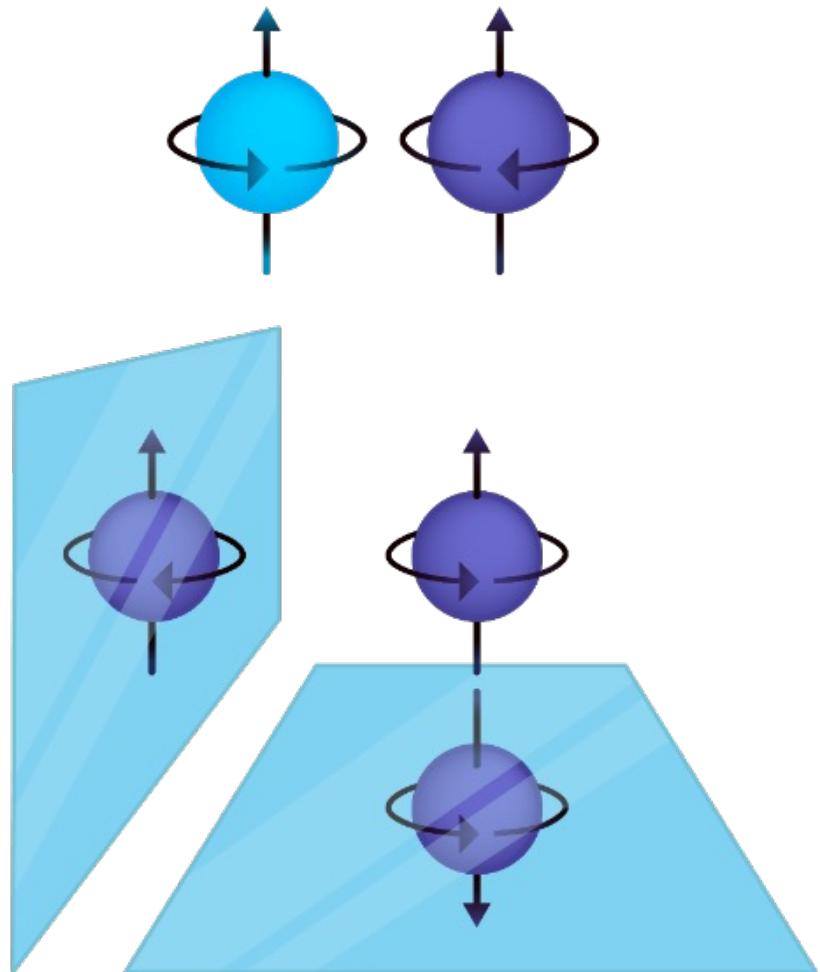
linkshändiges Fermion
 $\chi = -1$



rechtshändiges Fermion
 $\chi = +1$



chirality





Nobel Symposium 167 on Chiral matter

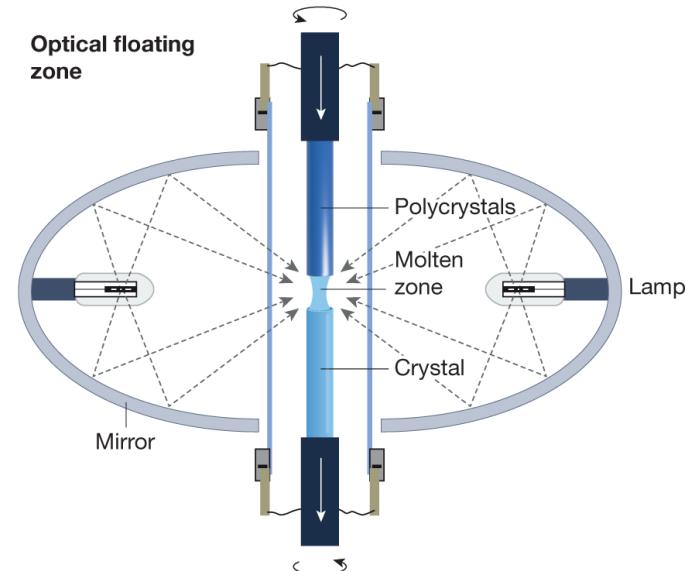
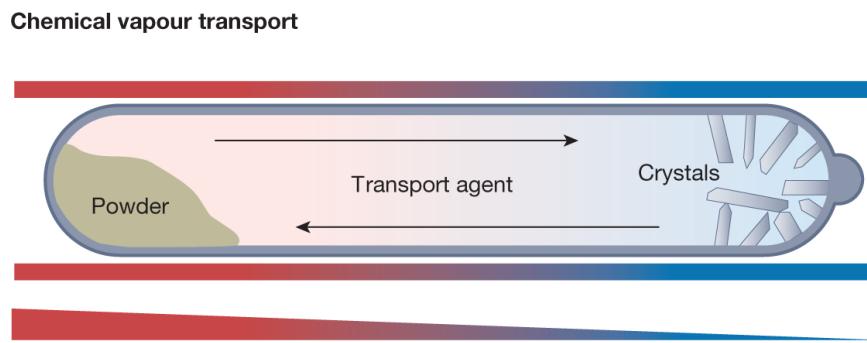
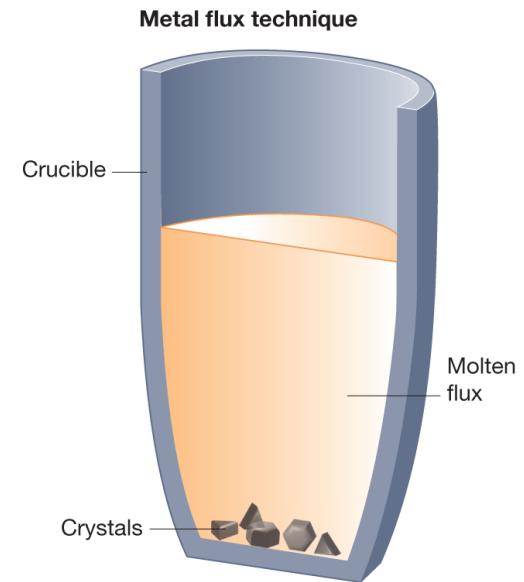
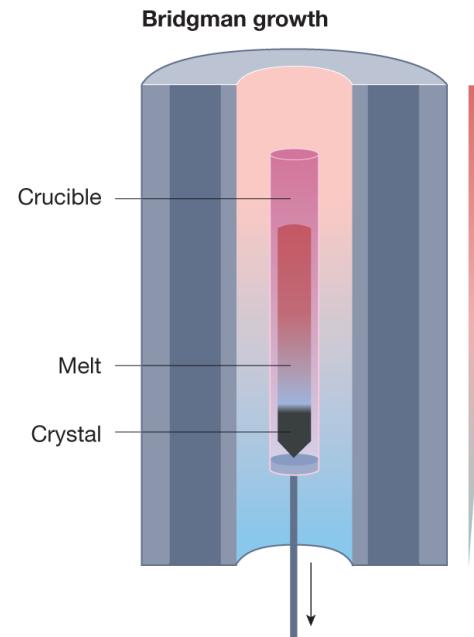
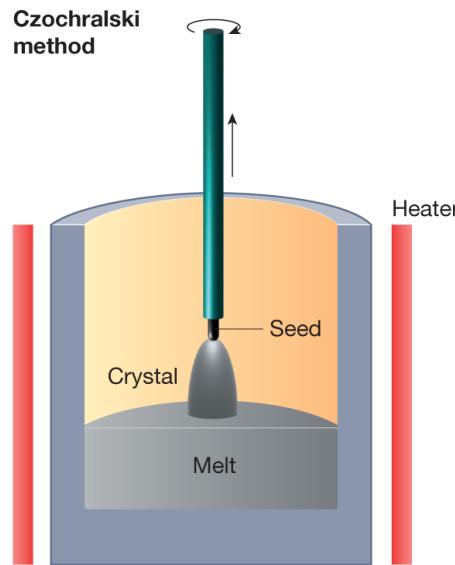
June 28 – July 1, 2021

Lidingö, Sweden

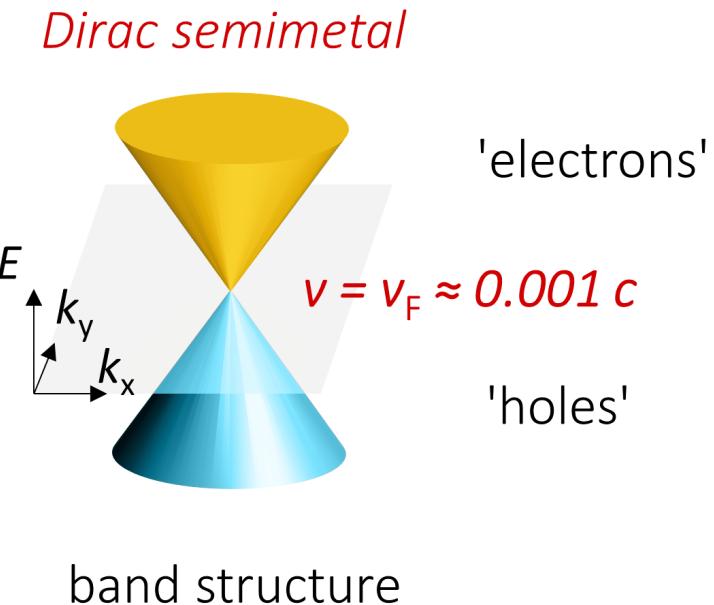
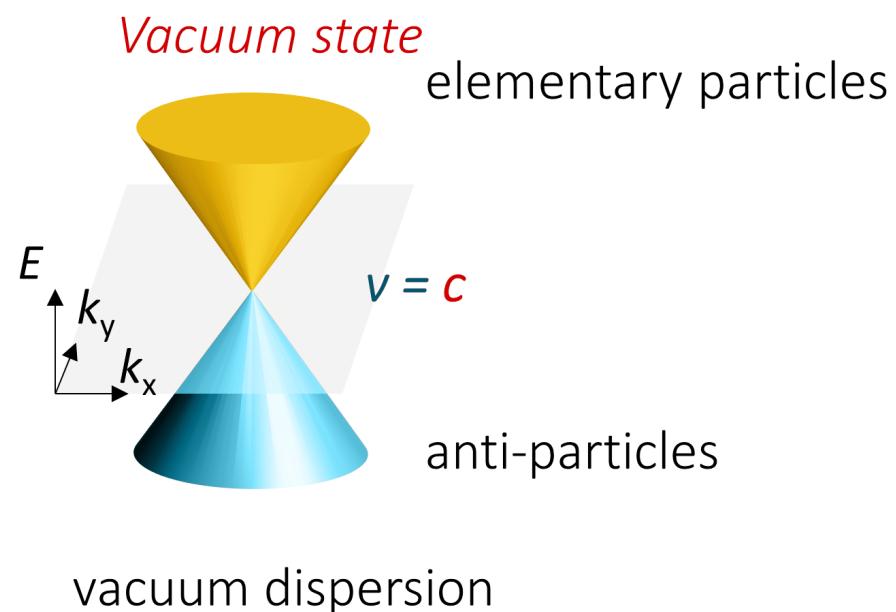
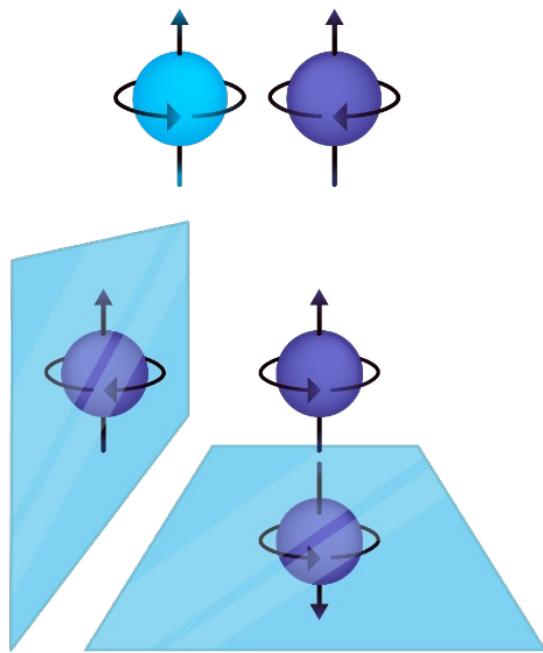


NOBEL SYMPOSIA

synthesis



chiral anomaly



Chiral anomaly is the **anomalous non-conservation** of a chiral current.

A sealed box with equal numbers of positive and negative charged particles is found when it is opened to have more positive than negative particles, or vice-versa.

Prohibited from classical conservation laws, but can be **broken in a quantum world**.

Universe contains more matter than antimatter

Wikipedia

Adler, Phys. Rev. **177**, 2426 (1969)

Bell and. Jackiw, Nuovo Cim. **A60**, 47 (1969)

Zyuzin, Burkov - Physical Review B (2012)

Burkov, Balents, PRL **107**, 12720 (2012)

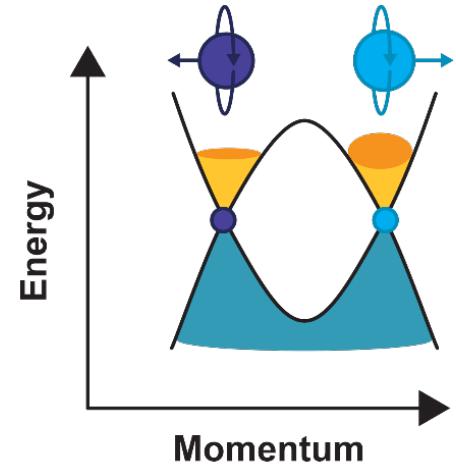
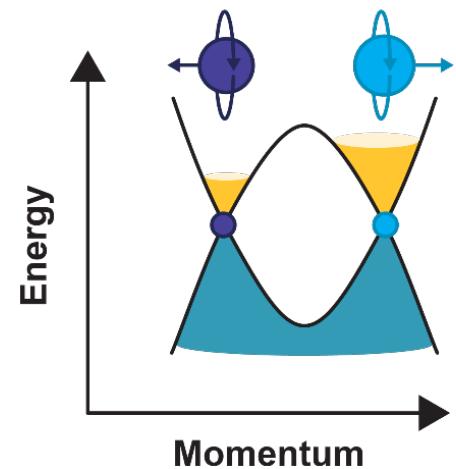
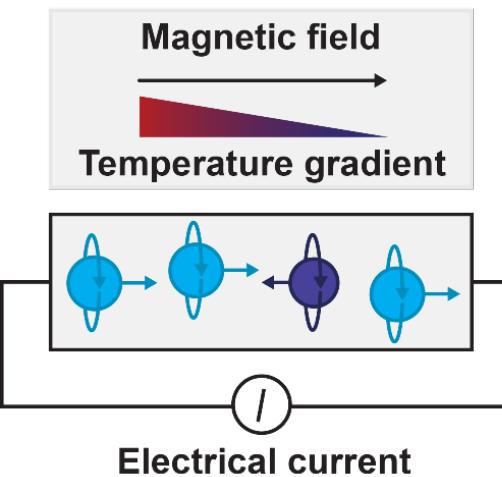
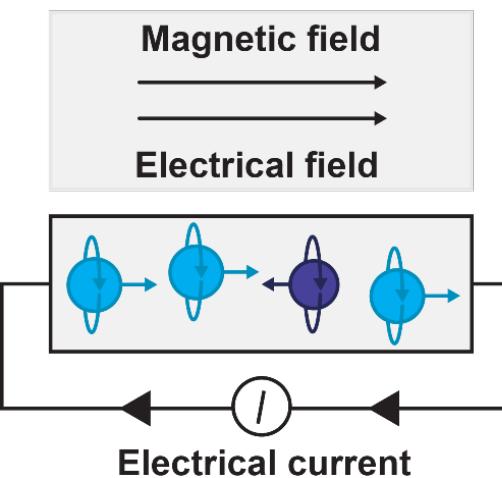
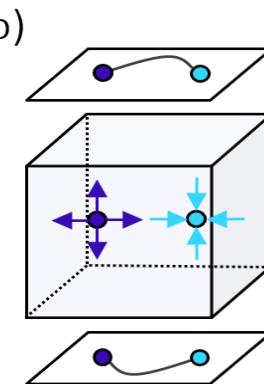
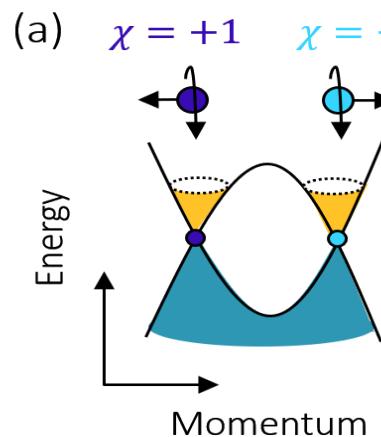
Burkov, J. Phys.: Condens. Matter **27**, 113201 (2015)

Volovik, The Universe in a Helium Droplet (International Series of Monographs on Physics, Band 117) ISBN: 9780199564842

Weyl semimetals

3D topological Weyl - transport measurements:

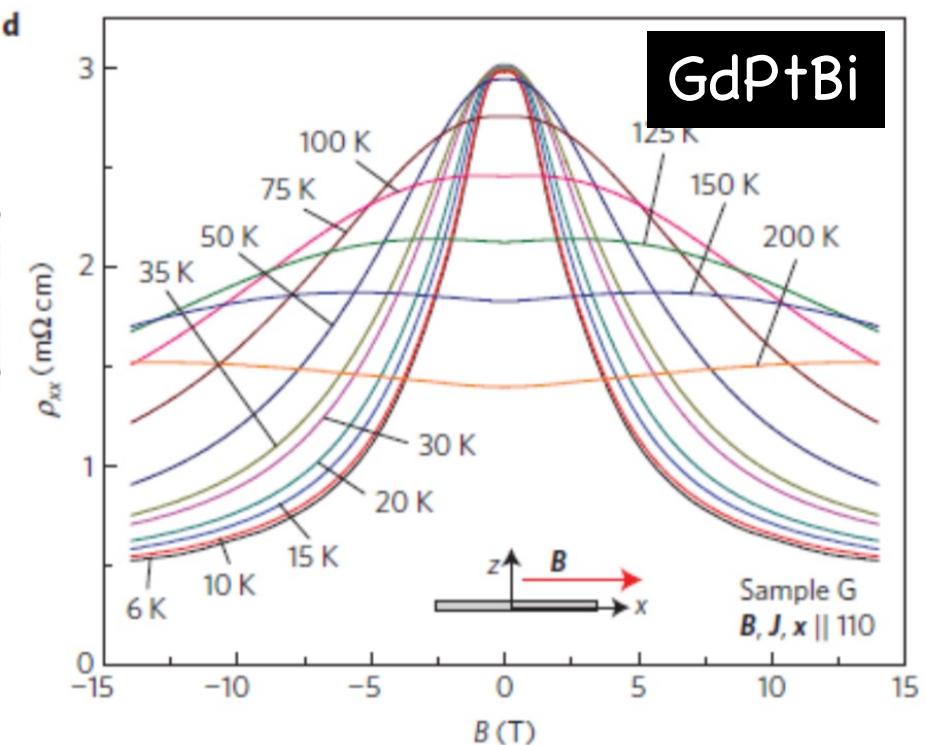
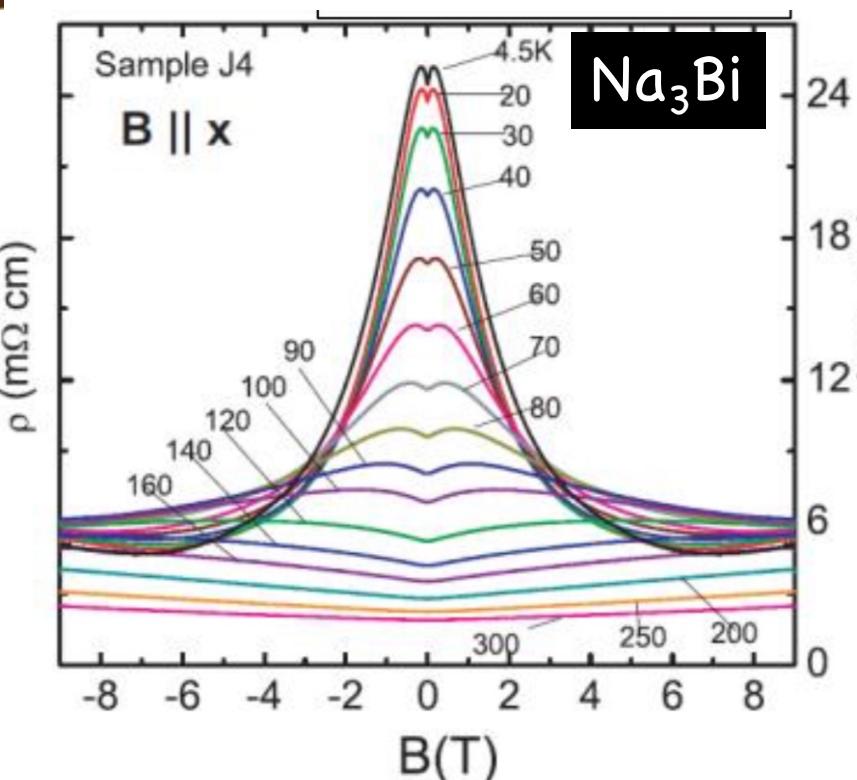
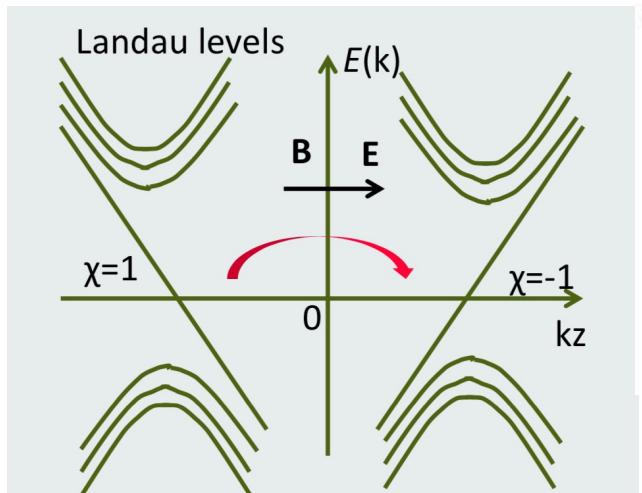
1. Giant responses to an external stimuli
2. Fermi arc
3. Chiral anomaly
4. Axial gravitational anomaly



chiral anomaly



$$A = -\left(\frac{L^2}{2\pi\ell_B^2}\right)\left(\frac{Le\dot{k}_z}{2\pi}\right) = -V \frac{e^3}{4\pi^2\hbar^2} \mathbf{E} \cdot \mathbf{B}$$



Jun Xiong, S. Kushwaha et al., *Science* 2015

Hirschberger et al., *Nature Materials*, 2016

S. L. Adler, Phys. Rev. 177, 2426 (1969)

J. S. Bell and R. Jackiw, Nuovo Cim. A60, 47 (1969)

AA Zyuzin, AA Burkov - Physical Review B (2012)



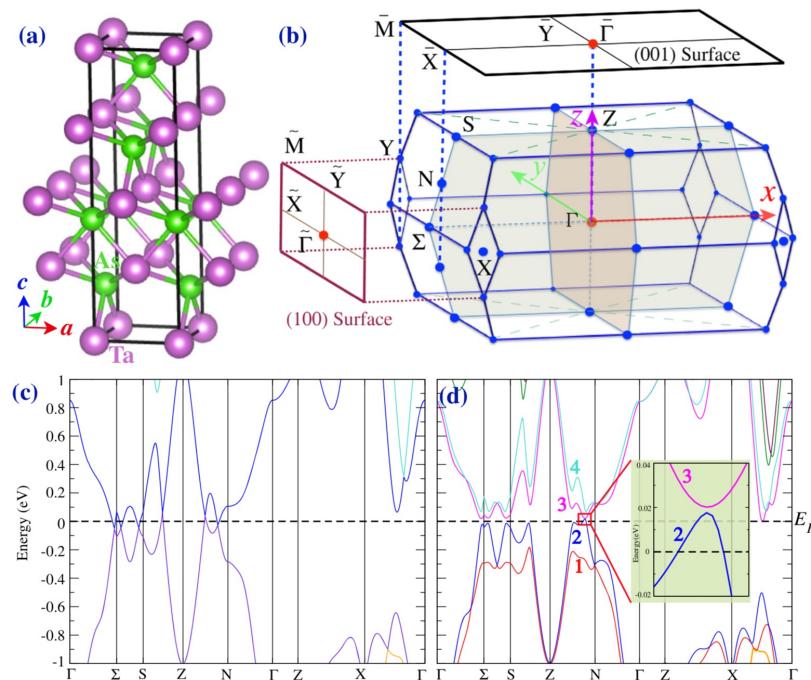
Weyl semimetals

NbP, NbAs, TaP, TaAs



Weyl Semimetal Phase in Noncentrosymmetric Transition-Metal Monophosphides

Hongming Weng, Chen Fang, Zhong Fang, B. Andrei Bernevig, and Xi Dai
Phys. Rev. X 5, 011029 – Published 17 March 2015



ARTICLE

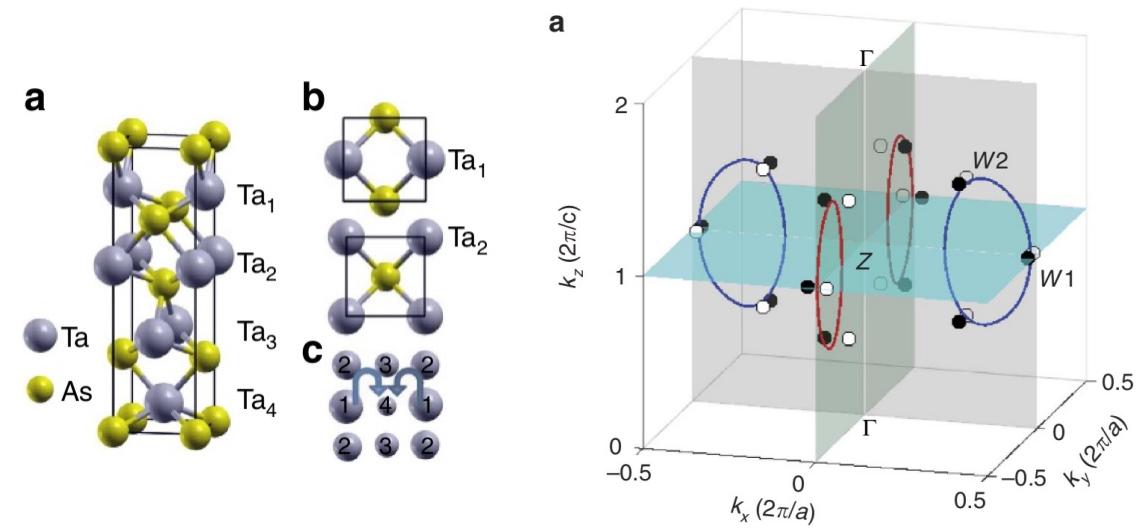
Received 24 Nov 2014 | Accepted 30 Apr 2015 | Published 12 Jun 2015

DOI: 10.1038/ncomms8373

OPEN

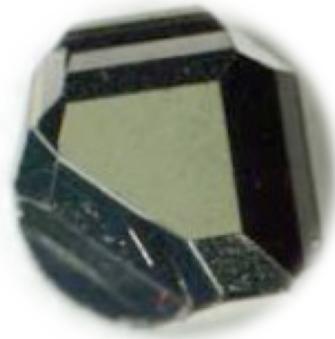
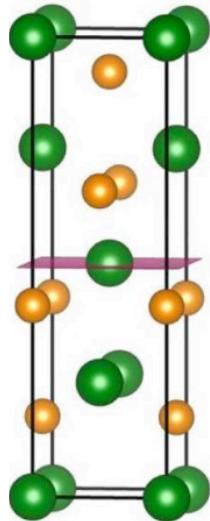
A Weyl Fermion semimetal with surface Fermi arcs in the transition metal monopnictide TaAs class

Shin-Ming Huang^{1,2,*}, Su-Yang Xu^{3,4,*}, Ilya Belopolski^{3,4,*}, Chi-Cheng Lee^{1,2}, Guoqing Chang^{1,2}, BaoKai Wang^{1,2,5}, Nasser Alidoust^{3,4}, Guang Bian³, Madhab Neupane^{3,4,6}, Chenglong Zhang⁷, Shuang Jia^{7,8}, Arun Bansil⁵, Hsin Lin^{1,2} & M. Zahid Hasan^{3,4,9}

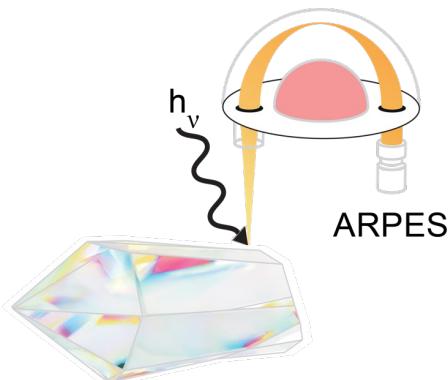


Weyl semimetals

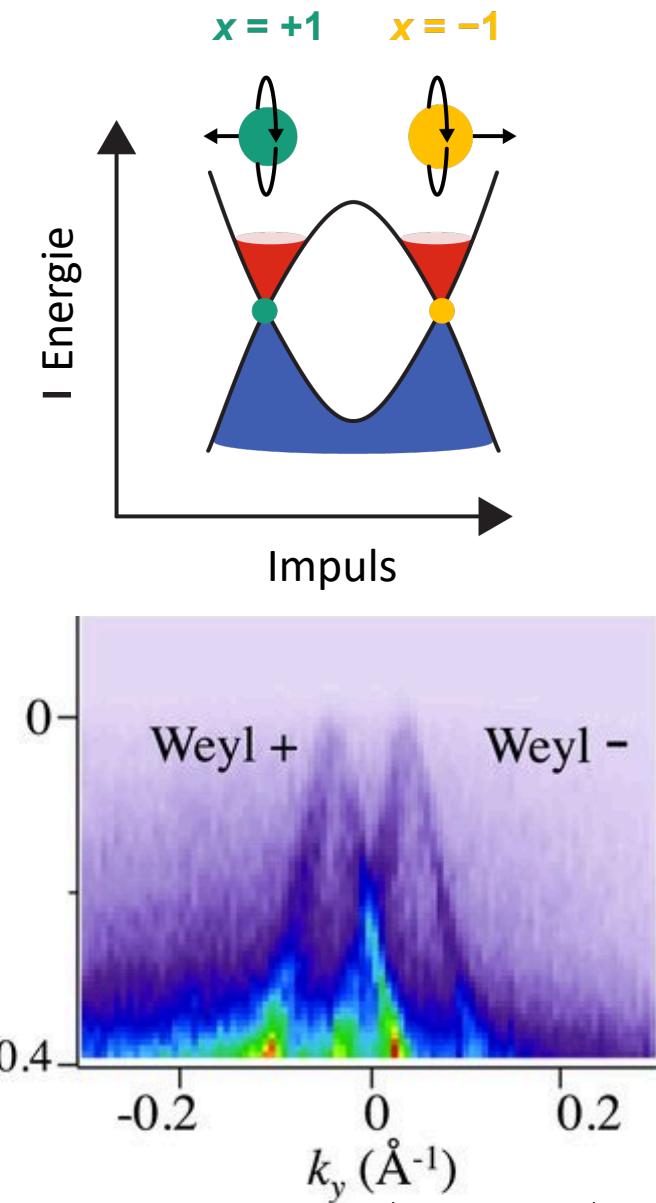
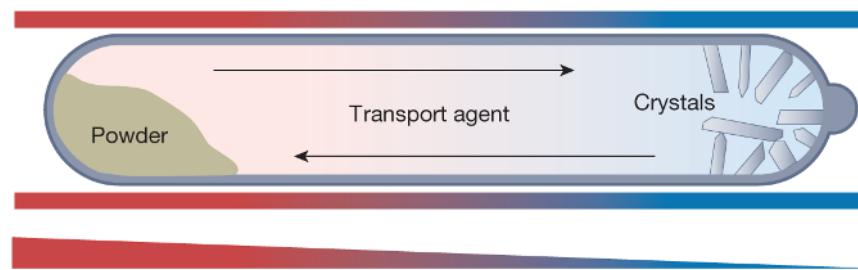
NbP, NbAs, TaP, TaAs



angle resolved Photoemission

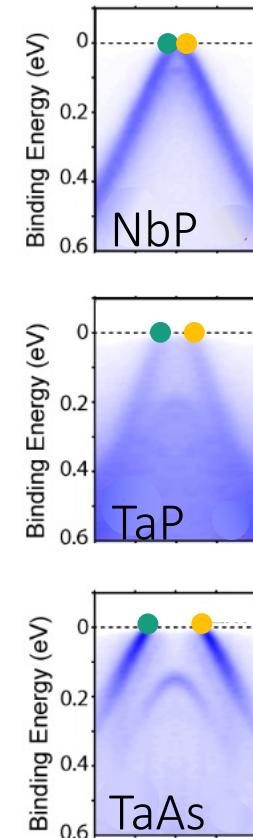
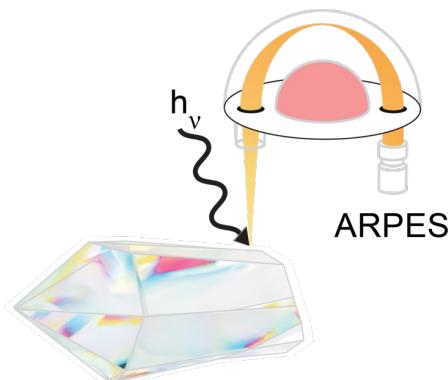
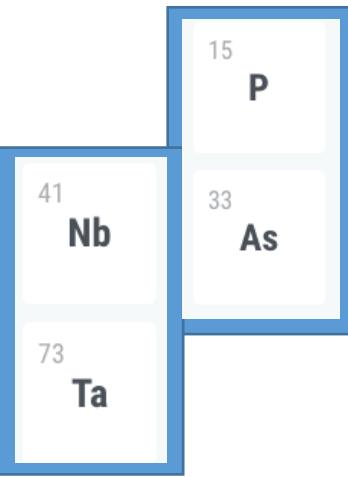
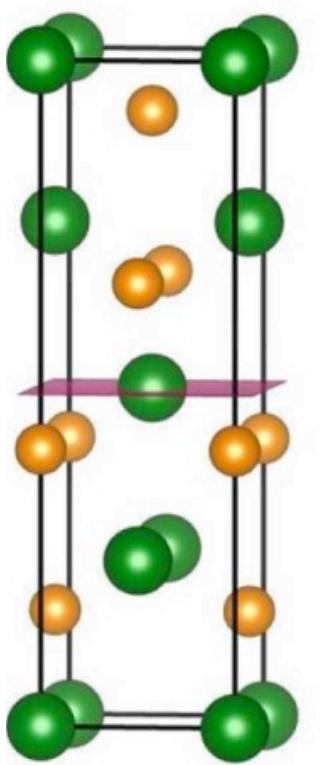


chemical vapor transport

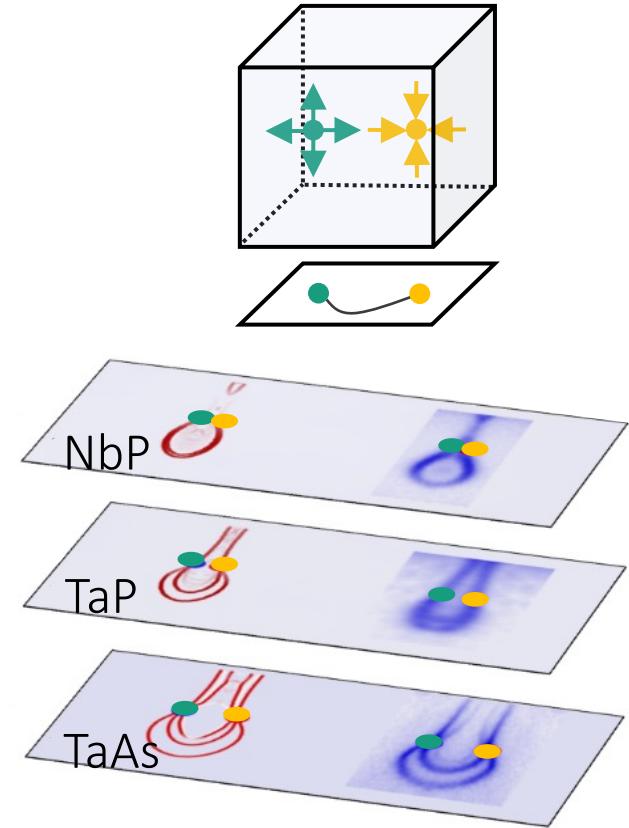


Weyl semimetals

NbP, NbAs, TaP, TaAs

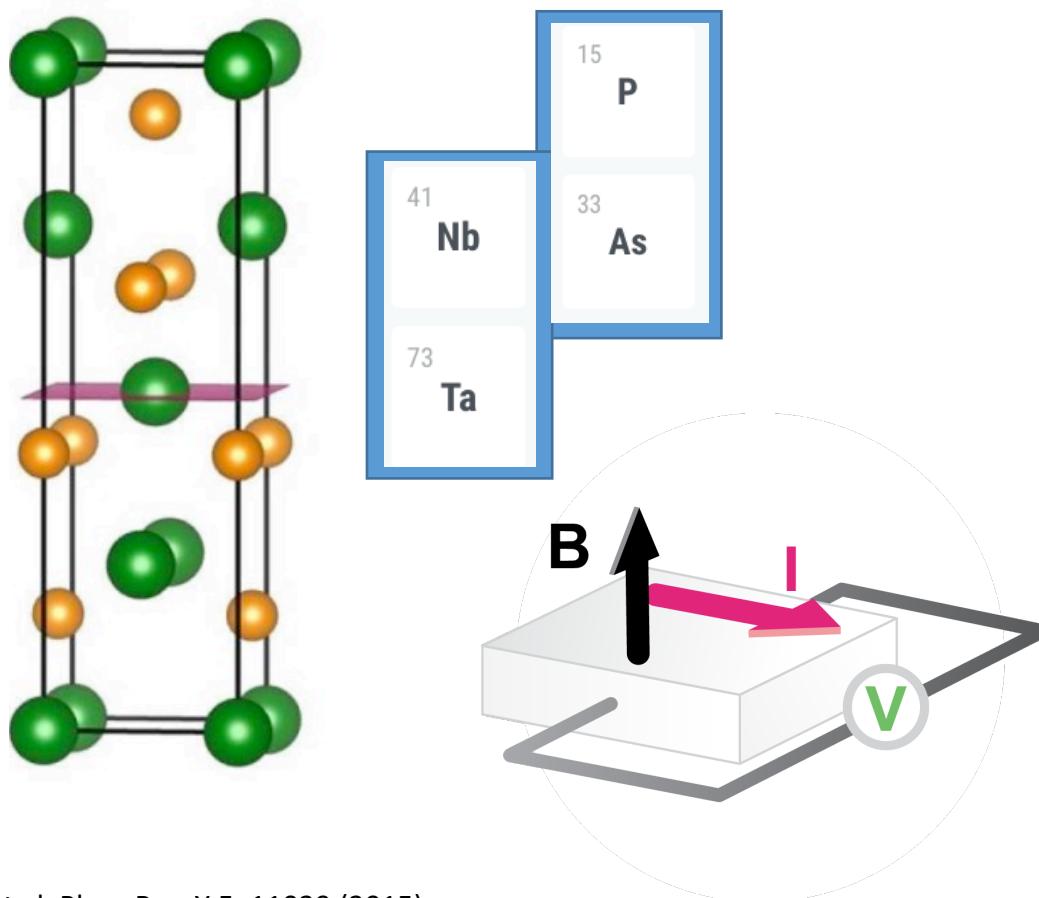


$$\chi = +1 \quad \chi = -1$$



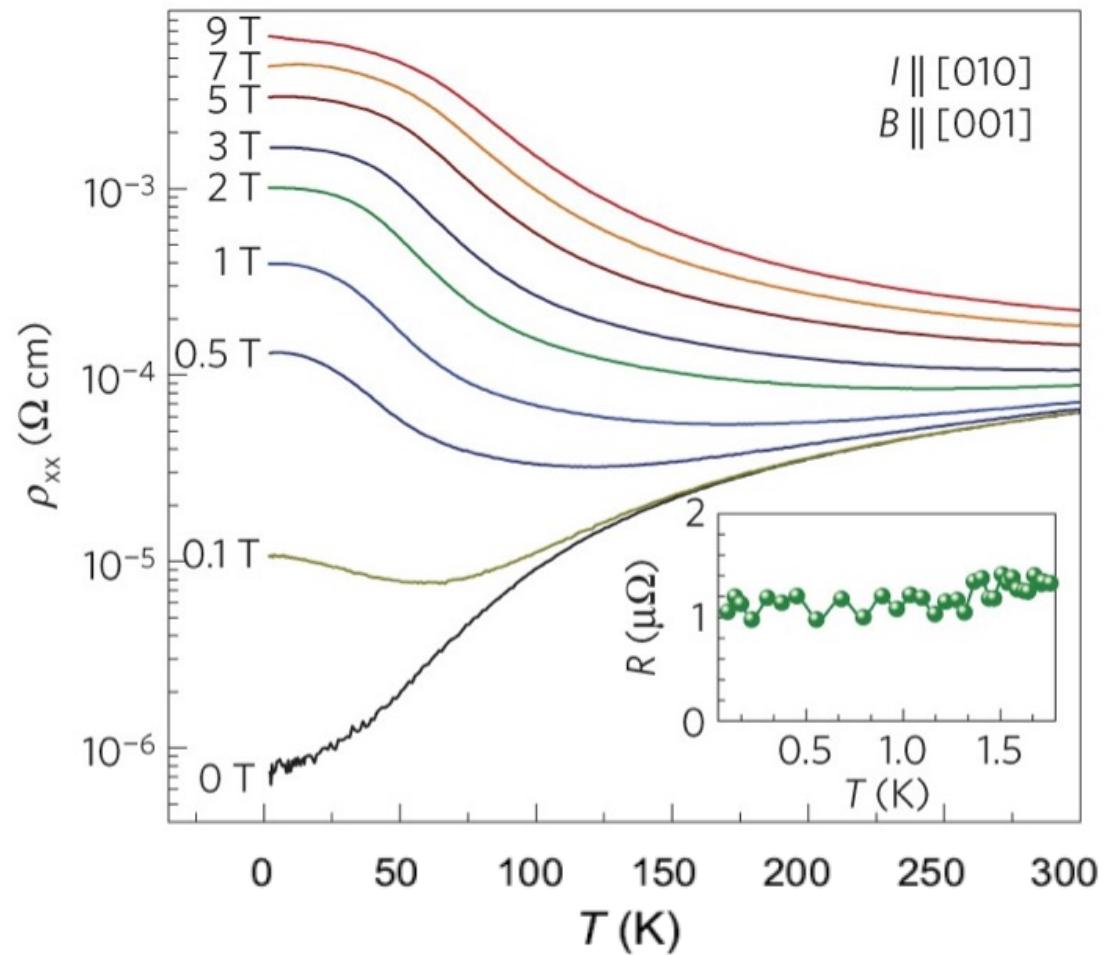
Weyl semimetals

NbP, NbAs, TaP, TaAs



Weng, et al. Phys. Rev. X 5, 11029 (2015)
Huang et al. preprint arXiv:1501.00755

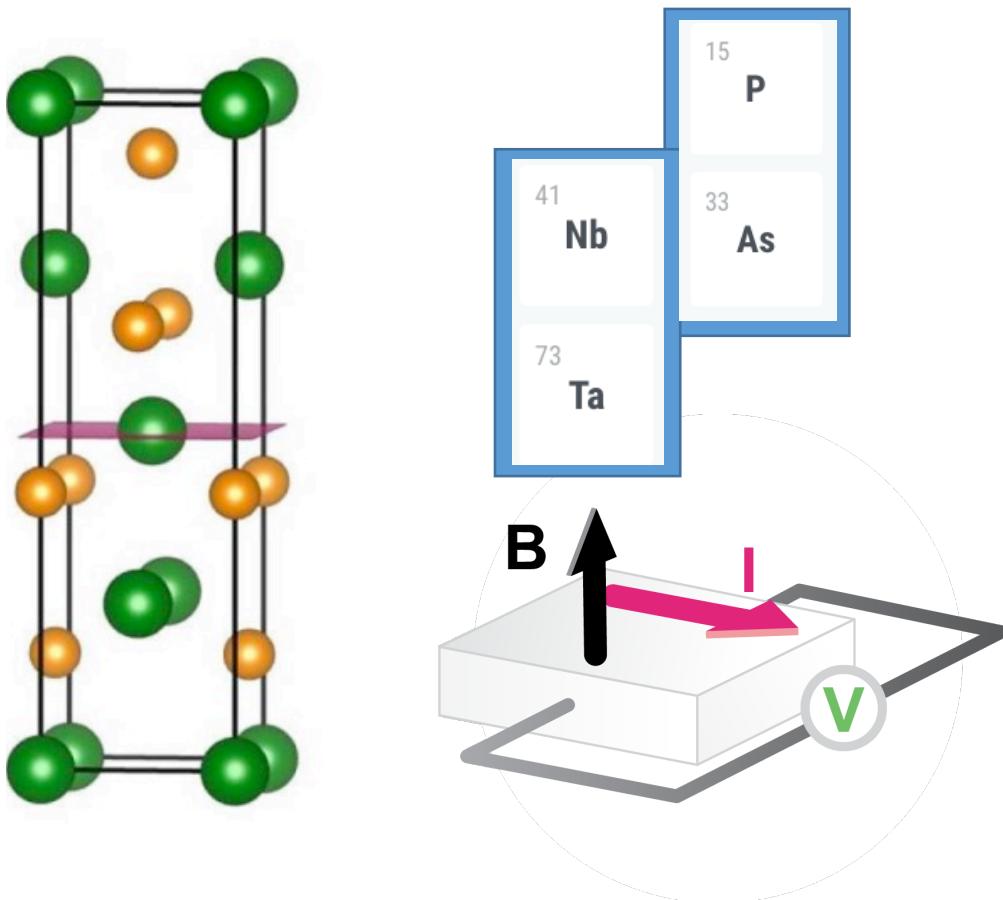
giant magnetoresistance



Shekhar, et al., Nature Physics 11 (2015) 645
Frank Arnold, et al. Nature Communication 7 (2016) 11615

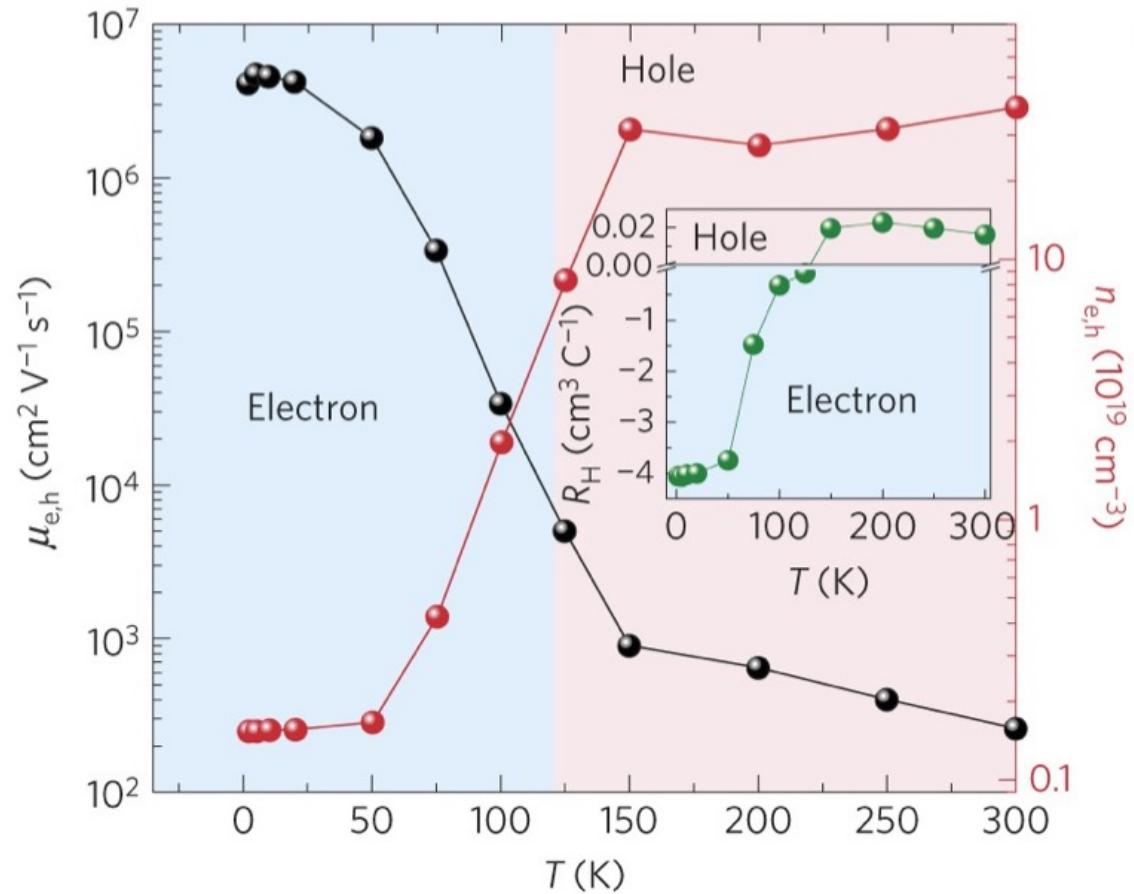
Weyl semimetals

NbP, NbAs, TaP, TaAs



Weng, et al. Phys. Rev. X 5, 11029 (2015)
Huang et al. preprint arXiv:1501.00755

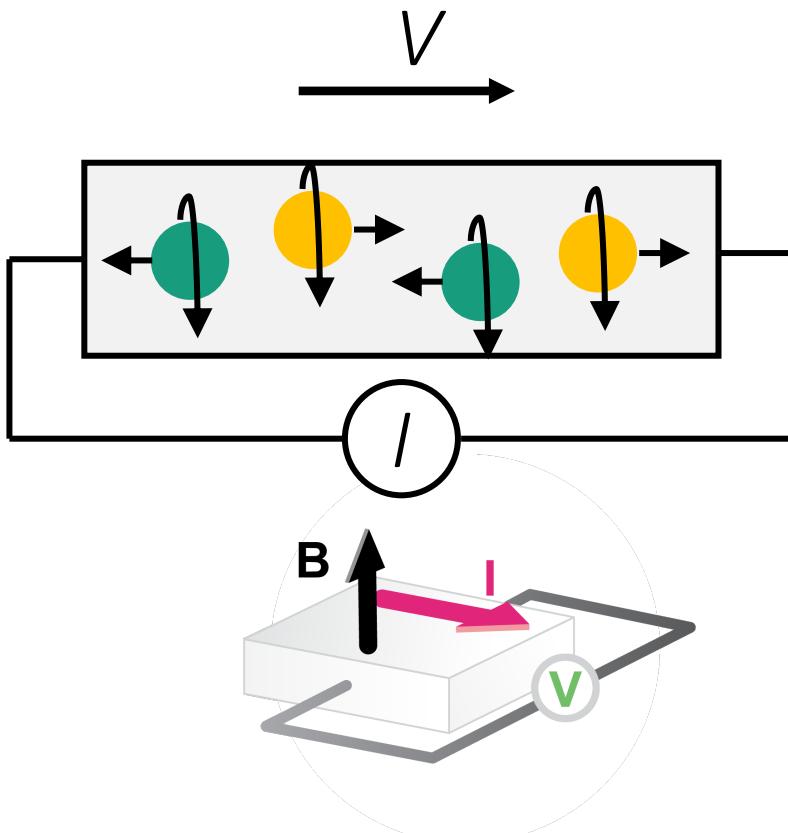
giant mobility



Shekhar, et al., Nature Physics 11 (2015) 645
Frank Arnold, et al. Nature Communication 7 (2016) 11615

chiral anomaly

conservation of chirality



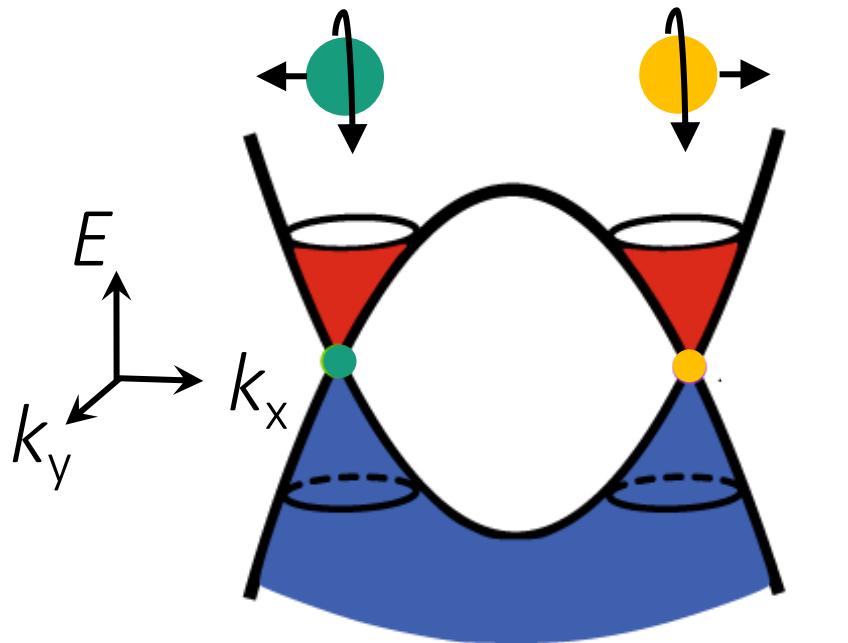
J. Gooth et al., Nature 547 (2017) 324 arXiv:1703.10682.

J. Gooth, J. Kübler, C. Felser, Physik Journal 20 (2021) 29.

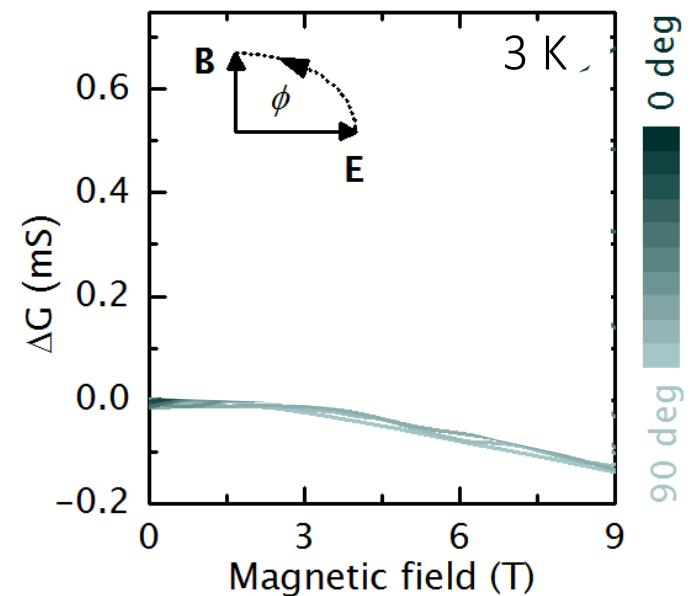
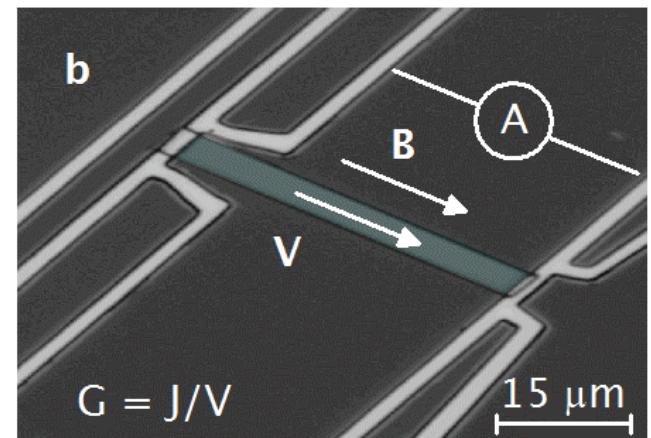
Lucas, A., Davison, R. A. & Sachdev, S. PNAS 113, 9463–9468 (2016).

Landsteiner K., et al. Gravitational anomaly and transport phenomena. Phys. Rev. Lett. 107, 021601 (2011). URL

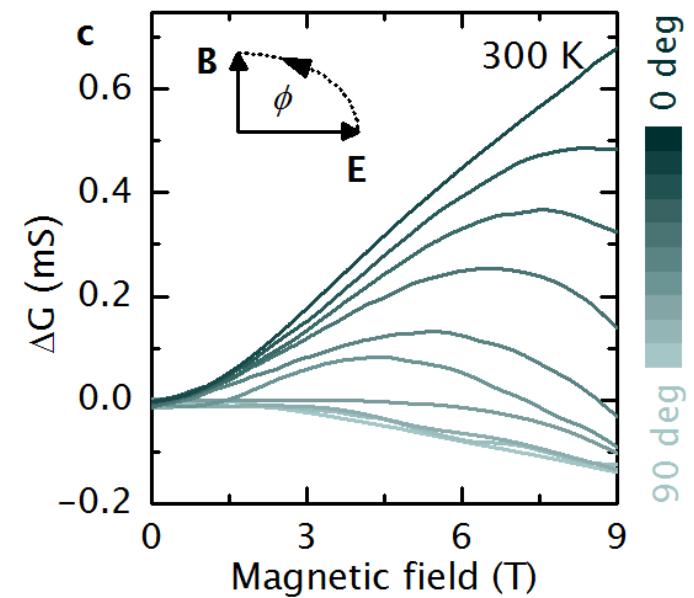
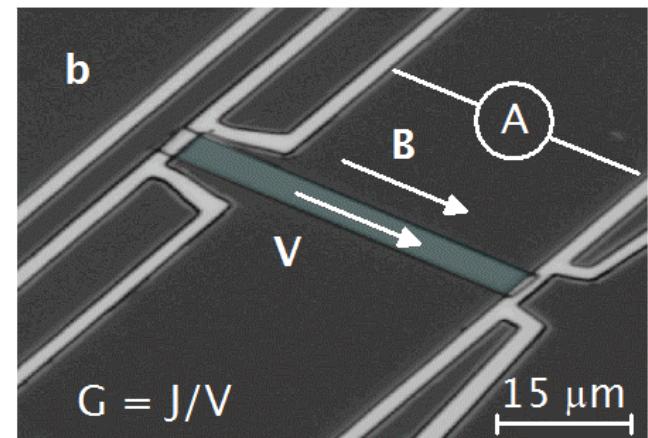
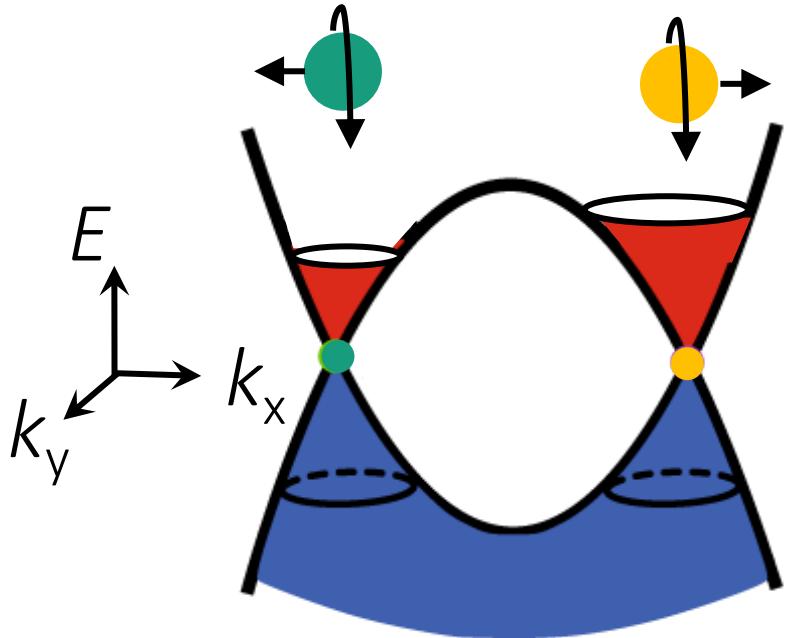
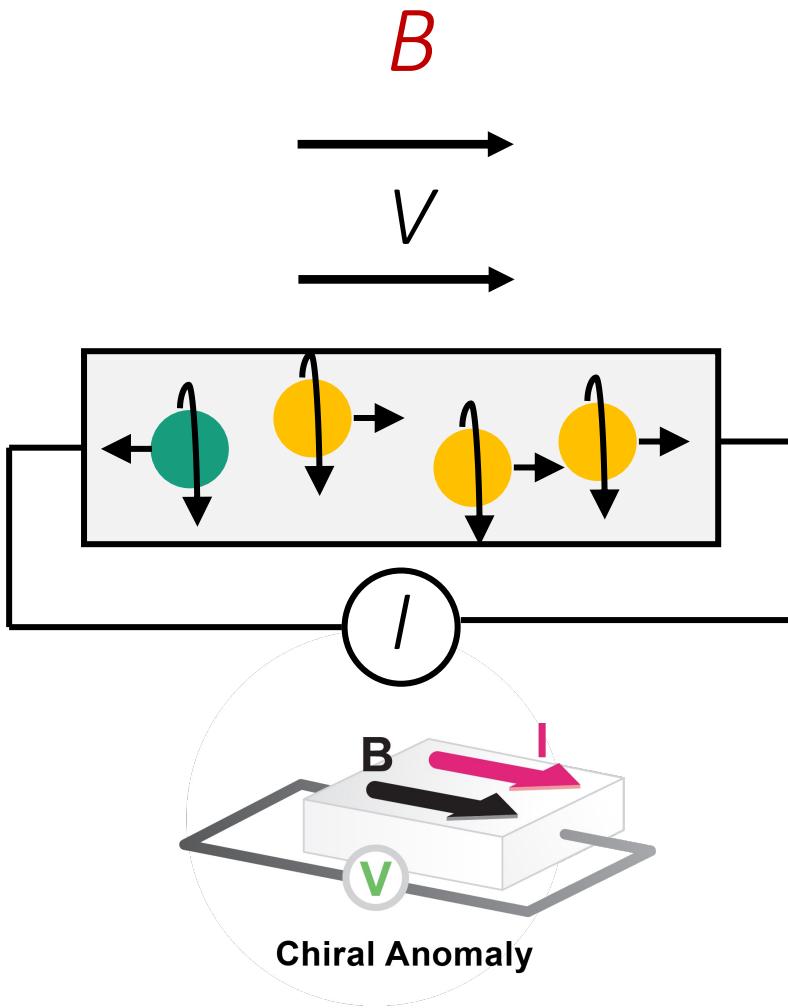
Jensen, et al. Thermodynamics, gravitational anomalies and cones. Journal of High Energy Physics 2013, 88 (2013).



$$\chi = +1 \quad \chi = -1$$



chirale Anomalie



J. Gooth et al., Nature 547 (2017) 324 arXiv:1703.10682.

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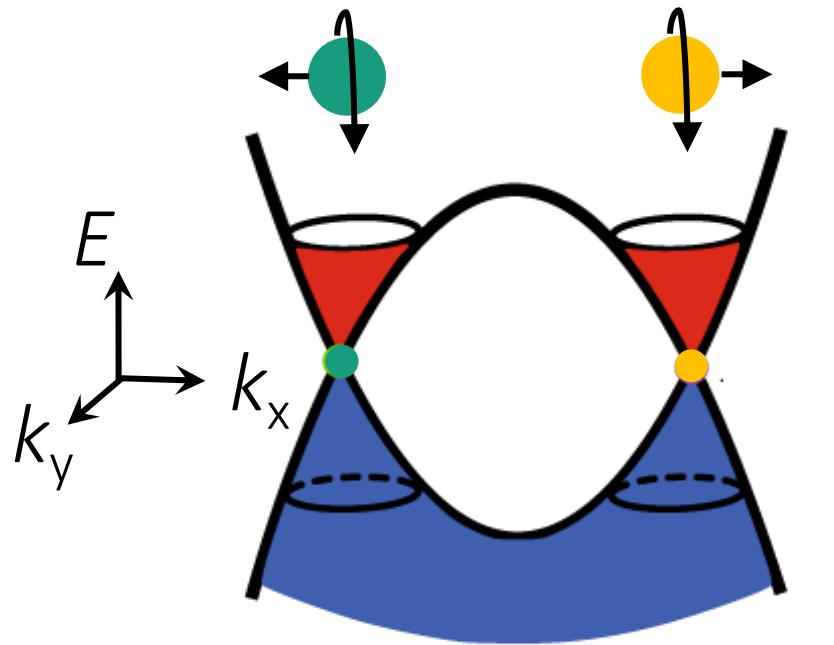
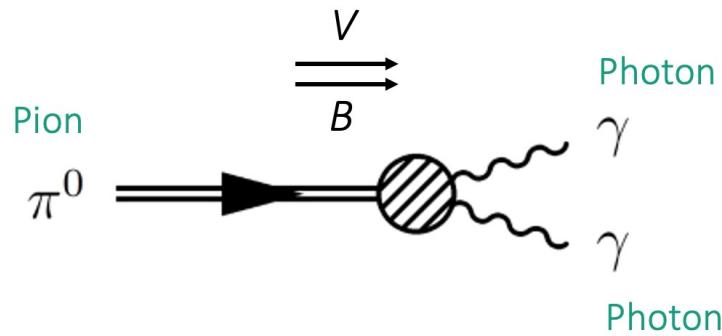
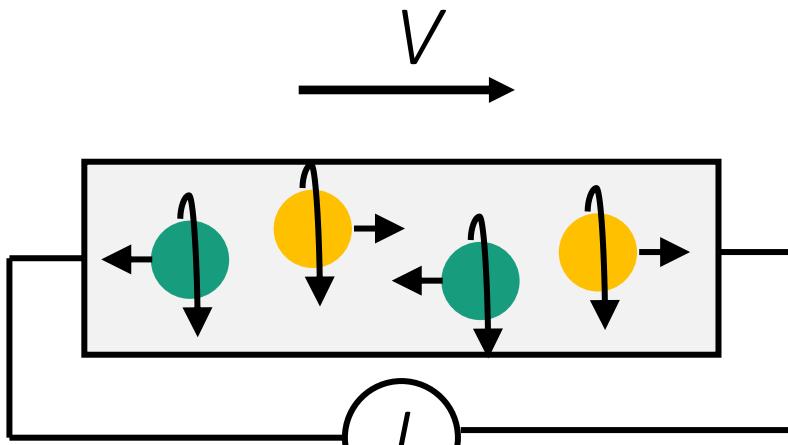
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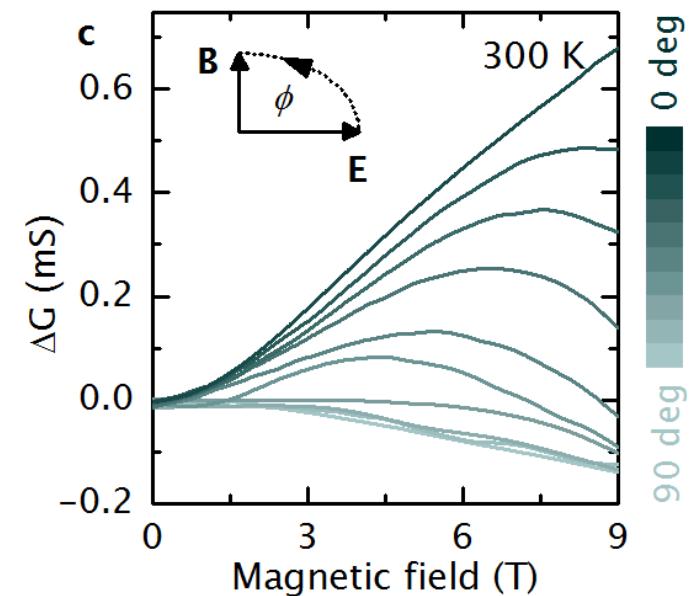
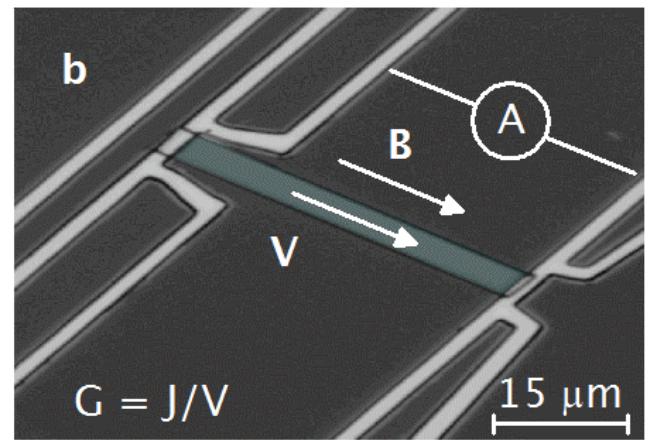
Jensen, et al. Thermodynamics, gravitational anomalies and cones. Journal of High Energy Physics 2013, 88 (2013).

chiral anomaly

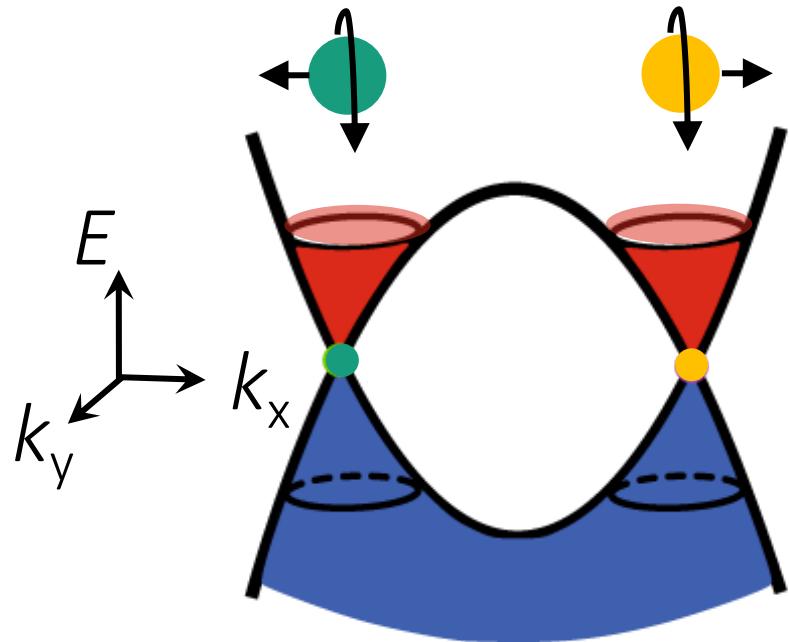
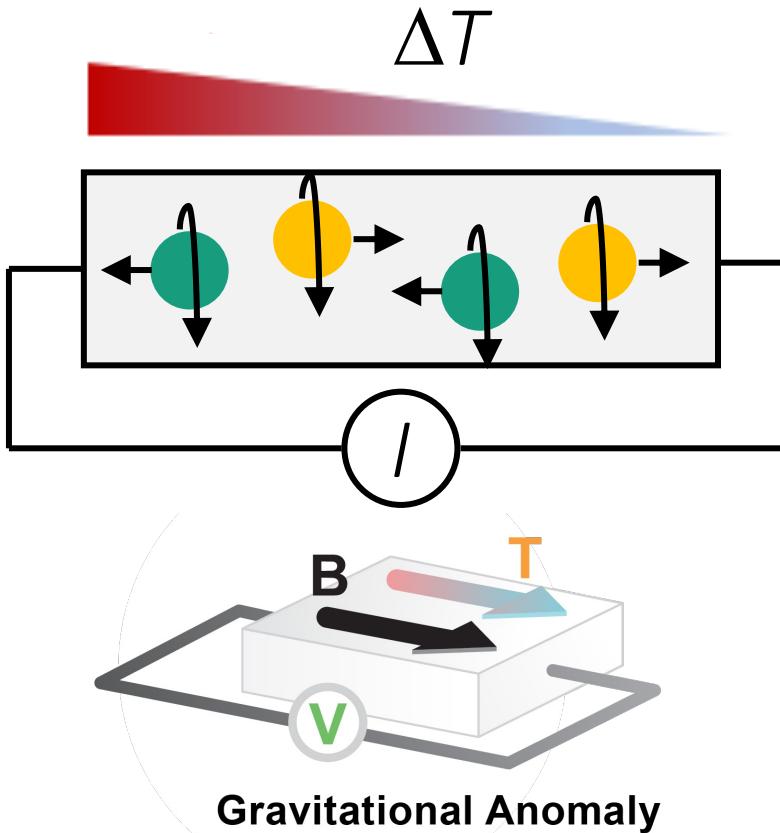
Erhaltung der Chiralität



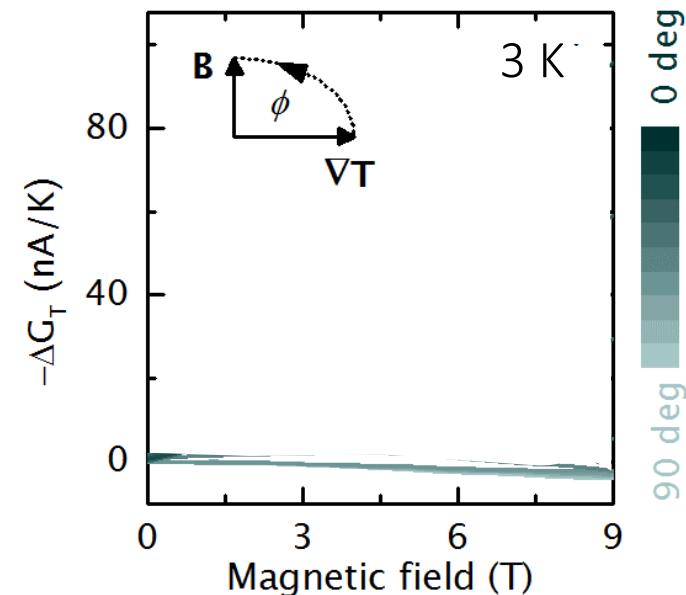
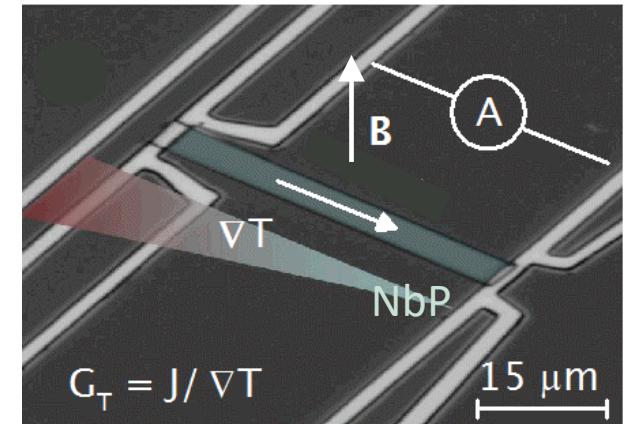
$$\chi = +1 \quad \chi = -1$$



axial gravitational anomaly



$$\chi = +1 \quad \chi = -1$$



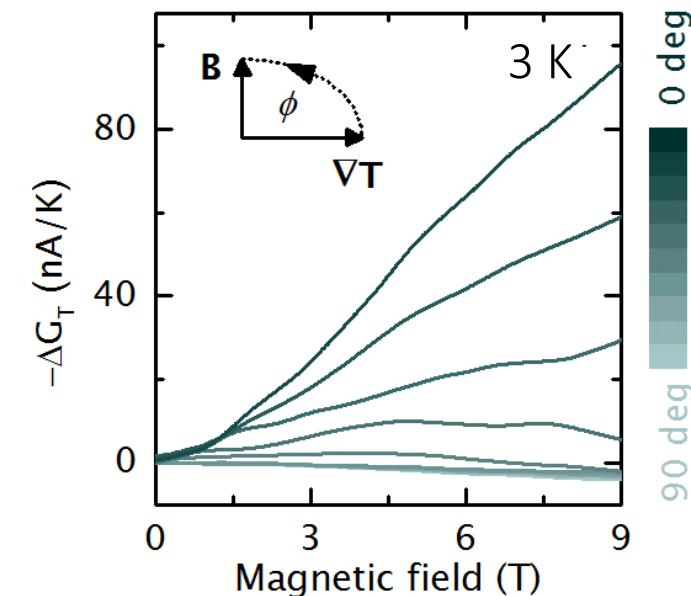
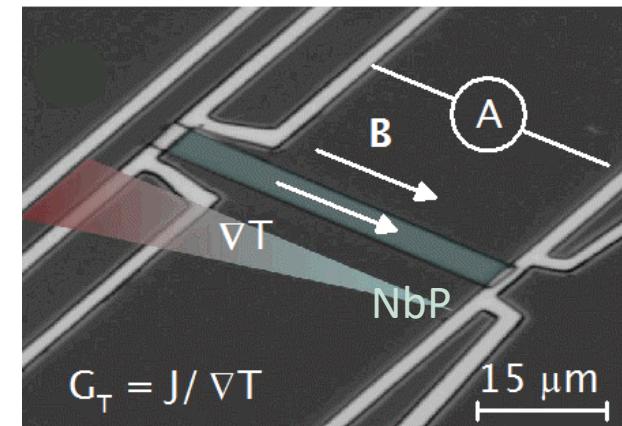
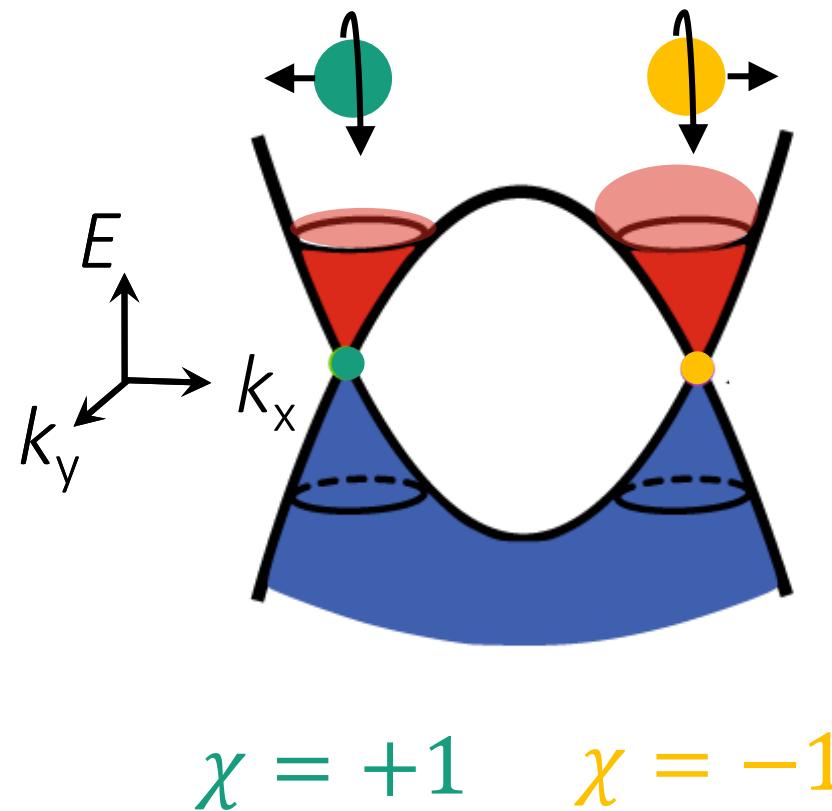
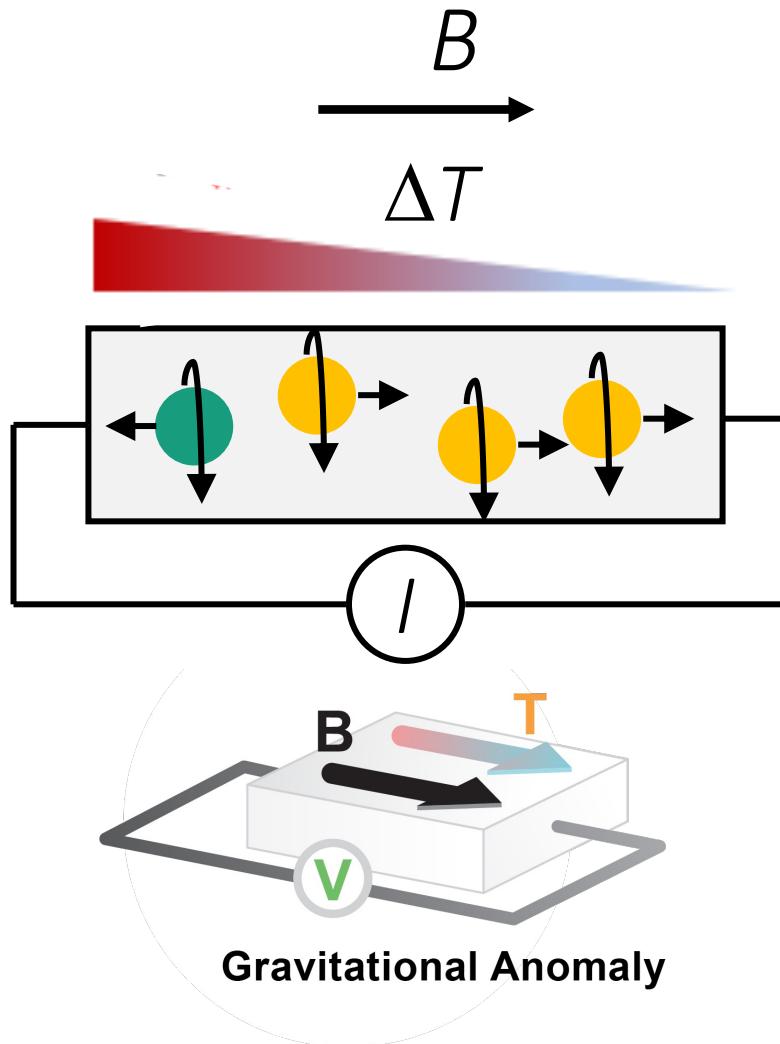
Johannes Gooth et al., Nature 547 (2017) 324 arXiv:1703.10682

Lucas, A., Davison, R. A. & Sachdev, S. PNAS 113, 9463–9468 (2016).

Landsteiner K., et al. Gravitational anomaly and transport phenomena. Phys. Rev. Lett. 107, 021601 (2011). URL

Jensen, et al. Thermodynamics, gravitational anomalies and cones. Journal of High Energy Physics 2013, 88 (2013).

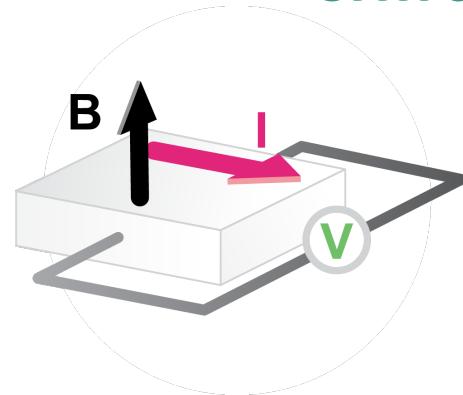
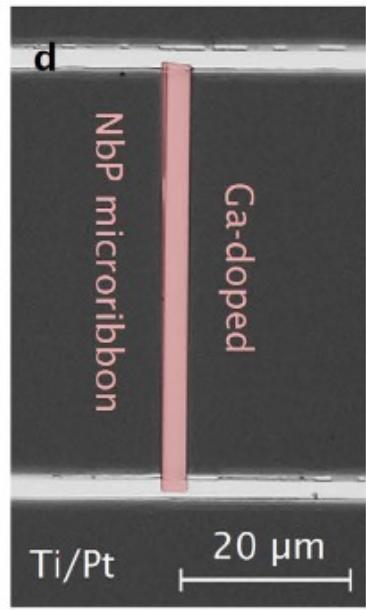
axial gravitational anomaly



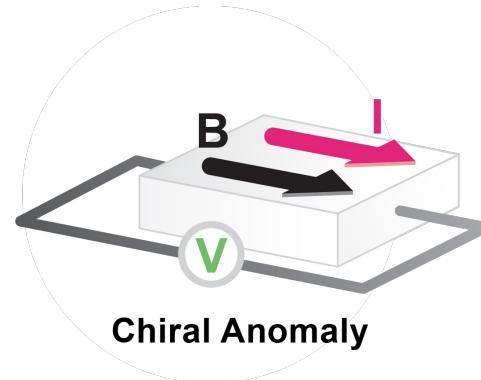
J. Gooth et al., Nature 547 (2017) 324 arXiv:1703.10682.

J. Gooth, J. Kübler, C. Felser, Physik Journal 20 (2021) 29.

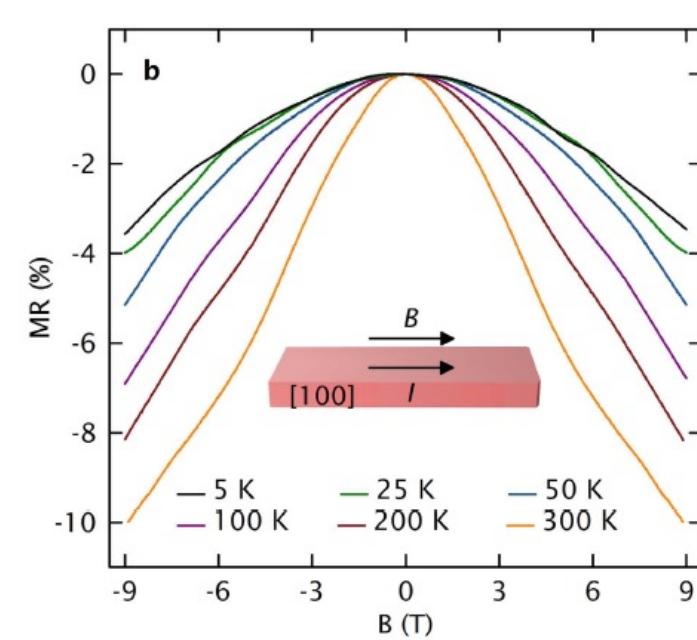
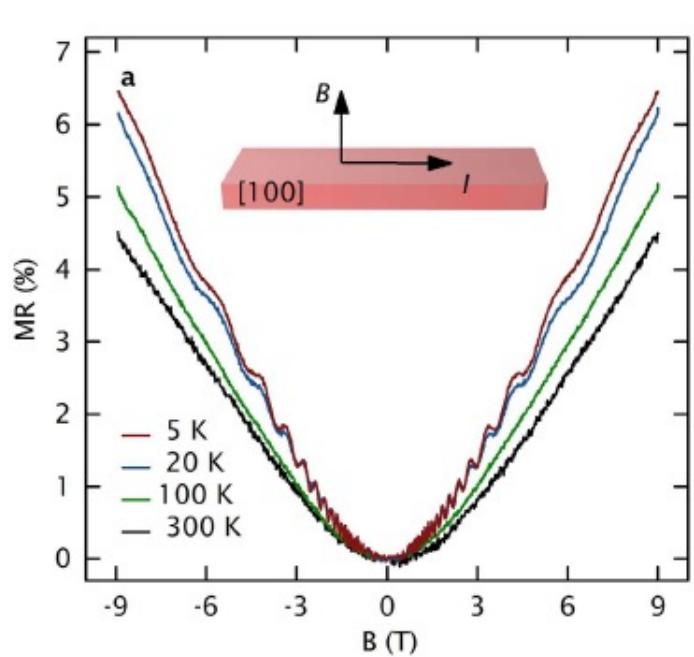
chiral anomaly



NbP



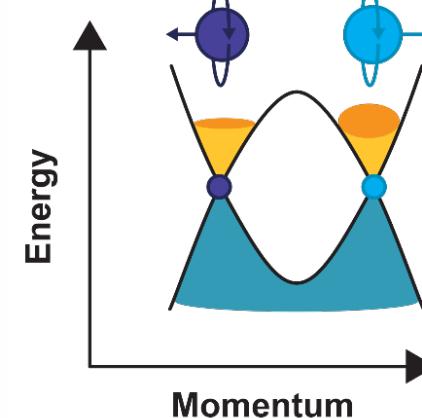
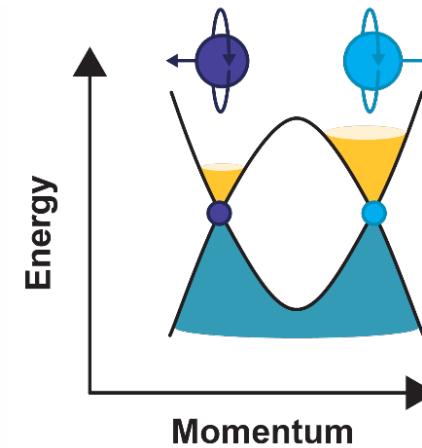
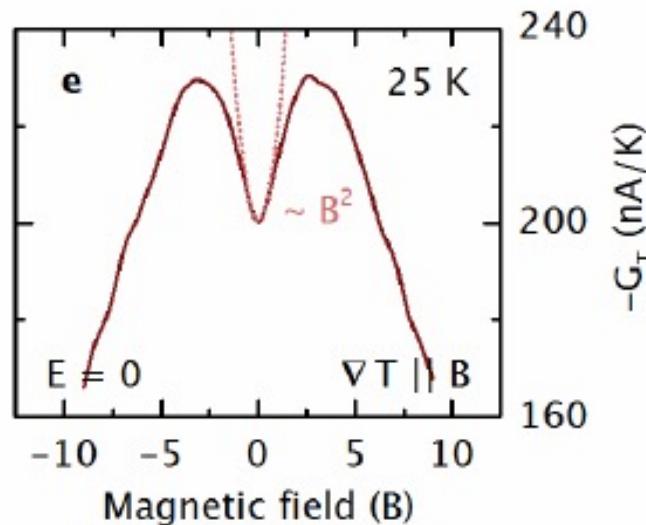
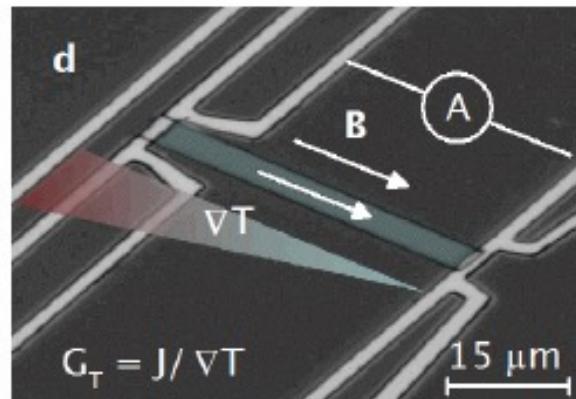
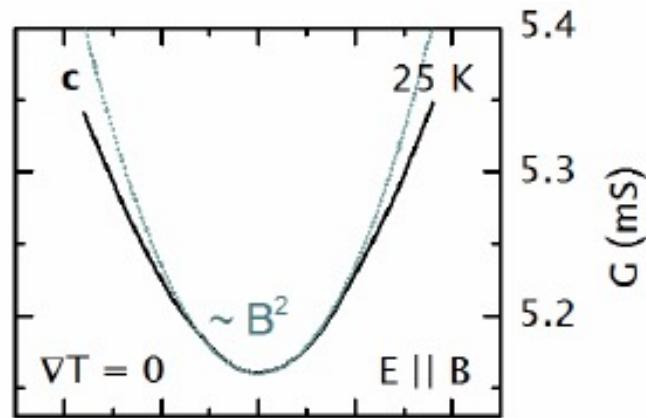
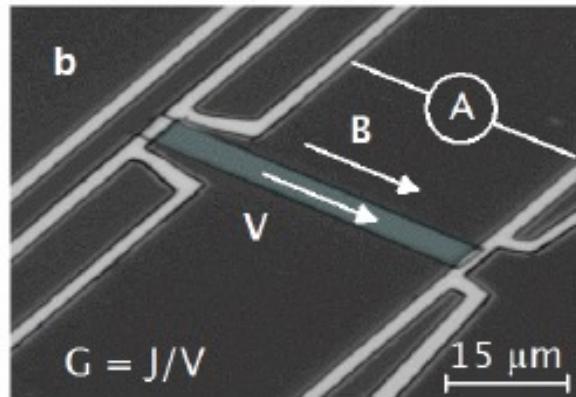
Chiral Anomaly



Ga-doping relocates the Fermi energy in NbP close to the W2 Weyl points
Therefore, we observe a negative MR as a signature of the chiral anomaly, that survives up to room temperature

NbP

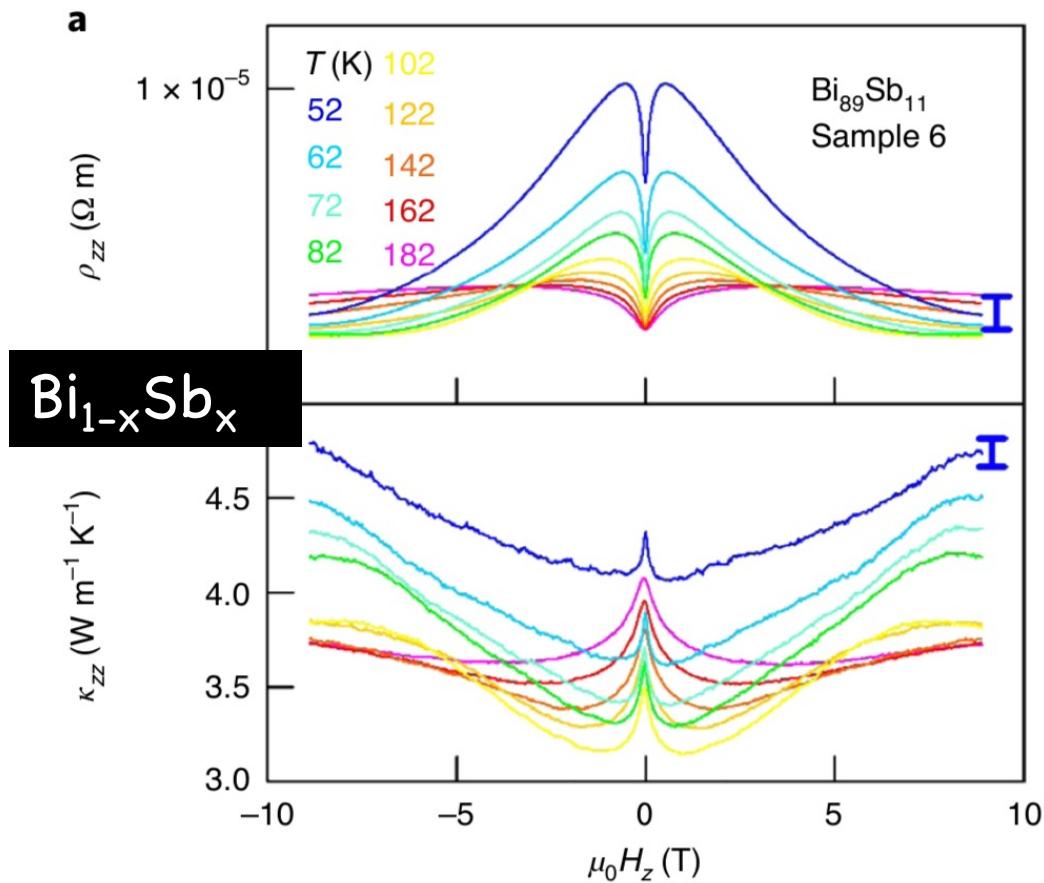
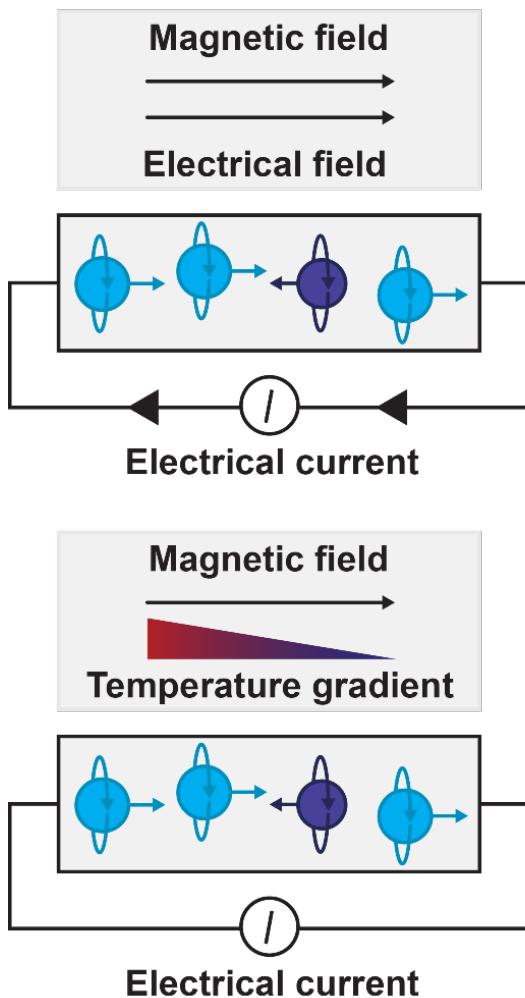
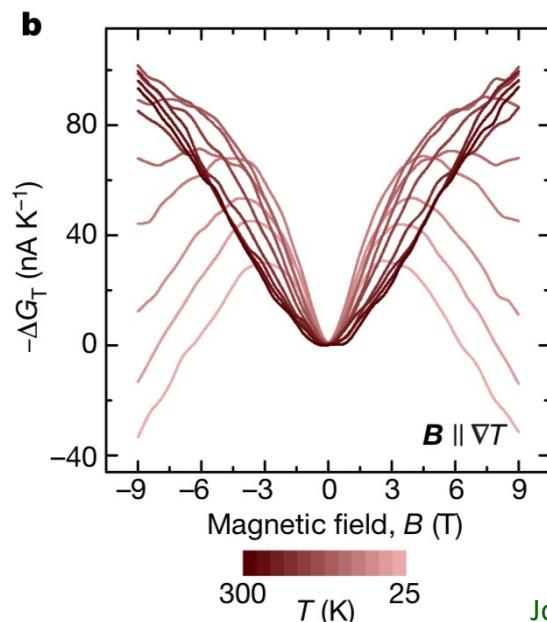
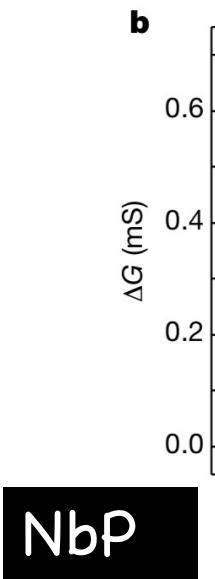
chiral and axial gravitational anomaly



A positive longitudinal magneto-thermoelectric conductance (PMTC) in the Weyl semimetal NbP for collinear temperature gradients and magnetic fields that vanishes in the ultra quantum limit.

Lucas, A., Davison, R. A. & Sachdev, S. PNAS **113**, 9463–9468 (2016).

chiral and axial gravitational anomaly



the universe in a crystal



The New York Times

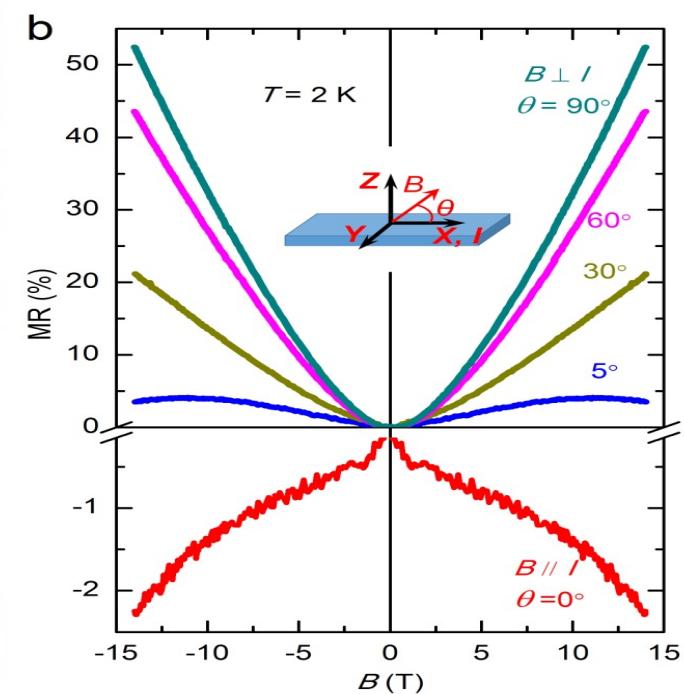
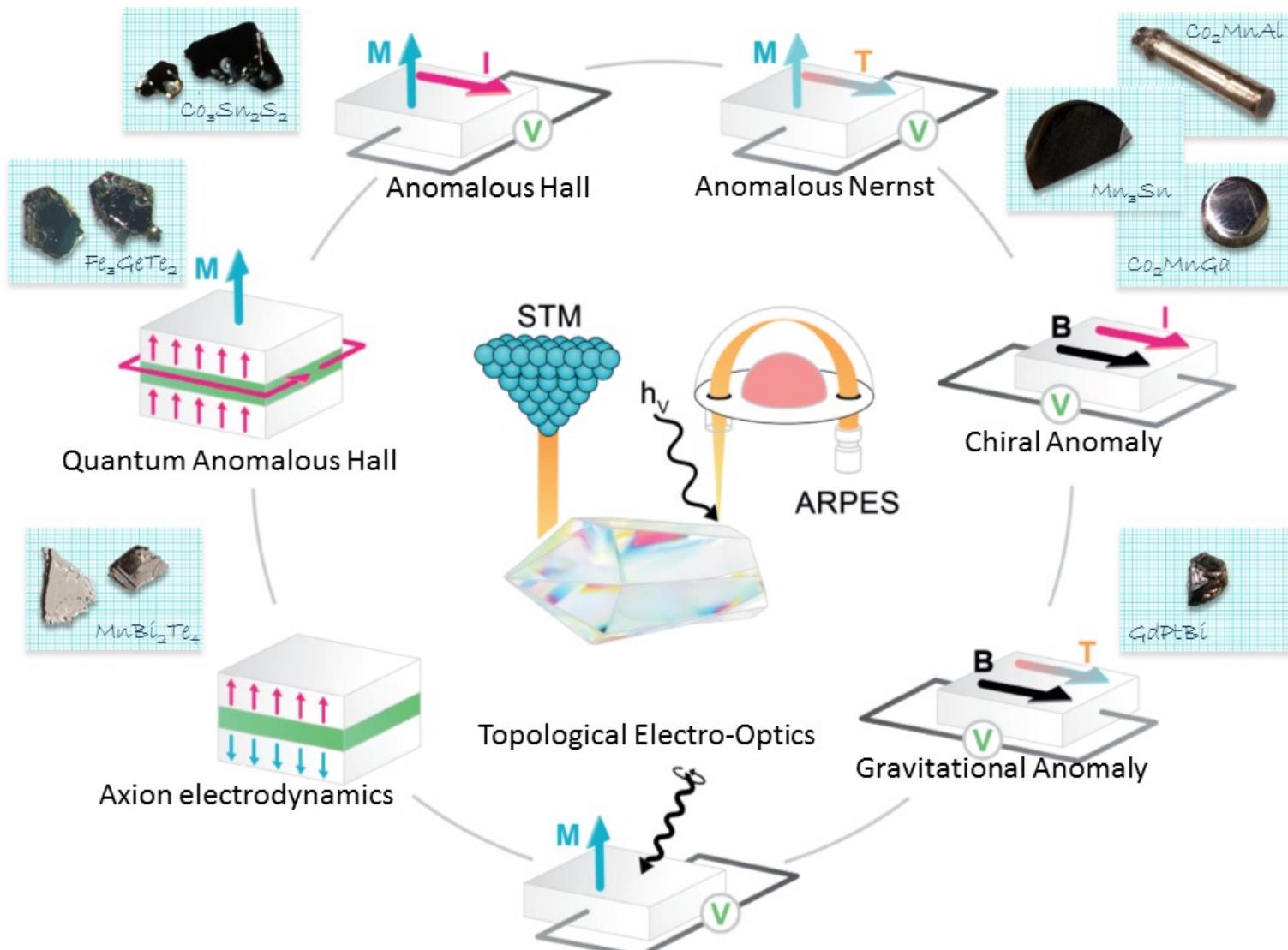
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SCIENCE

An Experiment in Zurich Brings Us Nearer to a Black Hole's Mysteries

By KENNETH CHANG JULY 19, 2017

magnetic Weyl



chiral anomaly in
 $\text{Co}_3\text{Sn}_2\text{S}_2$

antiferromagnetic topological materials

Article

High-throughput calculations of magnetic topological materials

Table 3 | The magnetic topological materials identified in this work

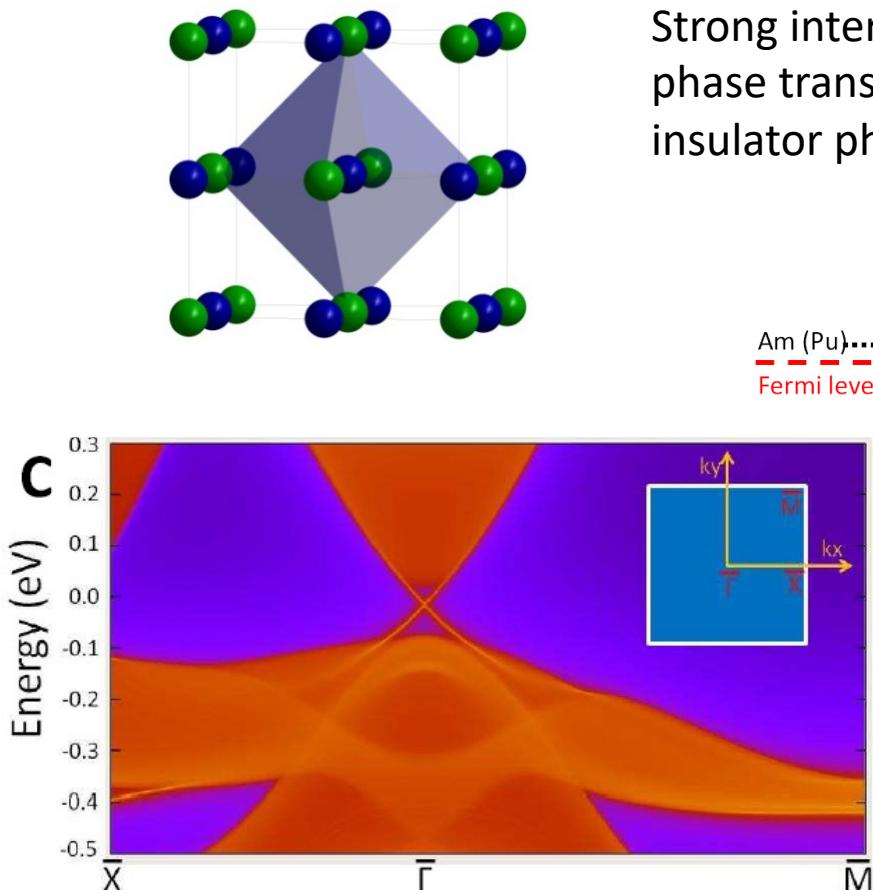
Categories	Properties	Materials
I-A	Non-collinear manganese compounds	Mn ₃ GaC, Mn ₃ ZnC, Mn ₃ CuN, Mn ₃ Sn, Mn ₃ Ge, Mn ₃ Ir, Mn ₃ Pt, Mn ₅ Si ₃
I-B	Actinide intermetallic	UNiGa ₅ , UPtGa ₅ , NpRhGa ₅ , NpNiGa ₅
I-C	Rare-earth intermetallic	NdCo ₂ , TbCo ₂ , NpCo ₂ , PrAg DyCu, NdZn, TbMg, NdMg, Nd ₅ Si ₄ , Nd ₅ Ge ₄ , Ho ₂ RhIn ₈ , Er ₂ CoGa ₈ , Nd ₂ RhIn ₈ , Tm ₂ CoGa ₈ , Ho ₂ RhIn ₈ , DyCo ₂ Ga ₈ , TbCo ₂ Ga ₈ , Er ₂ Ni ₂ In, CeRu ₂ Al ₁₀ , Nd ₃ Ru ₄ Al ₁₂ , Pr ₃ Ru ₄ Al ₁₂ , ScMn ₆ Ge ₆ , YFe ₄ Ge ₄ , LuFe ₄ Ge ₄ , CeCoGe ₃
II-A	Metallic iron pnictides	LaFeAsO, CaFe ₂ As ₂ , EuFe ₂ As ₂ , BaFe ₂ As ₂ , Fe ₂ As, CaFe ₄ As ₃ , LaCrAsO, Cr ₂ As, CrAs, CrN
II-B	Semiconducting manganese pnictides	BaMn ₂ As ₂ , BaMn ₂ Bi ₂ , CaMnBi ₂ , SrMnBi ₂ , CaMn ₂ Sb ₂ , CuMnAs, CuMnSb, Mn ₂ As
II-C	Rare-earth intermetallic compounds with the composition 1:2:2	PrNi ₂ Si ₂ , YbCo ₂ Si ₂ , DyCo ₂ Si ₂ , PrCo ₂ P ₂ , CeCo ₂ P ₂ , NdCo ₂ P ₂ , DyCu ₂ Si ₂ , CeRh ₂ Si ₂ , UAu ₂ Si ₂ , U ₂ Pd ₂ Sn, U ₂ Pd ₂ In, U ₂ Ni ₂ Sn, U ₂ Ni ₂ In, U ₂ Rh ₂ Sn
II-D	Rare-earth ternary compounds of the composition 1:1:1	CeMgPb, PrMgPb, NdMgPb, TmMgPb
III-A	Semiconducting actinides/rare-earth pnictides	HoP, UP, UP ₂ , UAs, NpS, NpSe, NpTe, NpSb, NpBi, U ₃ As ₄ , U ₃ P ₄
III-B	Metallic oxides	Ag ₂ NiO ₂ , AgNiO ₂ , Ca ₃ Ru ₂ O ₇ , Double perovskite Sr ₃ ColrO ₆
III-C	Metal-to-insulator transition compounds	NiS ₂ , Sr ₂ Mn ₃ As ₂ O ₂
III-D	Semiconducting and insulating oxides, borates, hydroxides, silicates and phosphate	LuFeO ₃ , PdNiO ₃ , ErVO ₃ , DyVO ₃ , MnGeO ₃ , Tm ₂ Mn ₂ O ₇ , Yb ₂ Sn ₂ O ₇ , Tb ₂ Sn ₂ O ₇ , Ho ₂ Ru ₂ O ₇ , Er ₂ Ti ₂ O ₇ , Tb ₂ Ti ₂ O ₇ , Cd ₂ Os ₂ O ₇ , Ho ₂ Ru ₂ O ₇ , Cr ₂ ReO ₆ , NiCr ₂ O ₄ , MnV ₂ O ₄ , Co ₂ SiO ₄ , Fe ₂ SiO ₄ , PrFe ₃ (BO ₃) ₄ , KCo ₄ (PO ₄) ₃ , CoPS ₃ , SrMn(VO ₄)(OH), Ba ₅ Co ₅ ClO ₁₃ , Fel ₂

antiferromagnetic topological materials

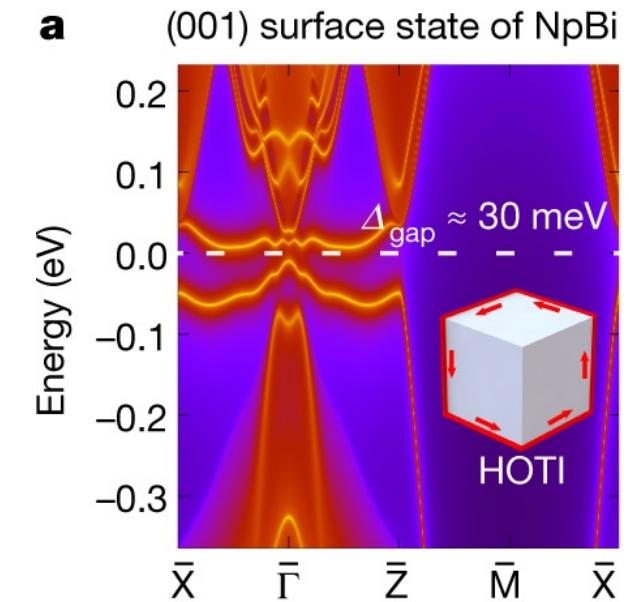
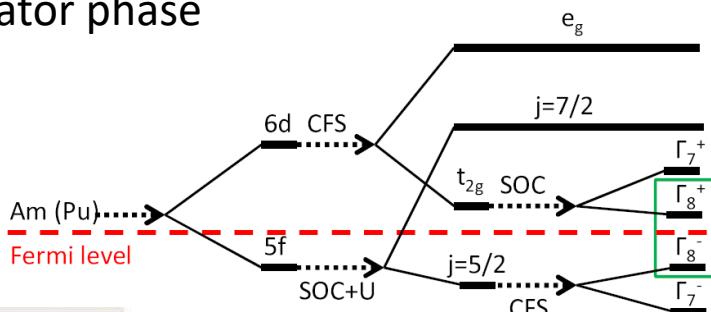
III-A

Semiconducting actinides/
rare-earth pnictides

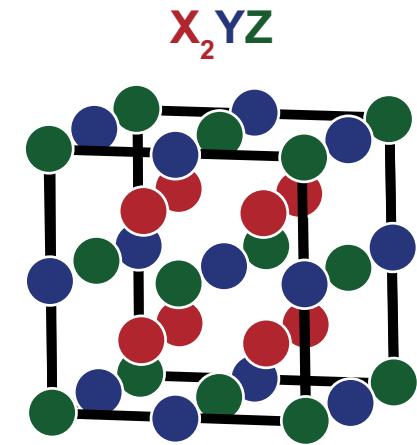
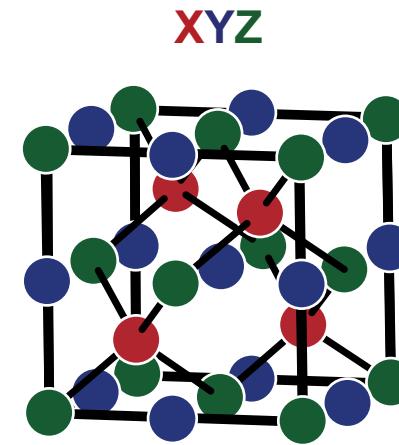
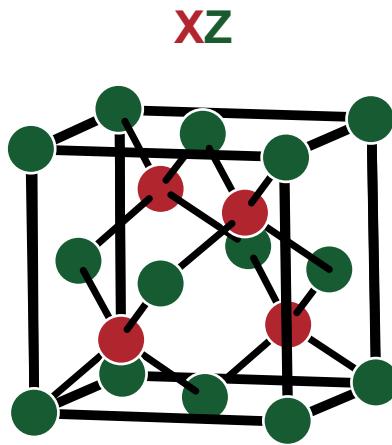
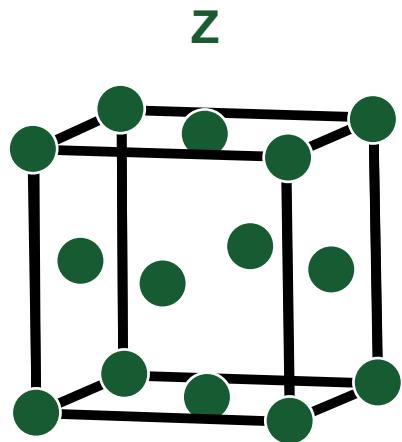
HoP, UP, UP₂, UAs, NpS, NpSe, NpTe, NpSb, NpBi, U₃As₄, U₃P₄



Strong interaction drives a quantum phase transition to a topological insulator phase



Heusler Verbindungen



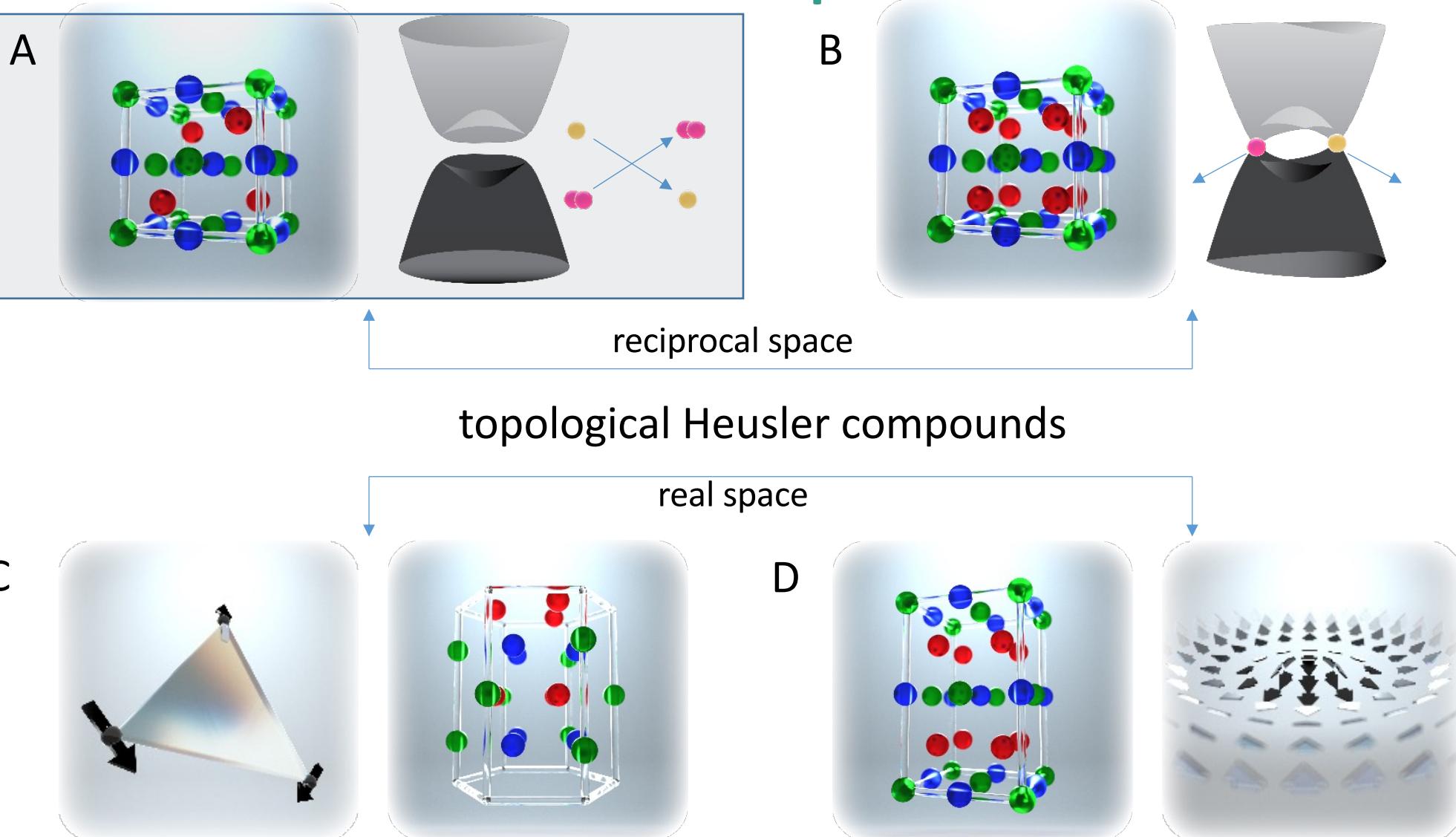
C	N
2.55	3.04
Si	P
1.90	2.19
Ge	As
2.01	2.18
Sn	Sb
1.96	2.05

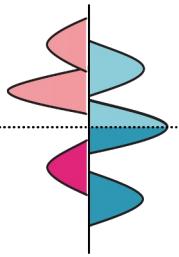
H	2.20
Li	0.98
Be	1.57
Na	0.93
Mg	1.31
K	0.82
Ca	1.00
Sc	1.36
Ti	1.54
V	1.63
Cr	1.66
Mn	1.55
Fe	1.83
Co	1.88
Ni	1.91
Cu	1.90
Zn	1.65
Ga	1.81
Ge	2.01
As	2.18
Se	2.55
Br	2.96
Kr	3.00

XYZ Heusler compounds

B	C	N	O	F	Ne
2.04	2.55	3.04	3.44	3.98	
Al	Si	P	S	Cl	Ar
1.61	1.90	2.19	2.58	3.16	
Sc	Ti	V	Cr	Mn	Fe
1.36	1.54	1.63	1.66	1.55	1.83
Y	Zr	Nb	Mo	Tc	Ru
1.22	1.33	1.60	2.16	1.90	2.20
Rh	Pd	Ag	Cd	In	Sn
2.28	2.20	1.93	1.69	1.78	1.96
Ir	Pt	Au	Hg	Tl	Pb
2.20	2.20	2.40	1.90	1.80	1.90
La	Ce	Pr	Nd	Pm	Sm
1.10	1.12	1.13	1.14	1.13	1.17
Eu	Gd	Tb	Dy	Ho	Er
1.20	1.20	1.10	1.22	1.23	1.24
Tm	Er	Tm	Yb	Lu	
1.25	1.25	1.10	1.27		
Ac	Th	Pa	U	Np	Pu
1.10	1.30	1.50	1.70	1.30	1.28
Am	Cm	Bk	Cf	Es	Fm
1.13	1.13	1.28	1.30	1.30	1.30
Md	No	Md	Lr		
1.30	1.30	1.30	1.30		

Heusler compounds



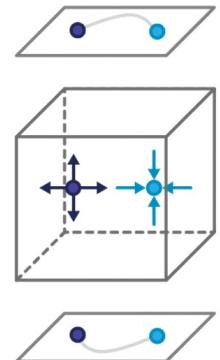
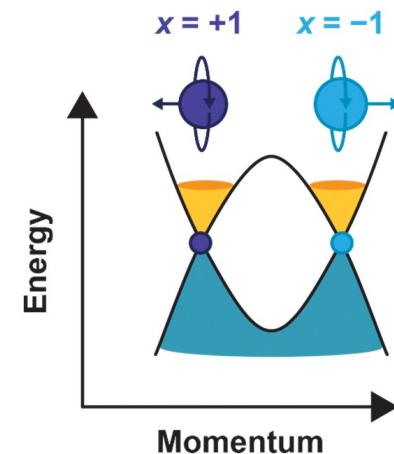
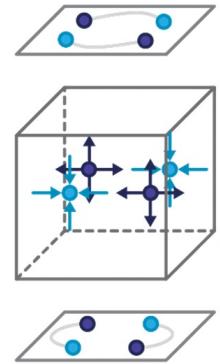
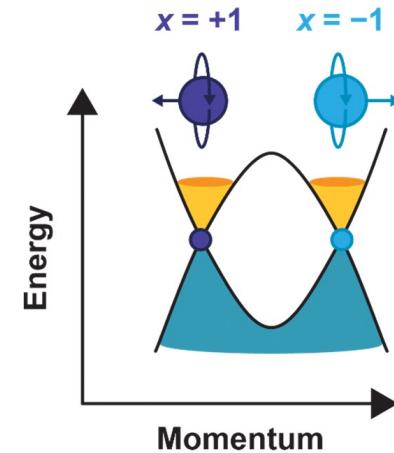
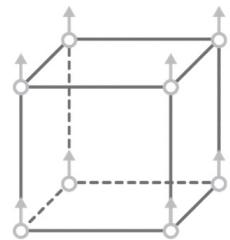


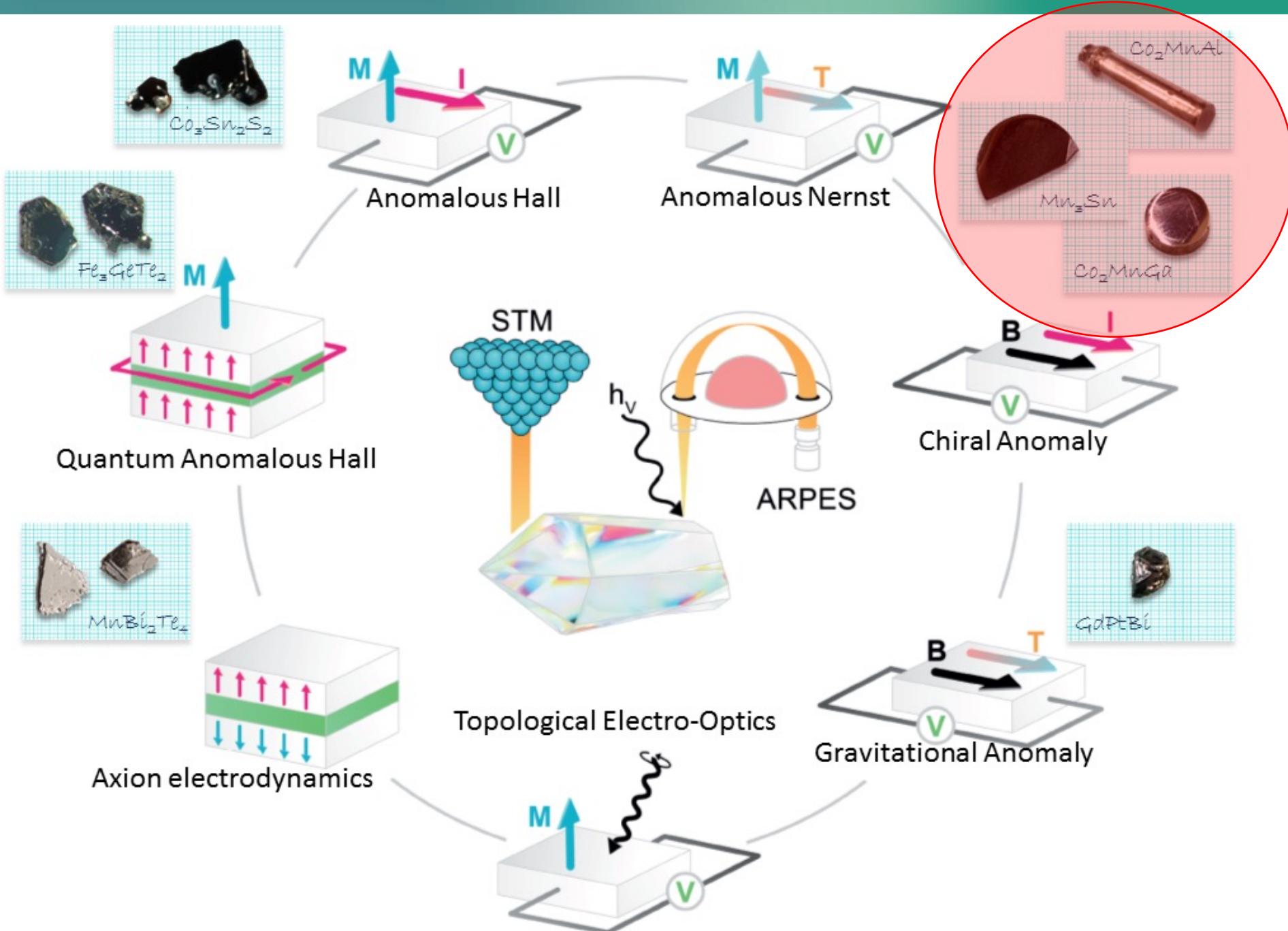
magnetic Weyl semimetals

Weyl points at low symmetry points

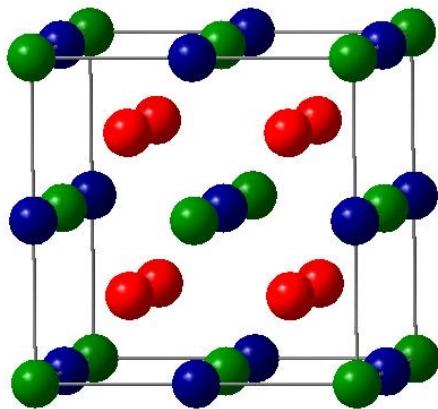
breaking time reversal symmetry

- magnetic field
- **all crossings in the band structure in ferromagnets are Weyl points**





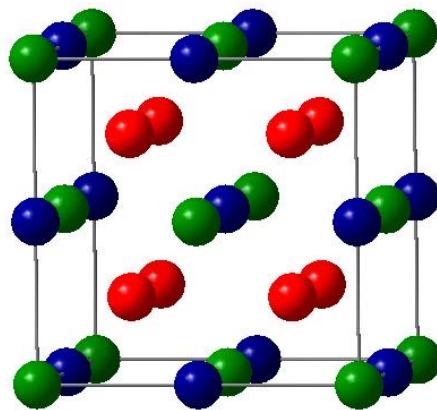
Heusler compounds



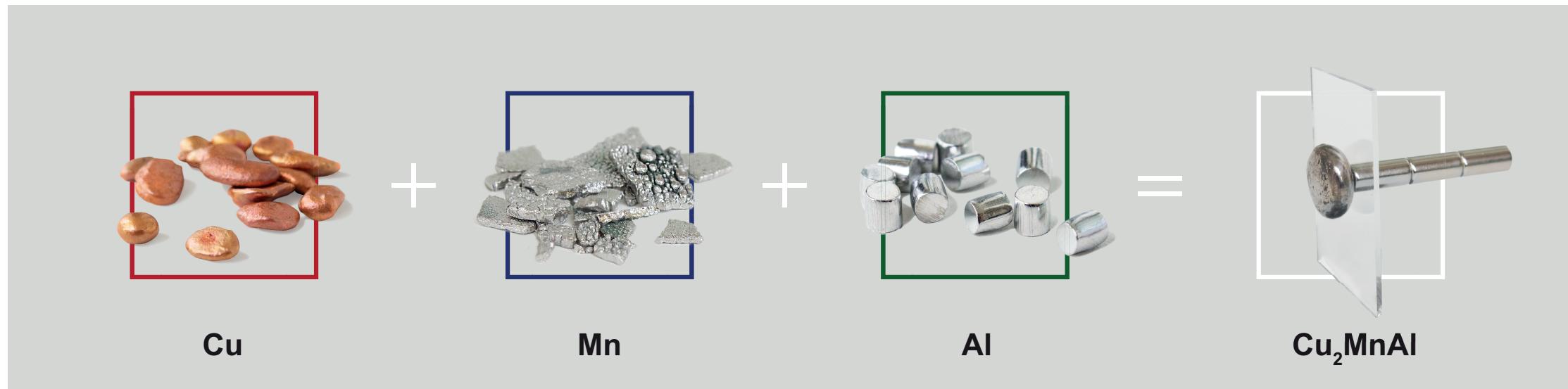
Heusler X_2YZ L₂1



Heusler Verbindungen

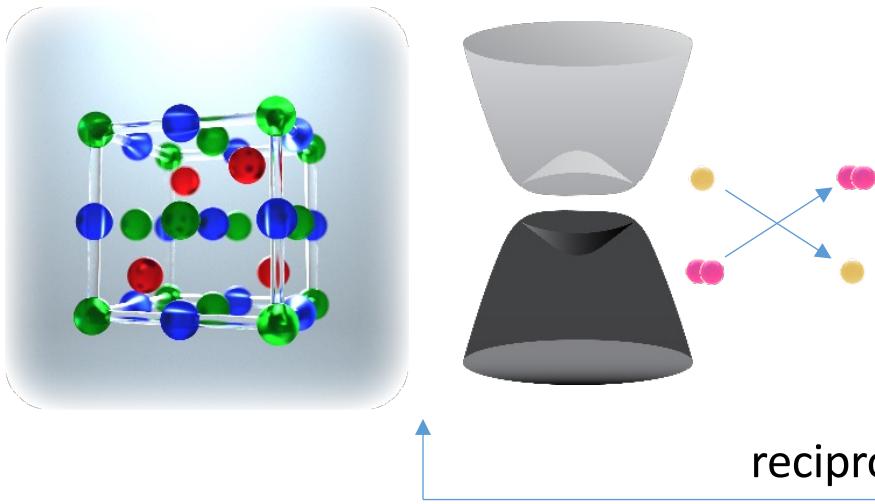


Heusler X_2YZ $L2_1$

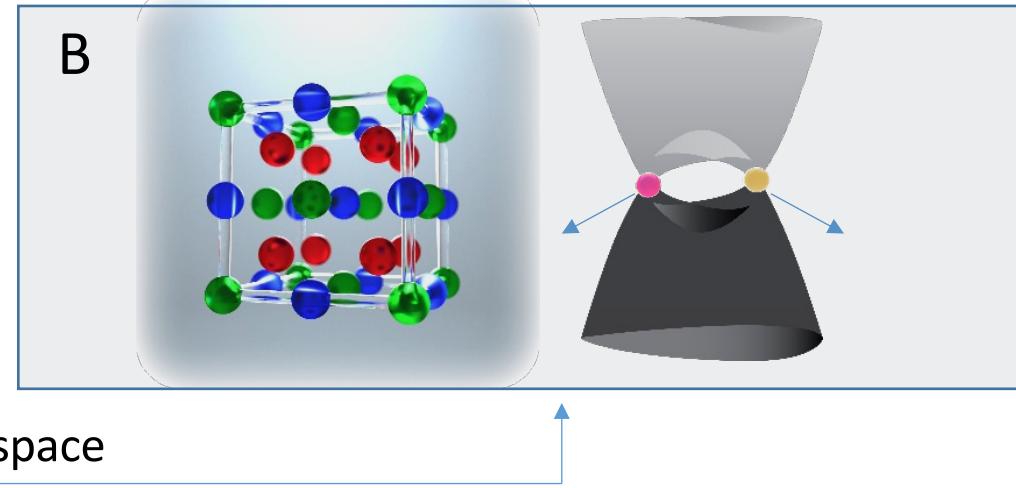


Heusler compounds

A



B

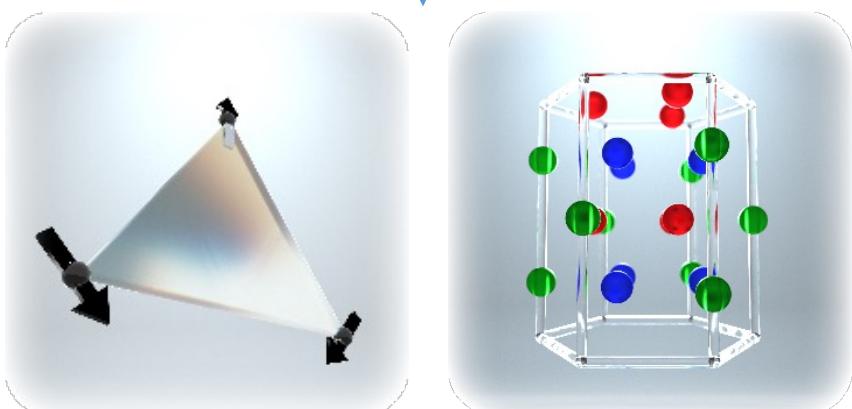


reciprocal space

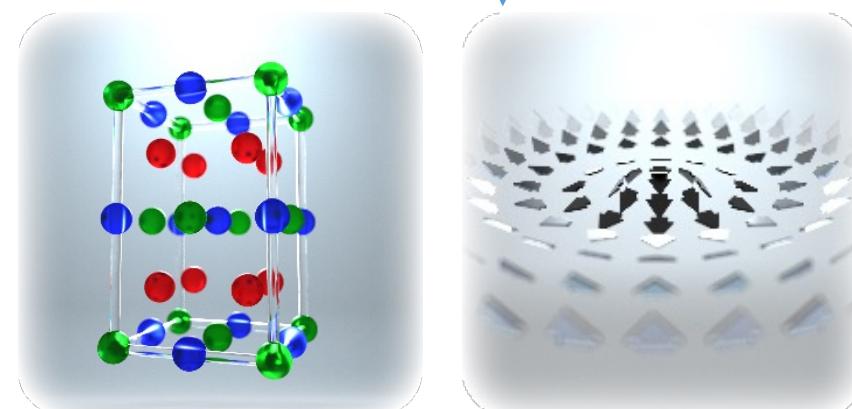
topological Heusler compounds

real space

C

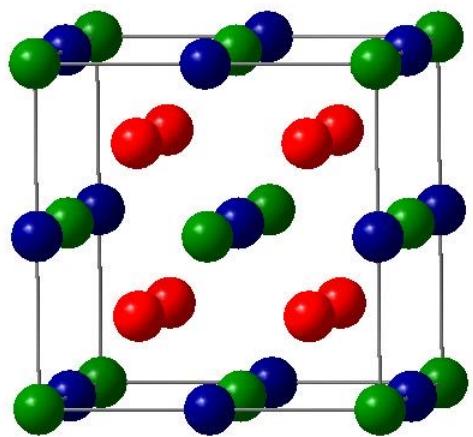


D



magnetic Heusler

X_2YZ



Fe₂

V

Al

$$2 * 8 + 5 + 3 = 24$$

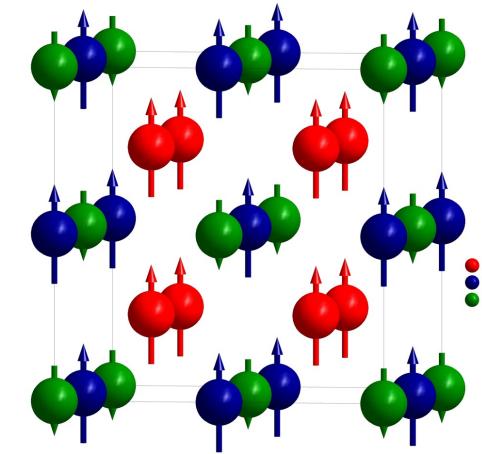
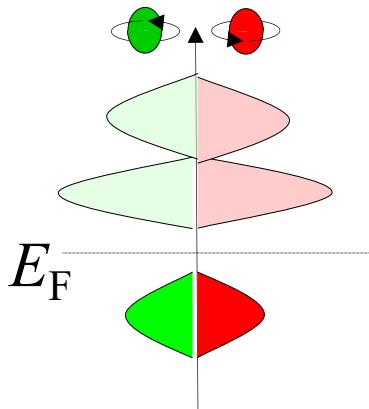


Kübler et al., PRB **28**, 1745 (1983)

de Groot RA, et al. PRL **50** 2024 (1983)

Galanakis et al., PRB **66**, 012406 (2002)

- magic valence electron number: 24
- valence electrons = 24 + magnetic moments

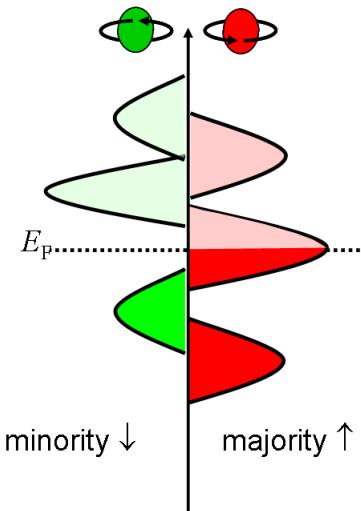


Co₂

Mn

Ga

$$2 * 9 + 7 + 3 = 24 + 4$$



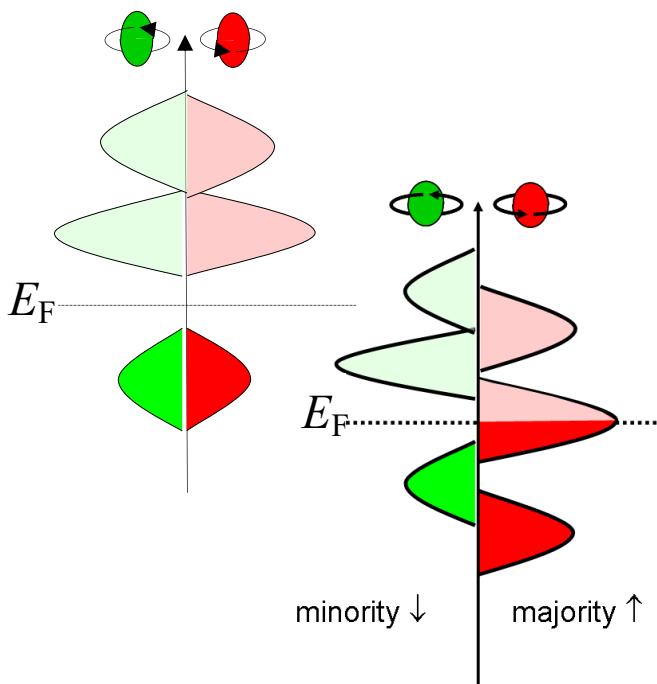
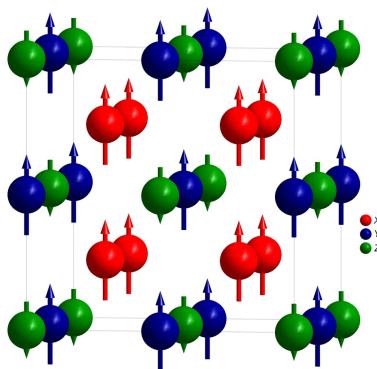
Kandpal et al., J. Phys. D **40** (2007) 1507

Balke et al. Solid State Com. **150** (2010) 529

Kübler et al., Phys. Rev. B **76** (2007) 024414

half metallic Heusler

X_2YZ



Kübler *et al.*, PRB **28**, 1745 (1983)

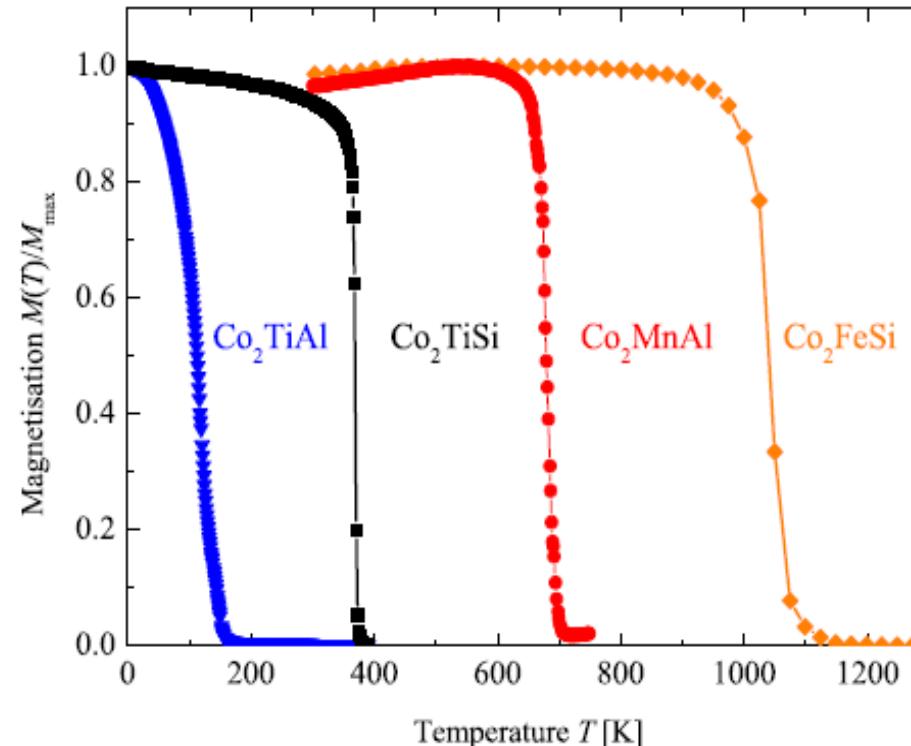
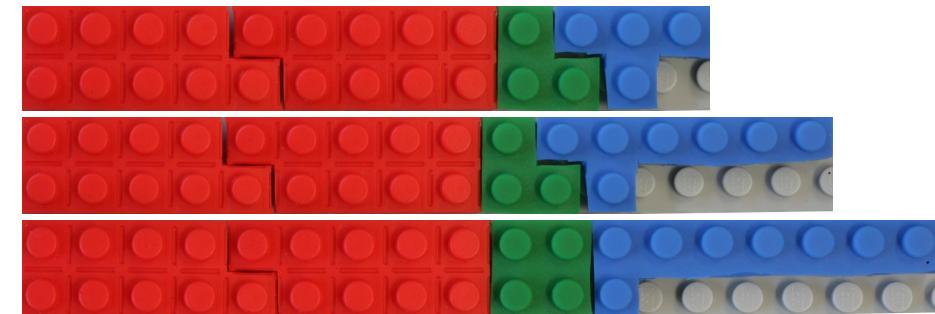
de Groot RA, *et al.* PRL **50** 2024 (1983)

Galanakis *et al.*, PRB **66**, 012406 (2002)

Co_2TiAl : $2 \times 9 + 4 + 3 = 25$ $M_s = 1m_B$

Co_2MnGa : $2 \times 9 + 7 + 3 = 28$ $M_s = 4m_B$

Co_2FeSi : $2 \times 9 + 8 + 4 = 30$ $M_s = 6m_B$

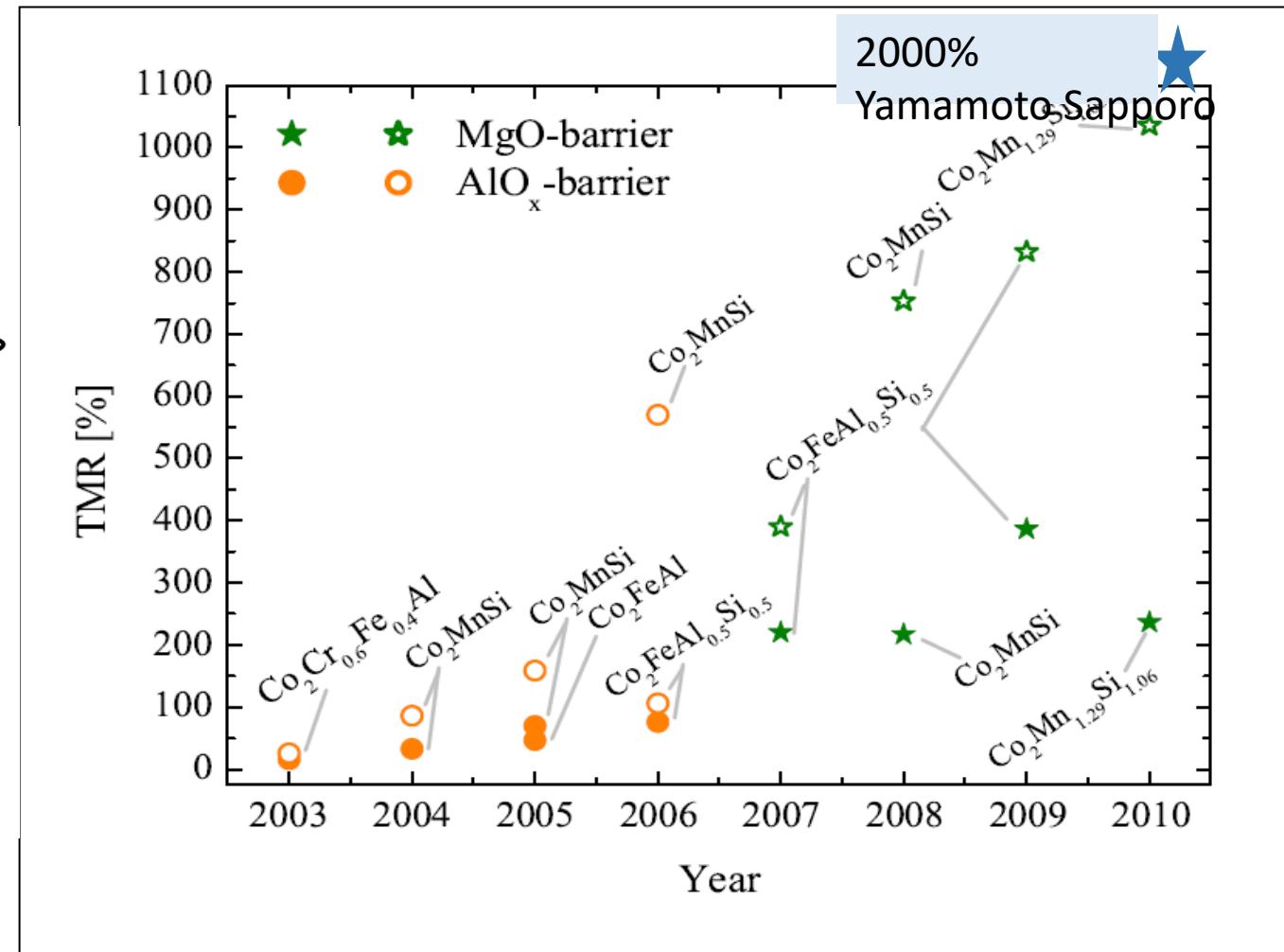
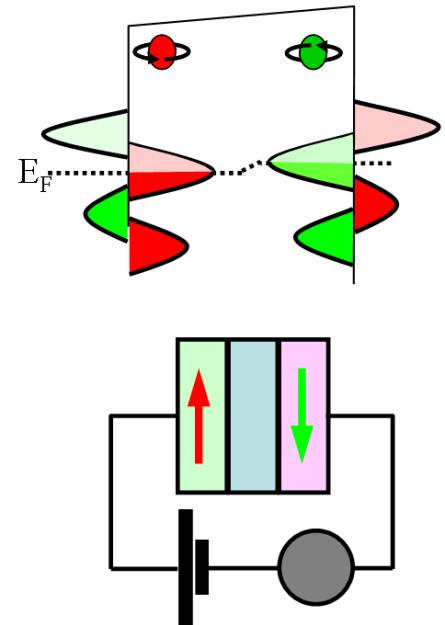
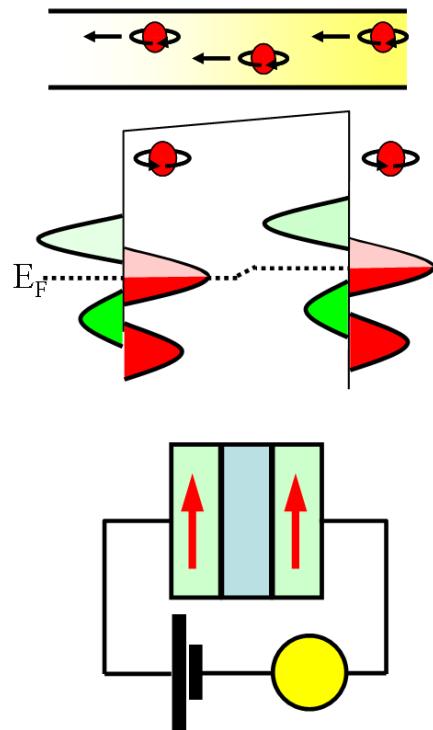


Il et al., J. Phys. D **40** (2007) 1507

Liu et al., Solid State Com. **150** (2010) 529

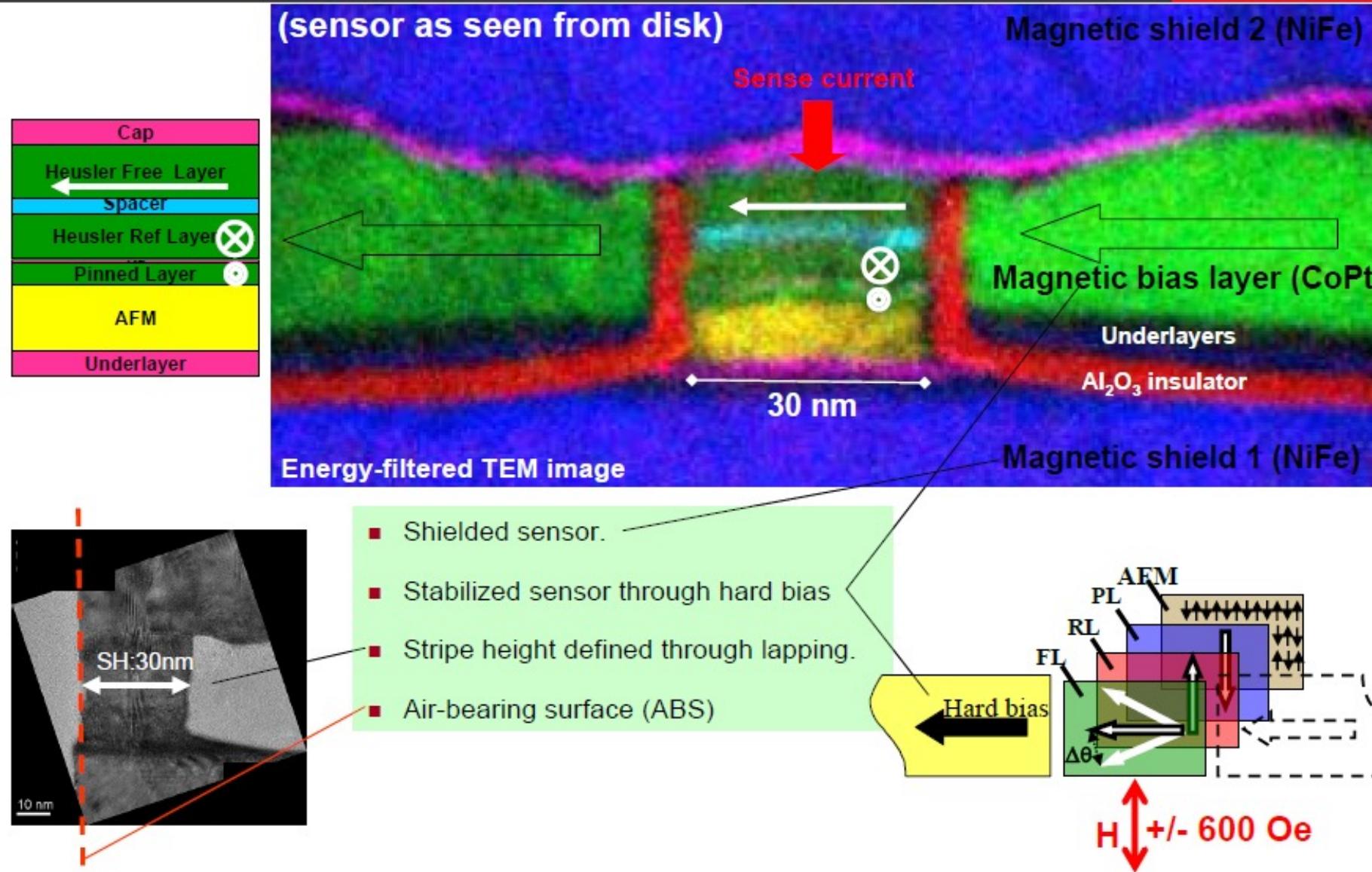
Kübler et al., Phys. Rev. B **76** (2007) 024414

Tunneling magnetoresistance effect in Heusler

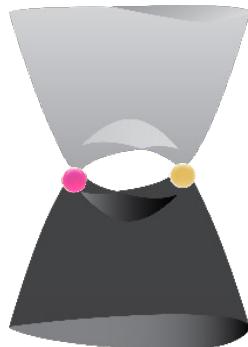


Read-Head Fabrication w/Heusler alloy

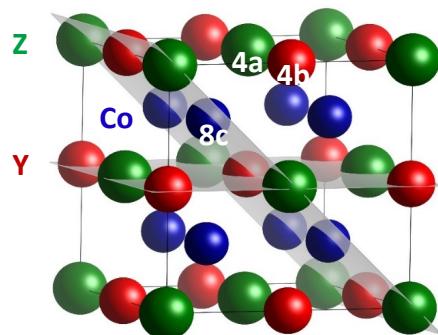
HITACHI
Inspire the Next



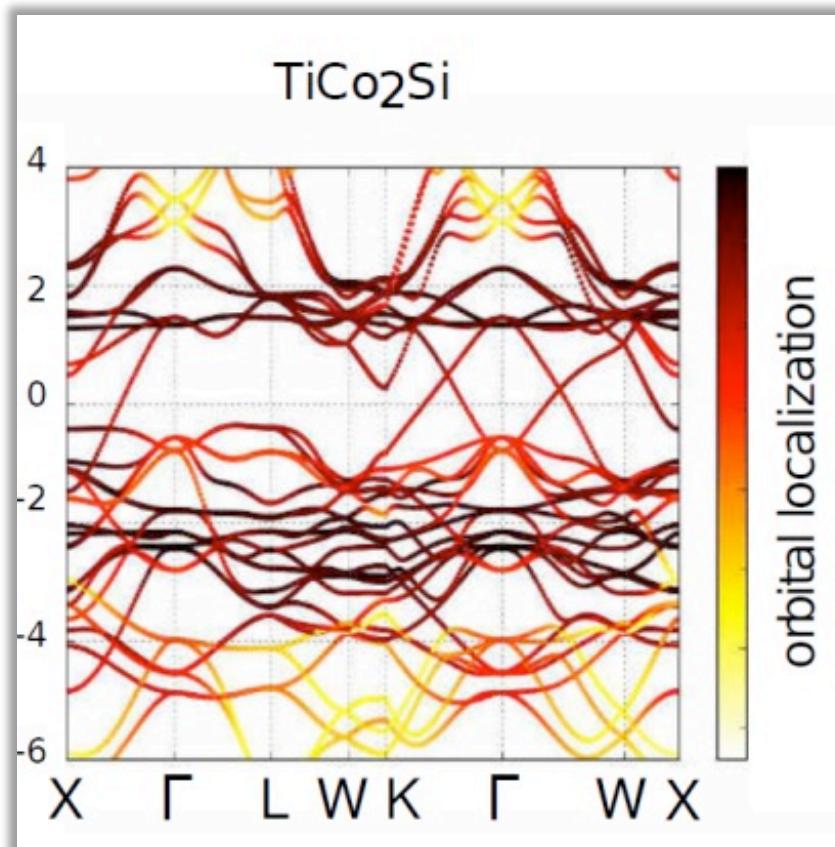
semimetallic Heusler compounds



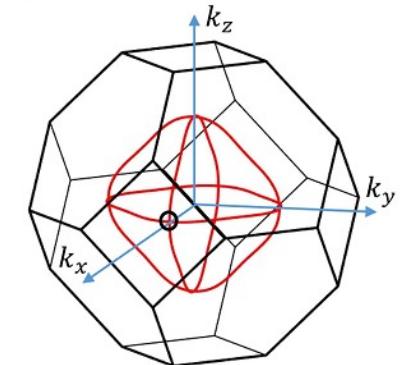
Co_2TiSi



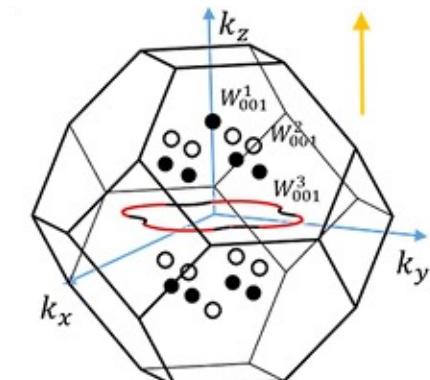
$$\text{Co}_2\text{TiSi}: 2 \times 9 + 4 + 4 = 26 \quad \text{Ms} = 2\mu_B$$



Phys. Rev. Lett. 117, 236401 (2016)
Sci. Rep. 6, 38839 (2016)



Symmetry and electronic structures depend on the magnetization direction
With SOC



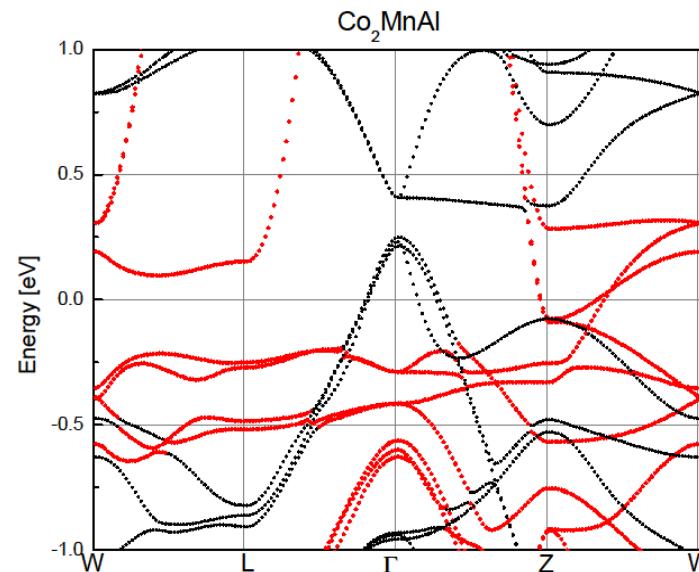
Heusler, Weyl and Berry



Giant AHE

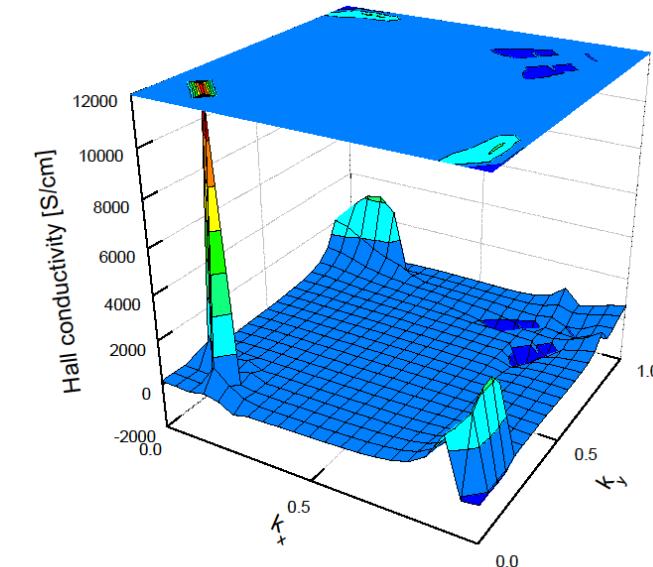
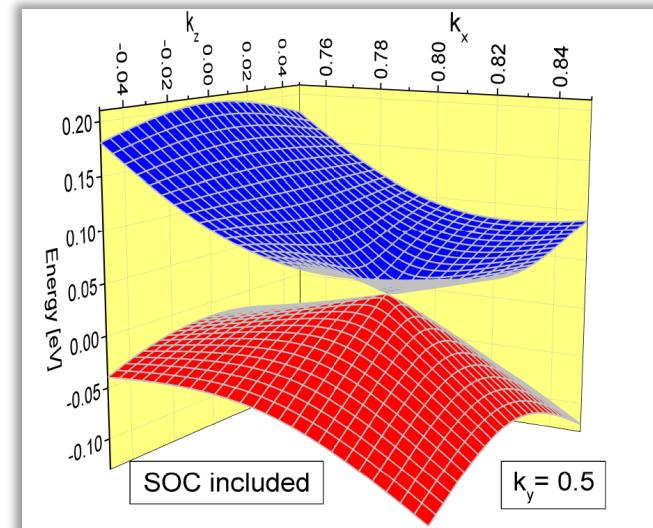
Co₂MnAl

$\sigma_{xy} = 1800 \text{ S/cm}$ calc.
 $\sigma_{xy} \approx 2000 \text{ S/cm}$ meas.

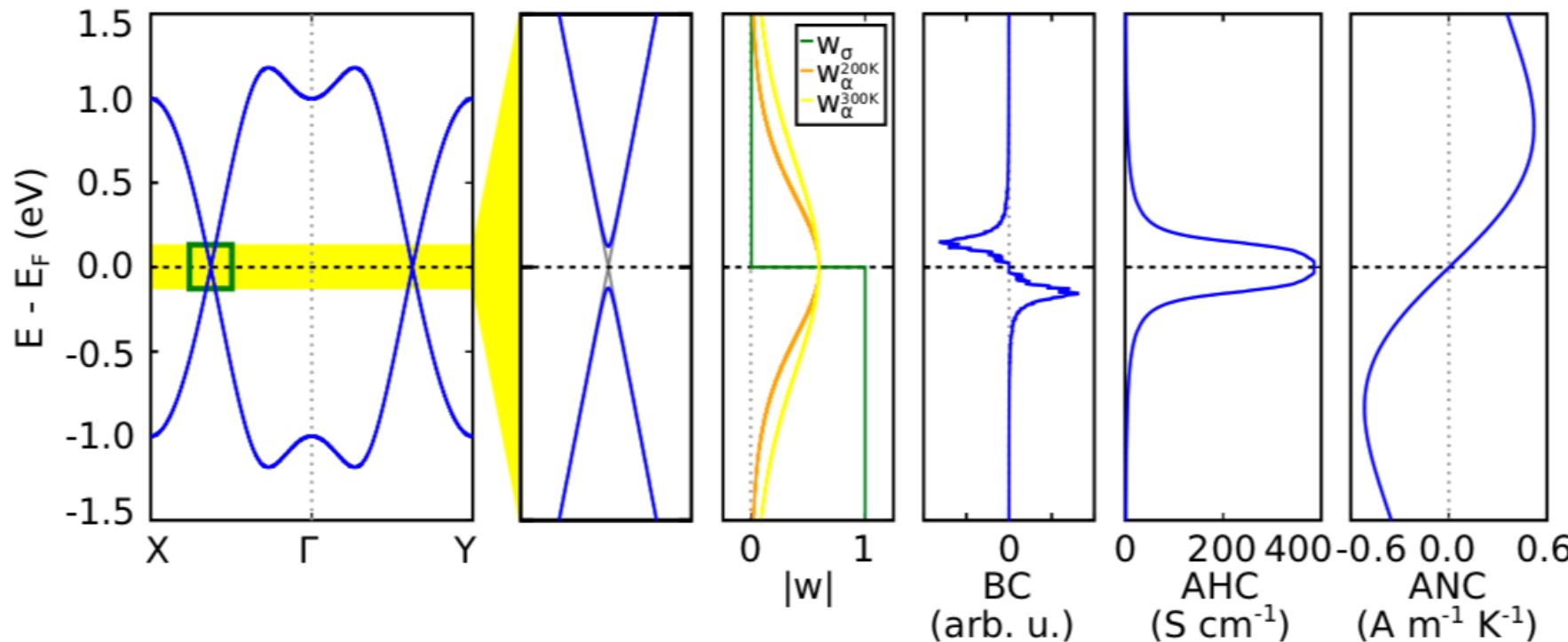


Nodal lines
and Weyl
points

Berry
curvature
enhancement

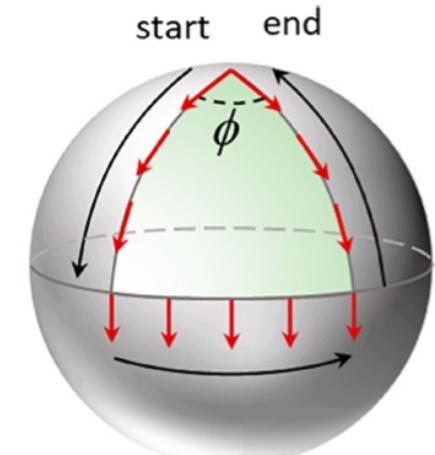


Berry curvature and anomalous Nernst



$$\sigma_{ij} = \frac{e^2}{\hbar} \sum_n \int \frac{d^3 k}{(2\pi)^3} \Omega_{ij}^{nk} f_{nk}$$

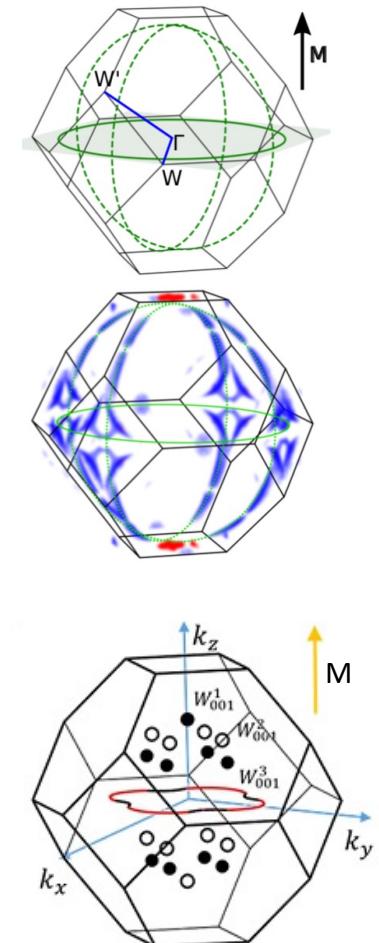
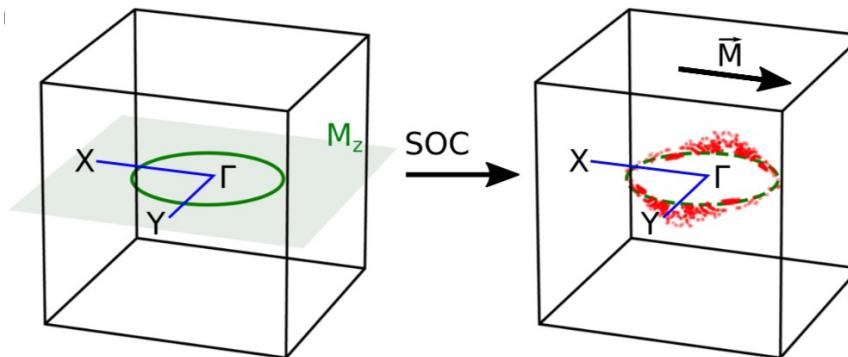
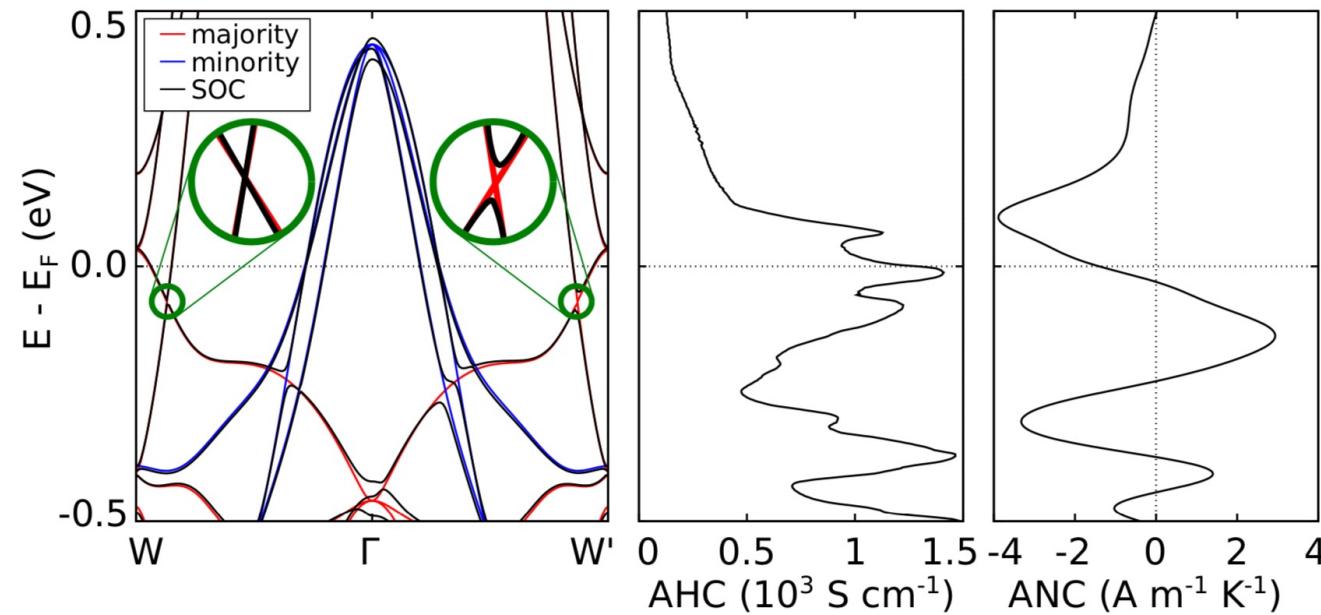
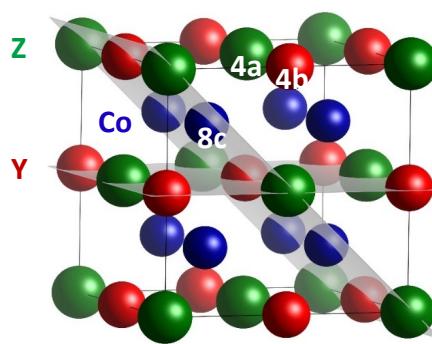
$$\alpha_{ij} = \frac{e^2}{\hbar} \sum_n \int \frac{d^3 k}{(2\pi)^3} \Omega_{ij}^{nk} - \frac{1}{eT} [(E_{nk} - E_F) f_{nk} + k_B T \ln (1 + \exp \frac{E_{nk} - E_F}{-k_B T})]$$



Berry curvature design

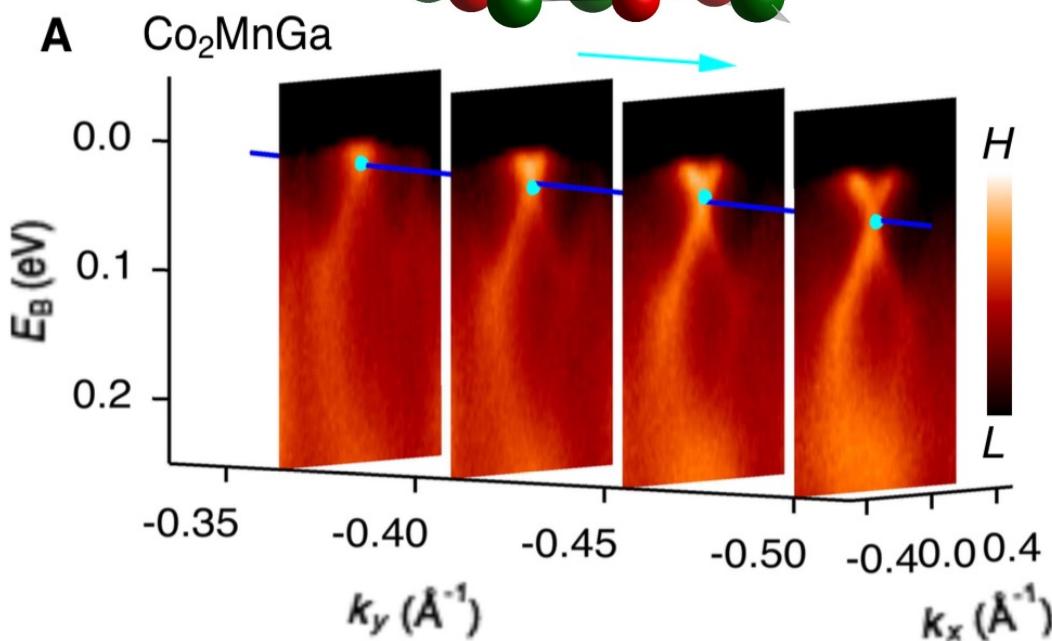
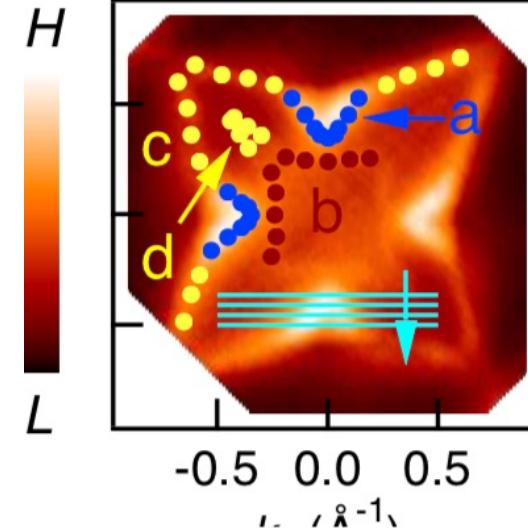
Co_2MnGa

- giant anomalous Hall
- giant anomalous Nernst



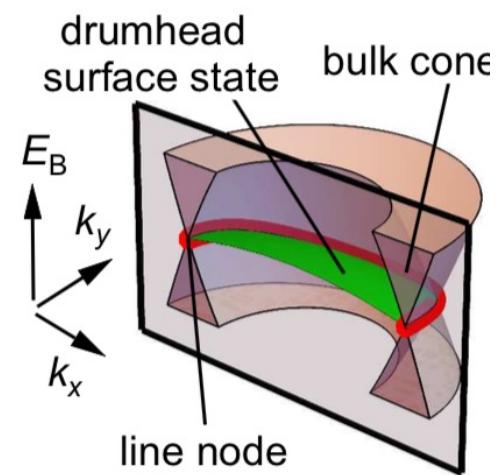
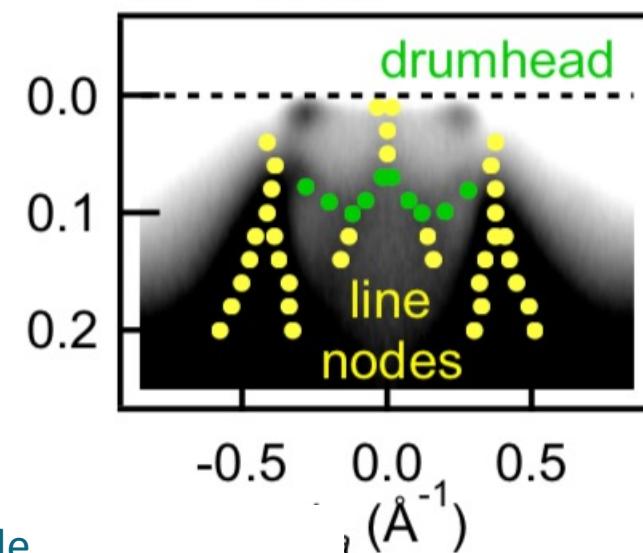
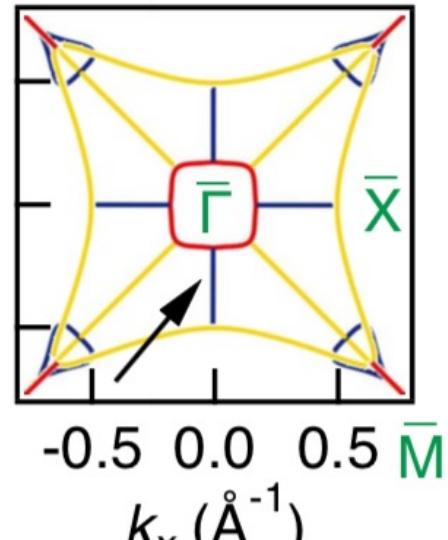


ARPES



Series of ARPES cuts through the candidate line node

DFT, Weyl lines

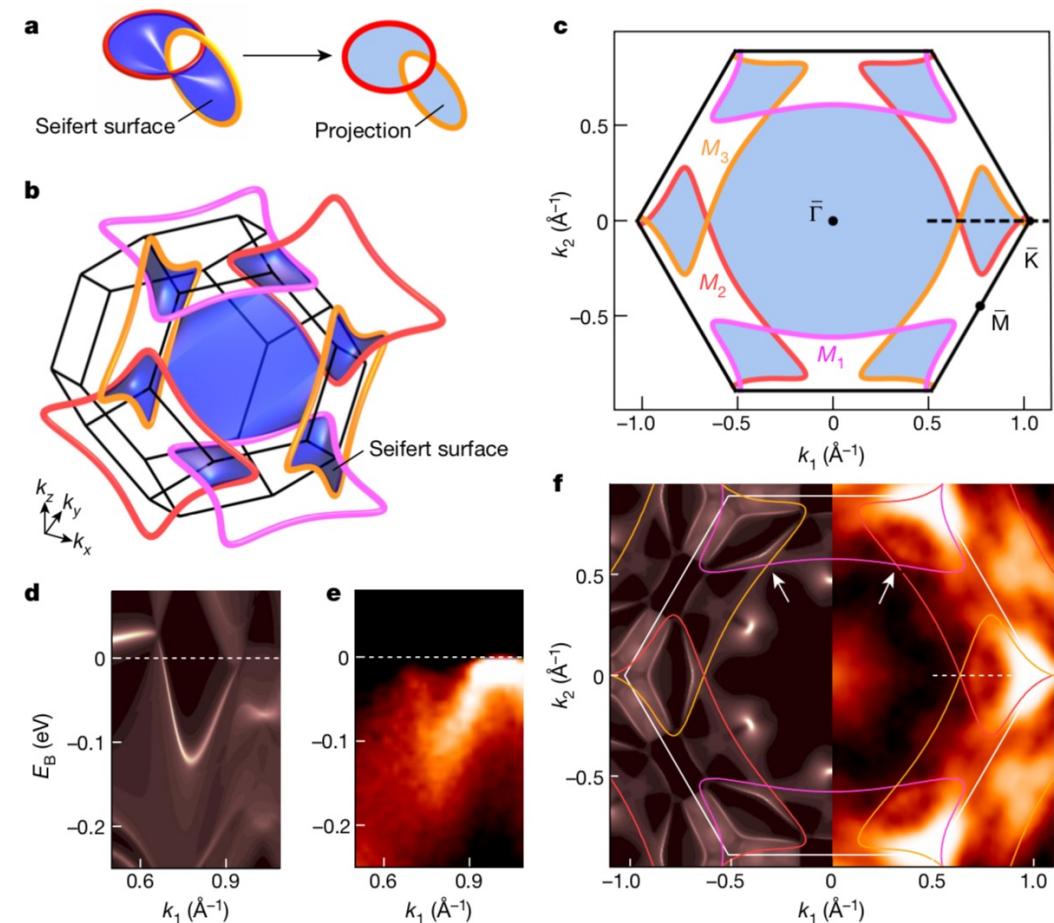
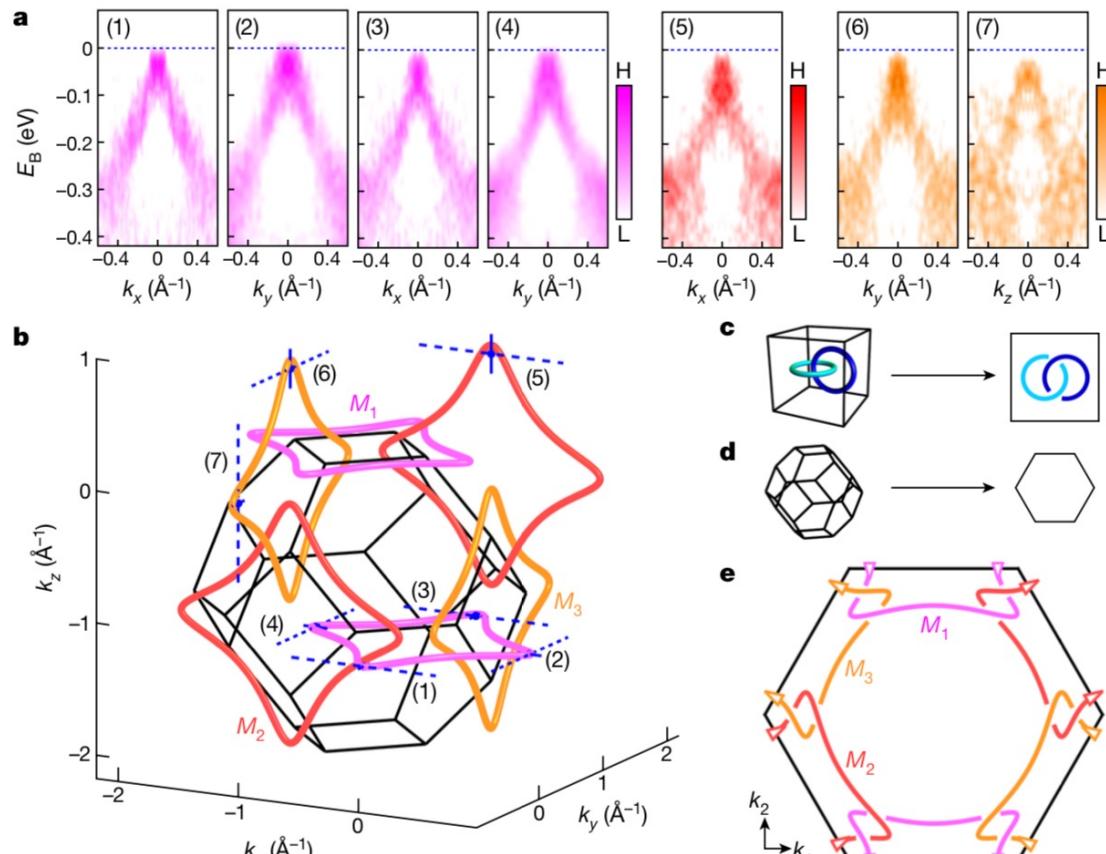




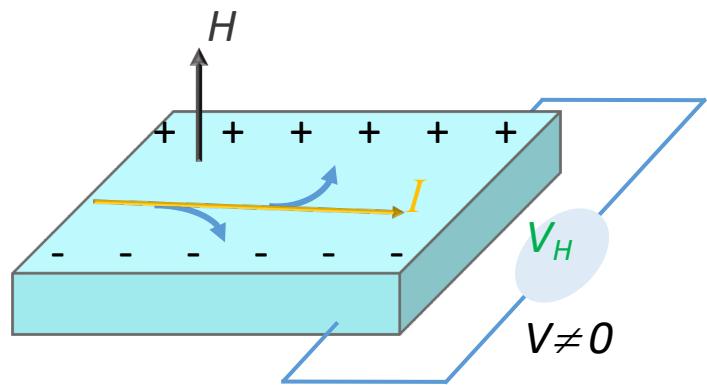
Article

Observation of a linked-loop quantum state in a topological magnet

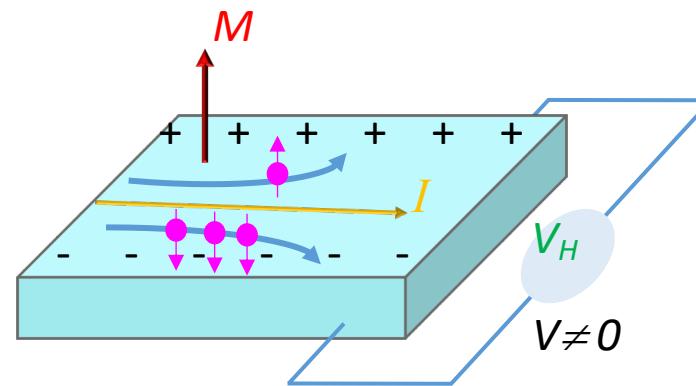
Co₂MnGa



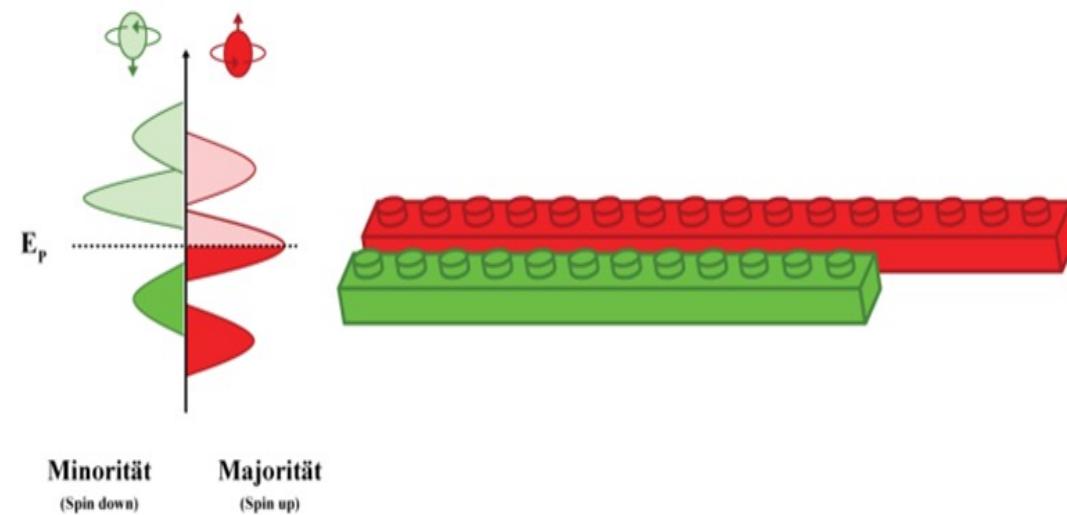
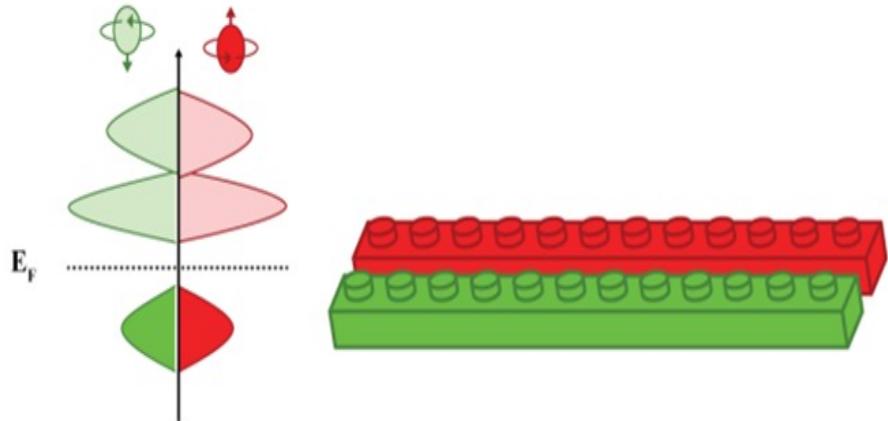
Hall effect



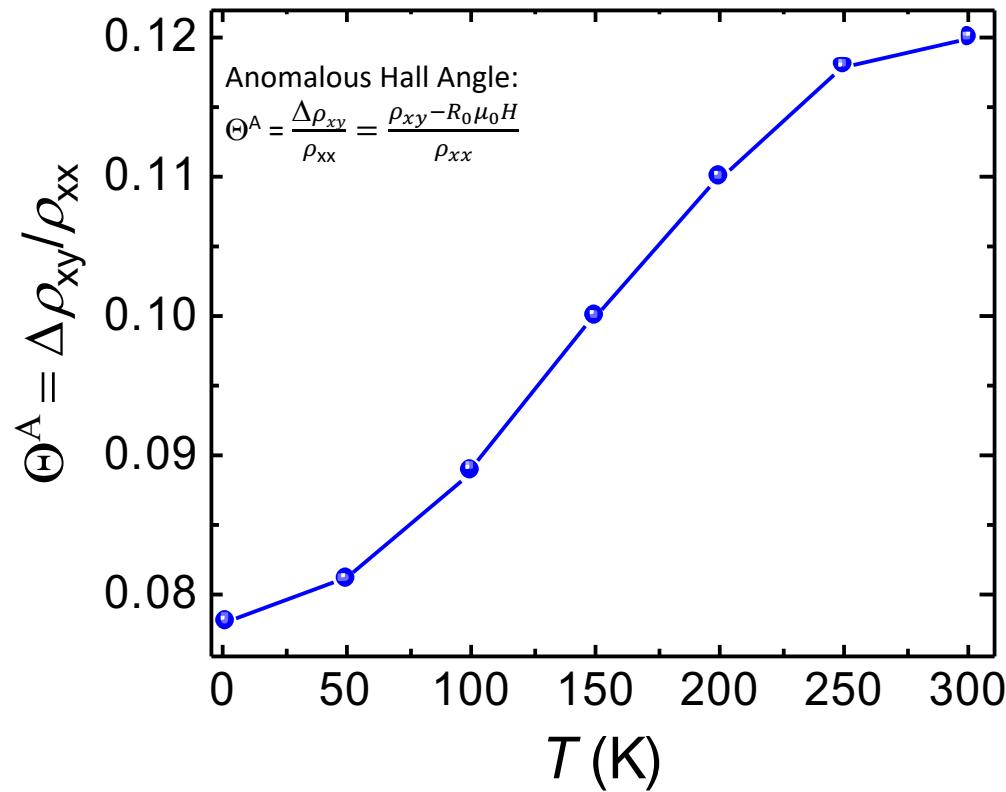
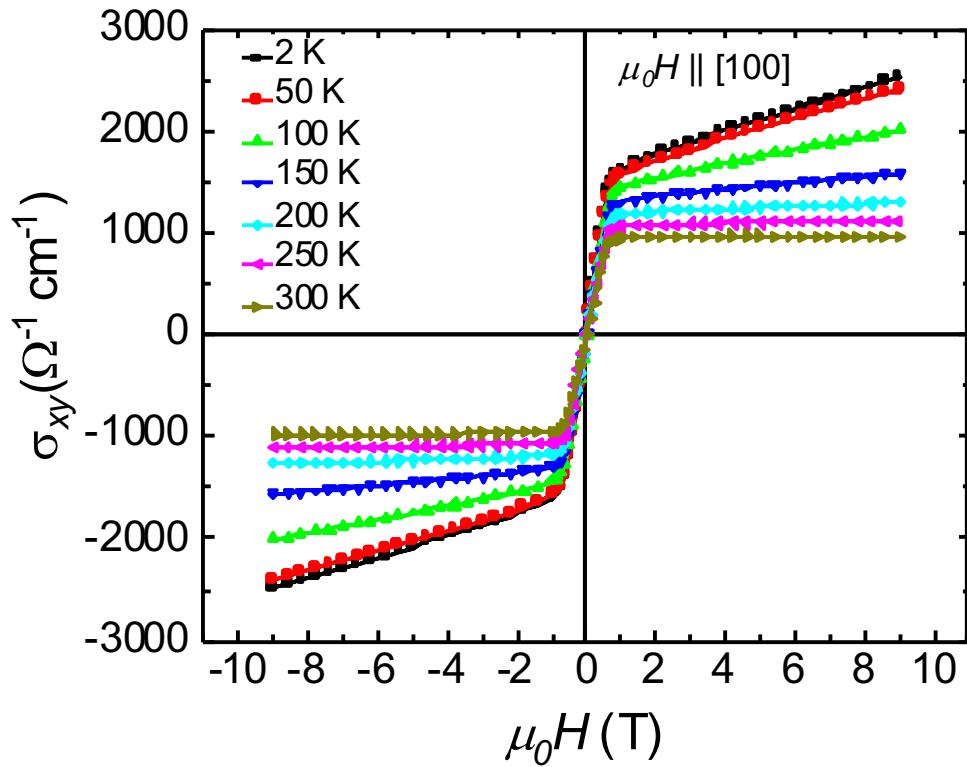
Hall Effect
diamagnetic semiconductor



Anomalous Hall Effect
ferromagnetic material



Anomalous Hall



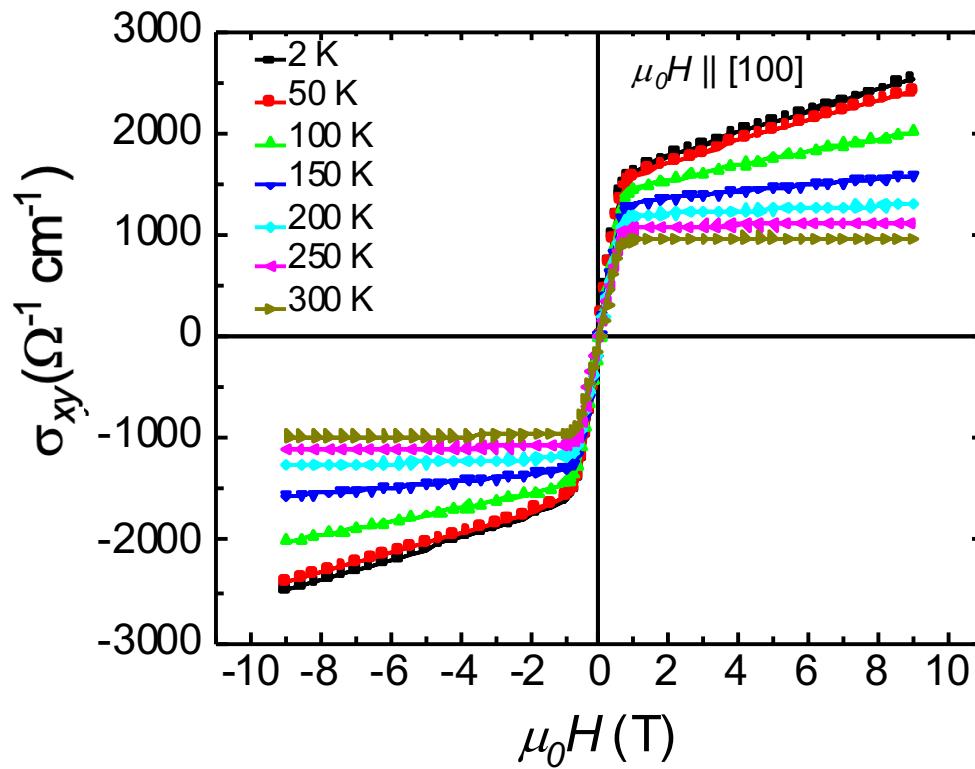
Co_2MnGa

Kübler, Felser, PRB 85 (2012) 012405

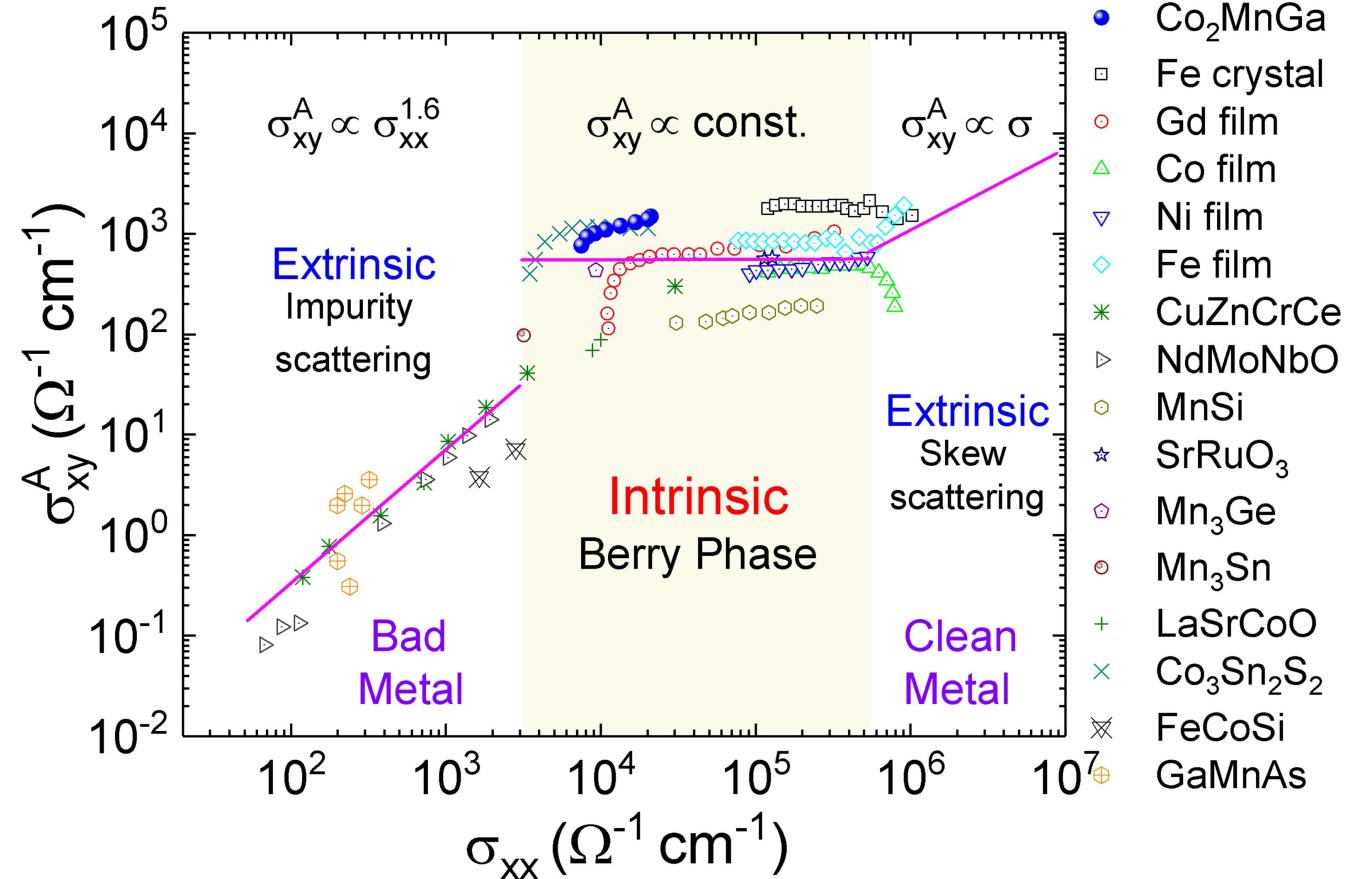
Vidal et al., Appl. Phys. Lett. 99 (2011) 132509

Manna et al., Phys. Rev. X 8 (2018) 041045, arXiv:1712.10174

Anomalous Hall



Co₂MnGa



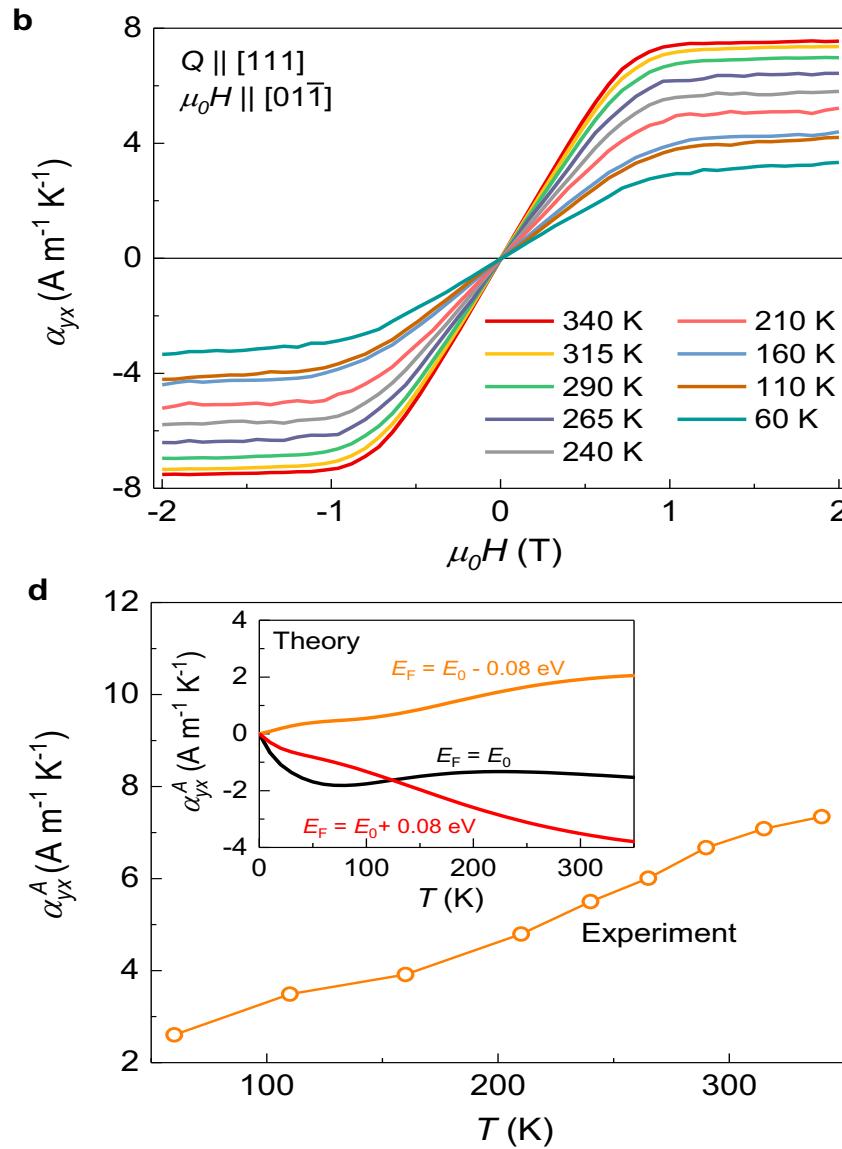
Naoto Nagaosa and Yoshinori Tokura 2012 Phys. Scr. 2012 014020

Kübler, Felser, PRB 85 (2012) 012405

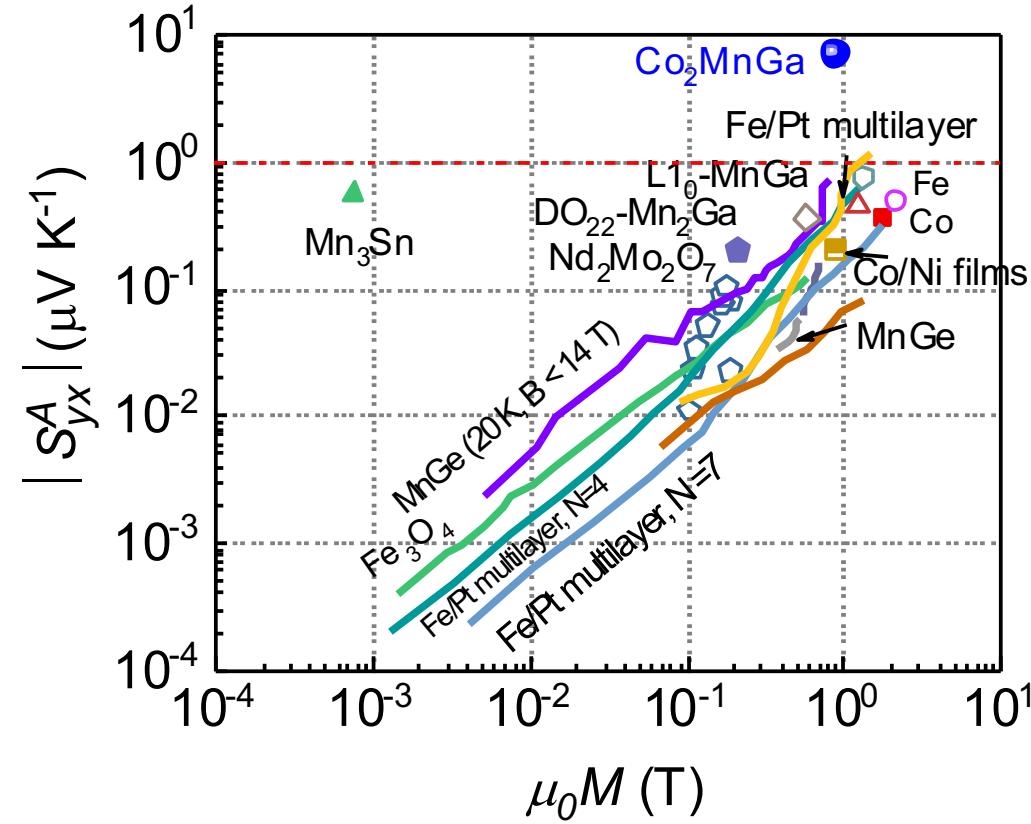
Vidal et al., Appl. Phys. Lett. 99 (2011) 132509

Manna et al., Phys. Rev. X 8 (2018) 041045, arXiv:1712.10174

anomalous Nernst Effect

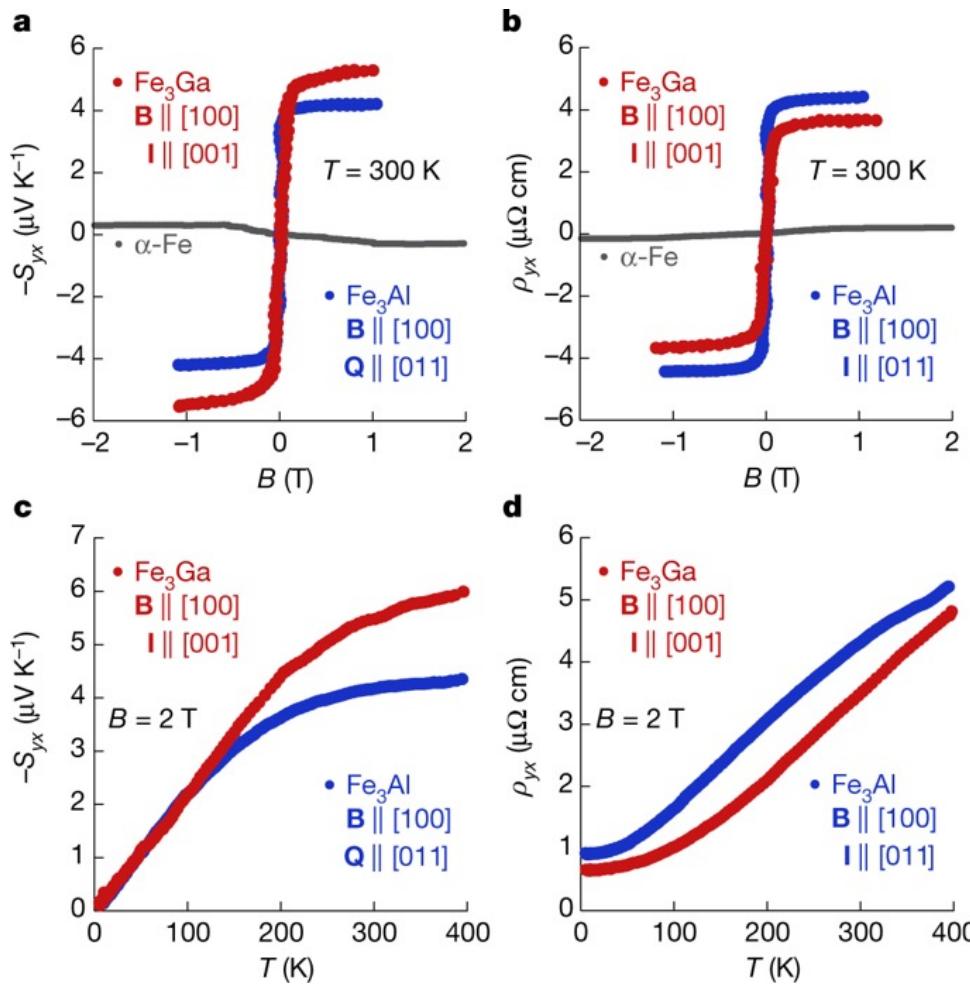


Co₂MnGa



Nernst Effect of iron compounds

Noky, J., Xu, Q., Felser, C. & Sun, Y. Large anomalous Hall and Nernst effects from nodal line symmetry breaking in Fe_2MnX ($X = \text{P}, \text{As}, \text{Sb}$). *Phys. Rev. B* 99, 165117 (2019).



Article

Iron-based binary ferromagnets for transverse thermoelectric conversion

<https://doi.org/10.1038/s41586-020-2230-z>

Received: 26 July 2019

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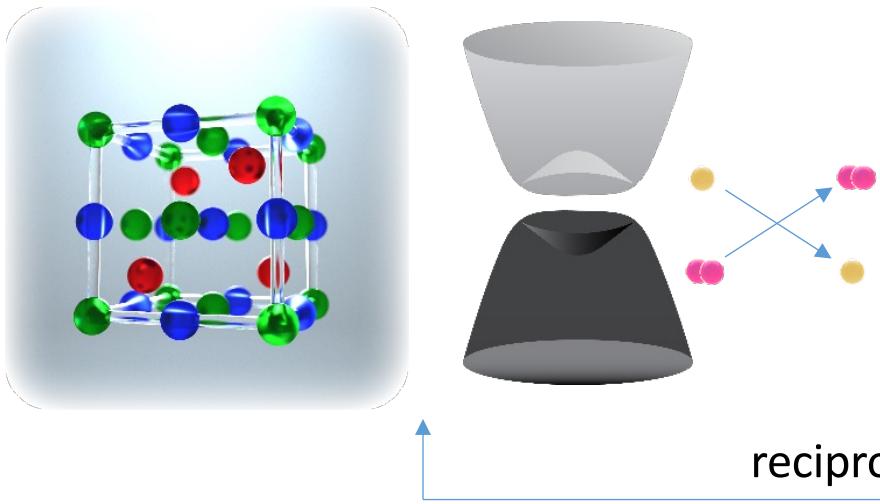
Check for updates

Akito Sakai^{1,2,3,10}, Susumu Minami^{4,5,10}, Takashi Koretsune^{6,10}, Taishi Chen^{1,3,10}, Tomoya Higo^{1,3,10}, Yangming Wang¹, Takuya Nomoto⁷, Motoaki Hirayama⁵, Shinji Miwa^{1,3,8}, Daisuke Nishio-Hamane¹, Fumiuki Ishii^{4,5}, Ryotaro Arita^{3,5,7} & Satoru Nakatsuji^{1,2,3,8,9,10}

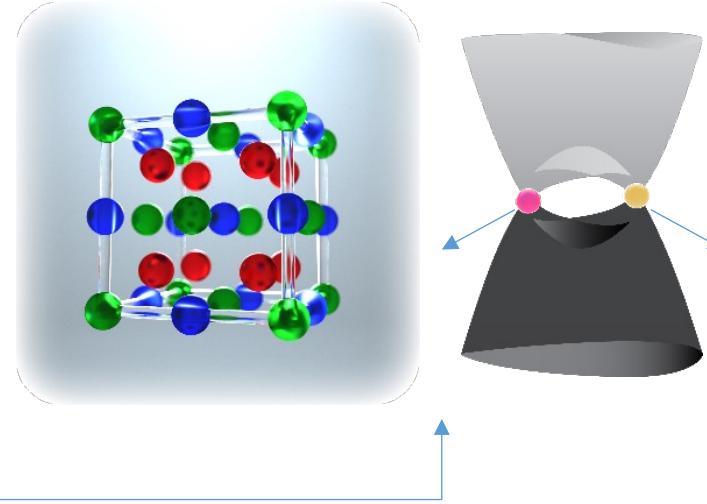
Thermoelectric generation using the anomalous Nernst effect (ANE) has great potential for application in energy harvesting technology because the transverse geometry of the Nernst effect should enable efficient, large-area and flexible coverage of a heat source. For such applications to be viable, substantial improvements will be necessary not only for their performance but also for the associated material costs,

Heusler compounds

A



B

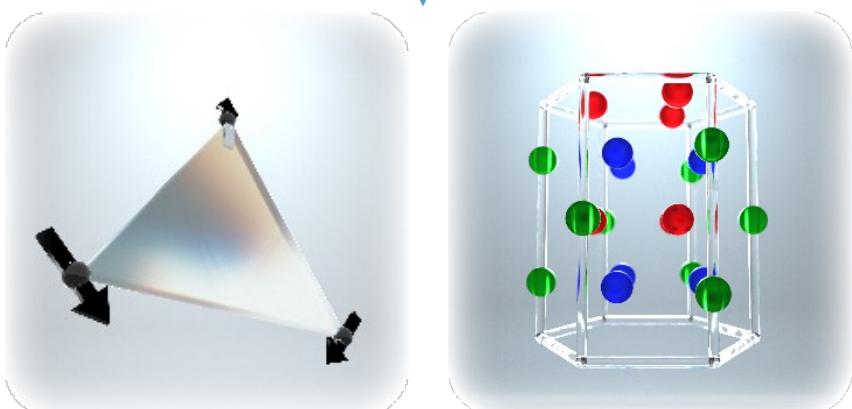


reciprocal space

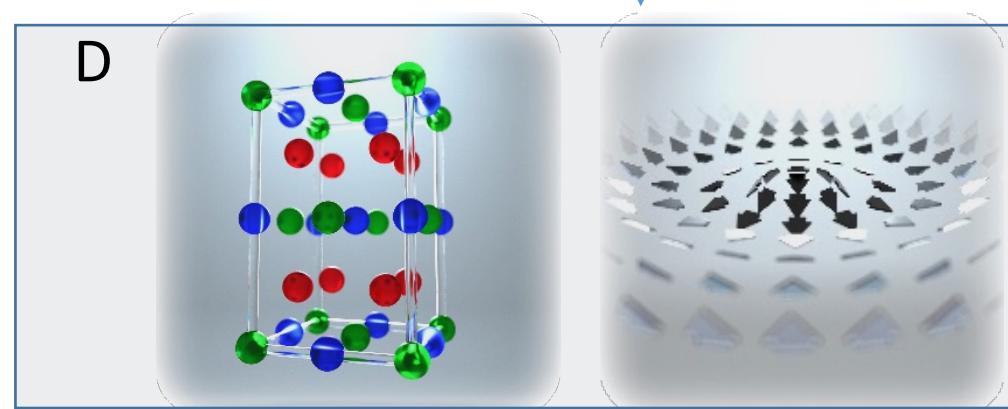
topological Heusler compounds

real space

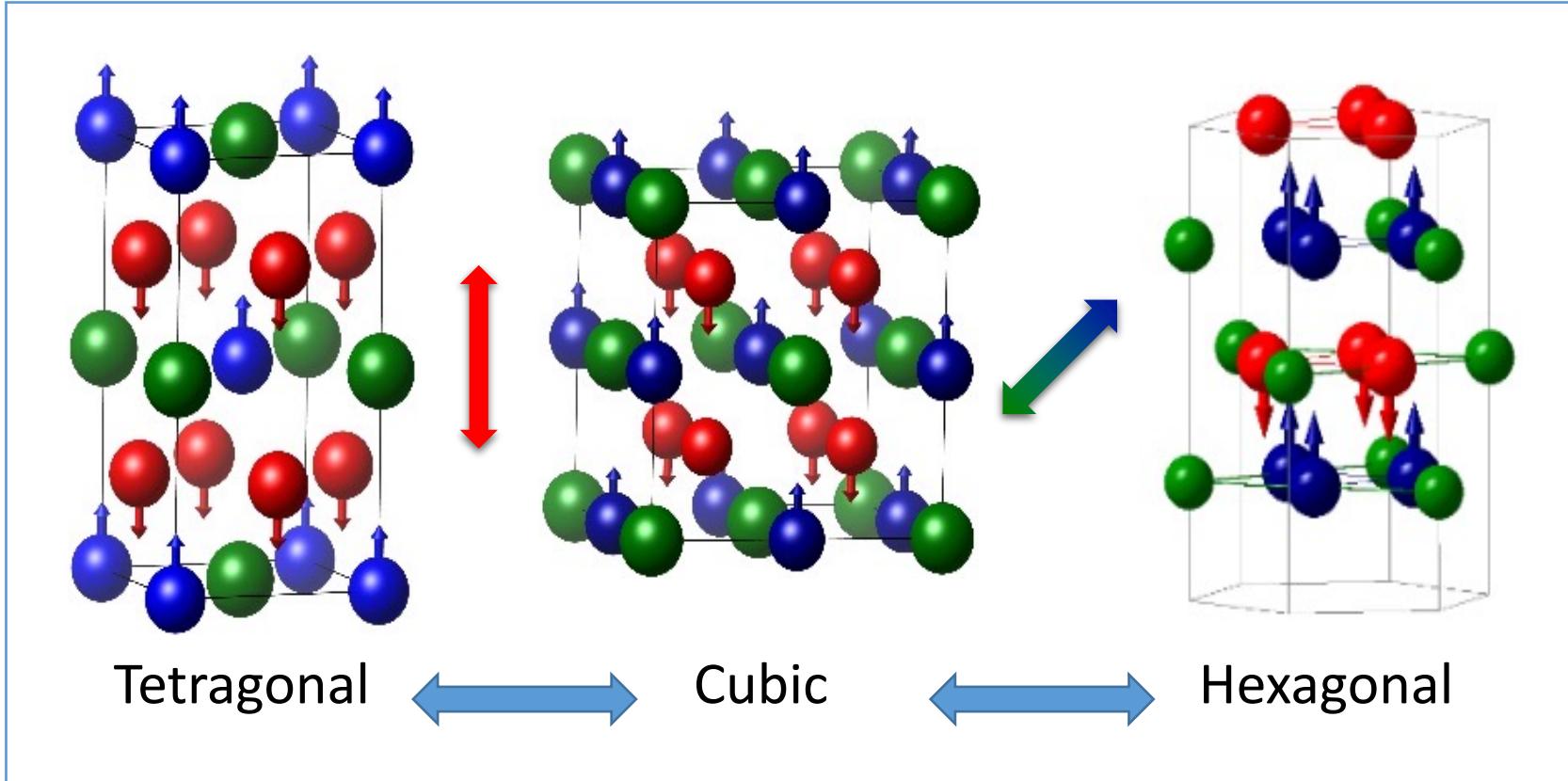
C



D



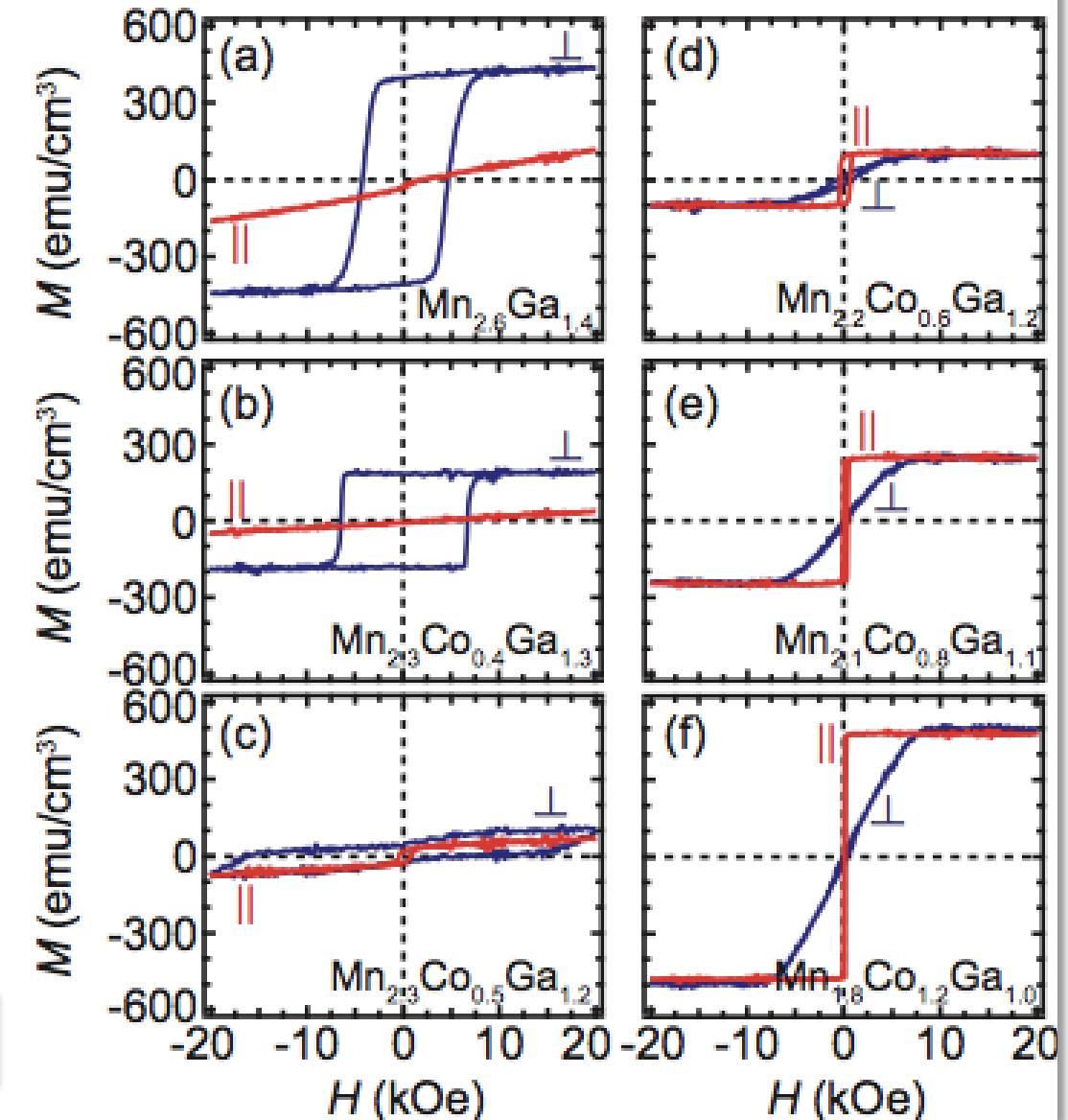
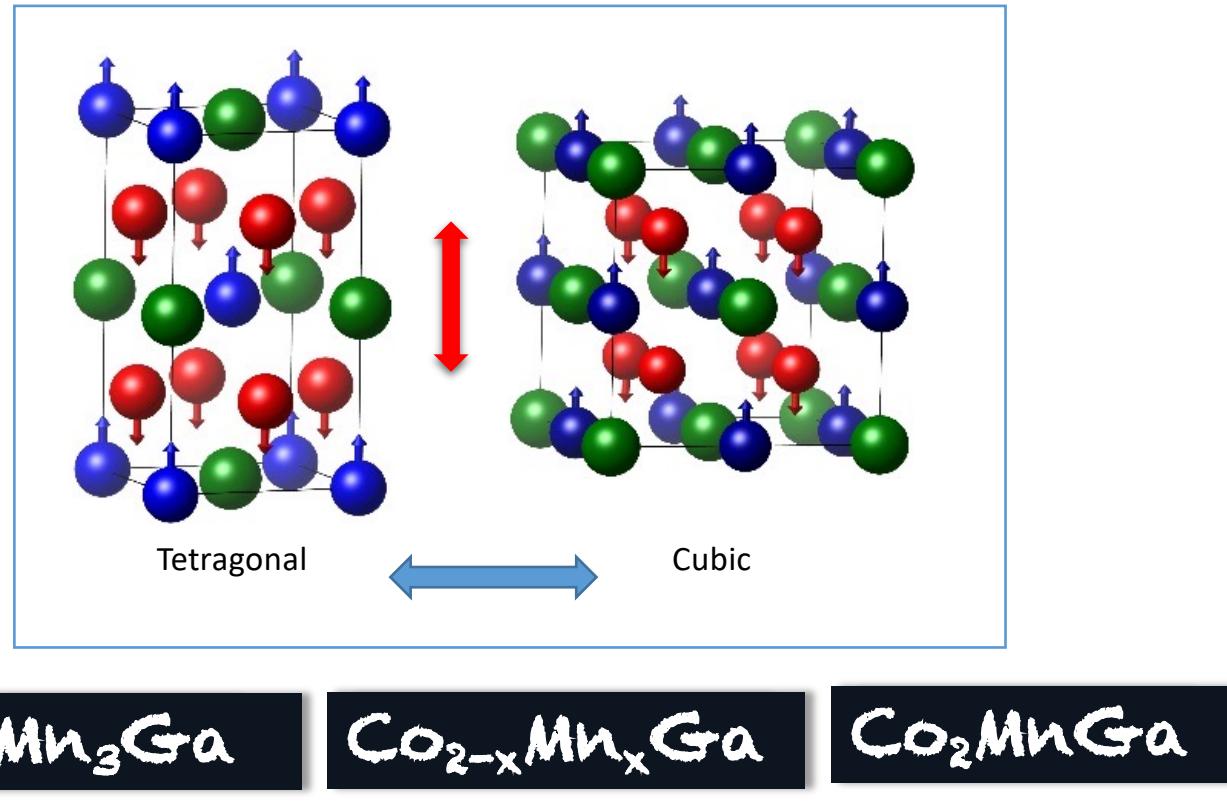
Heusler compounds with crystalline anisotropy



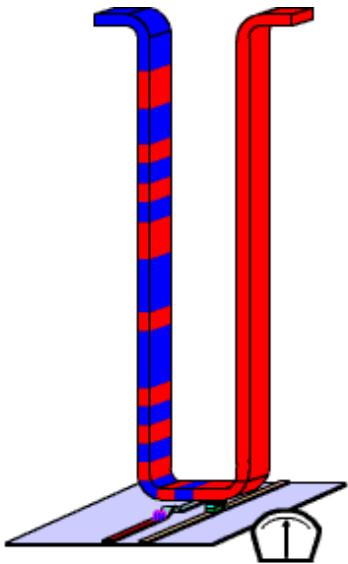
Mn_3Ga Mn_3Ge Mn_3Sn

Heusler compounds for STT MRAM

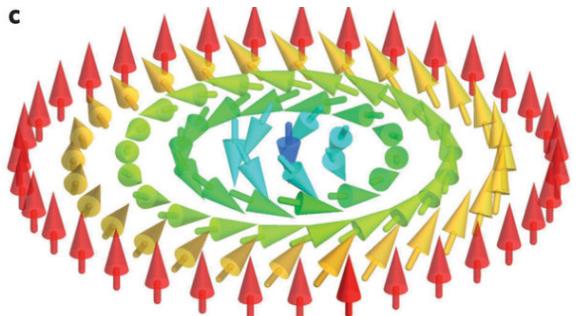
- Materials with low magnetic damping
- Materials with low magnetic moments
- Materials with high perpendicular anisotrop



Antiskyrmions in Heusler compounds

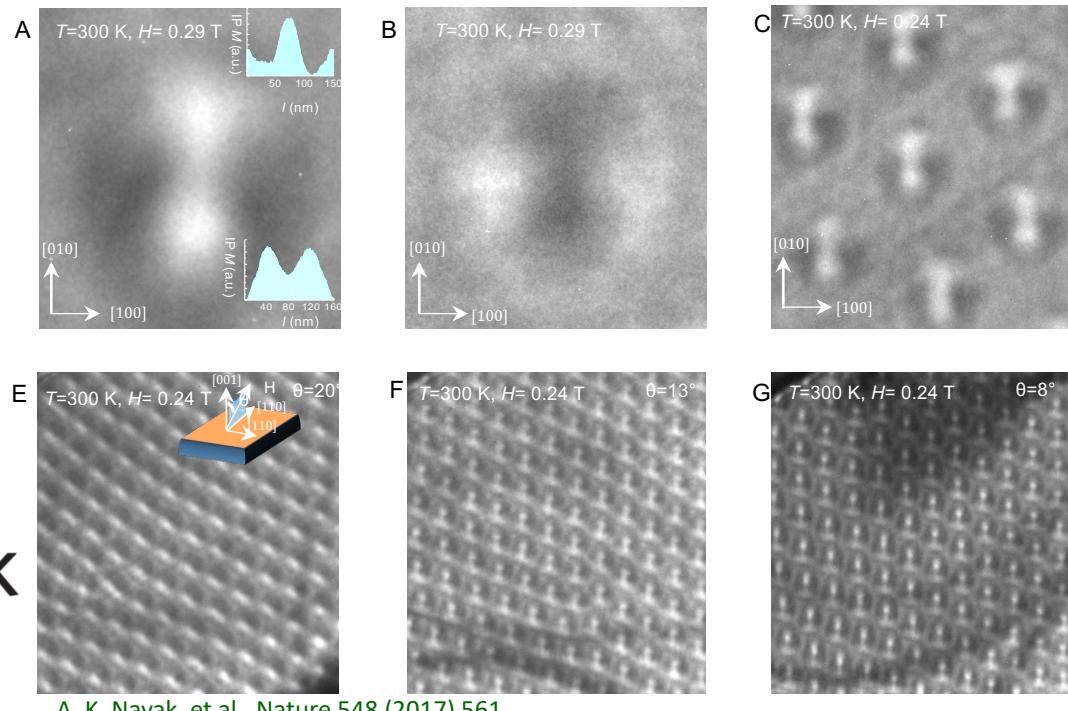
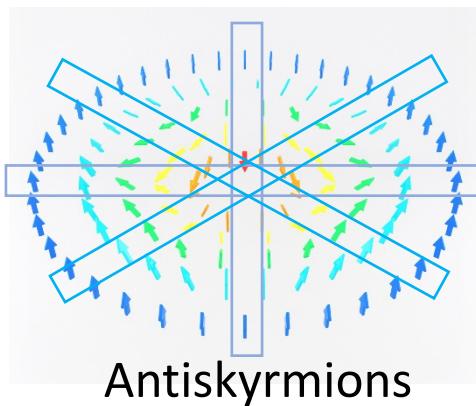


Stuart S. P. Parkin, et al.: *Magnetic Domain-Wall Racetrack Memory*, *Science* 320 (2008) 190–194

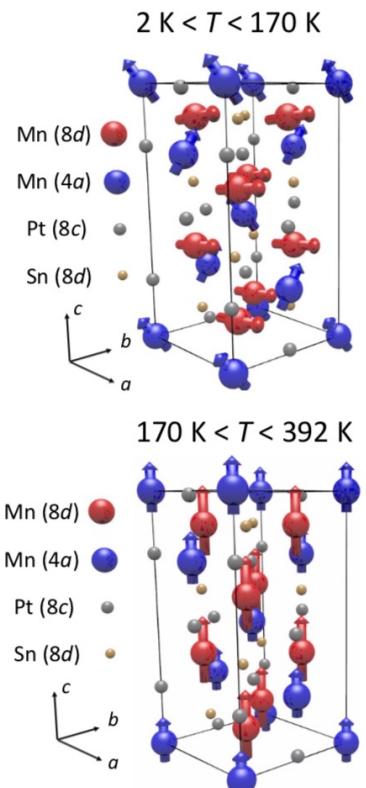
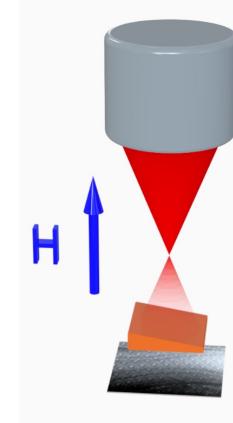


Skyrmions on the track

Albert Fert, Vincent Cros and João Sampaio

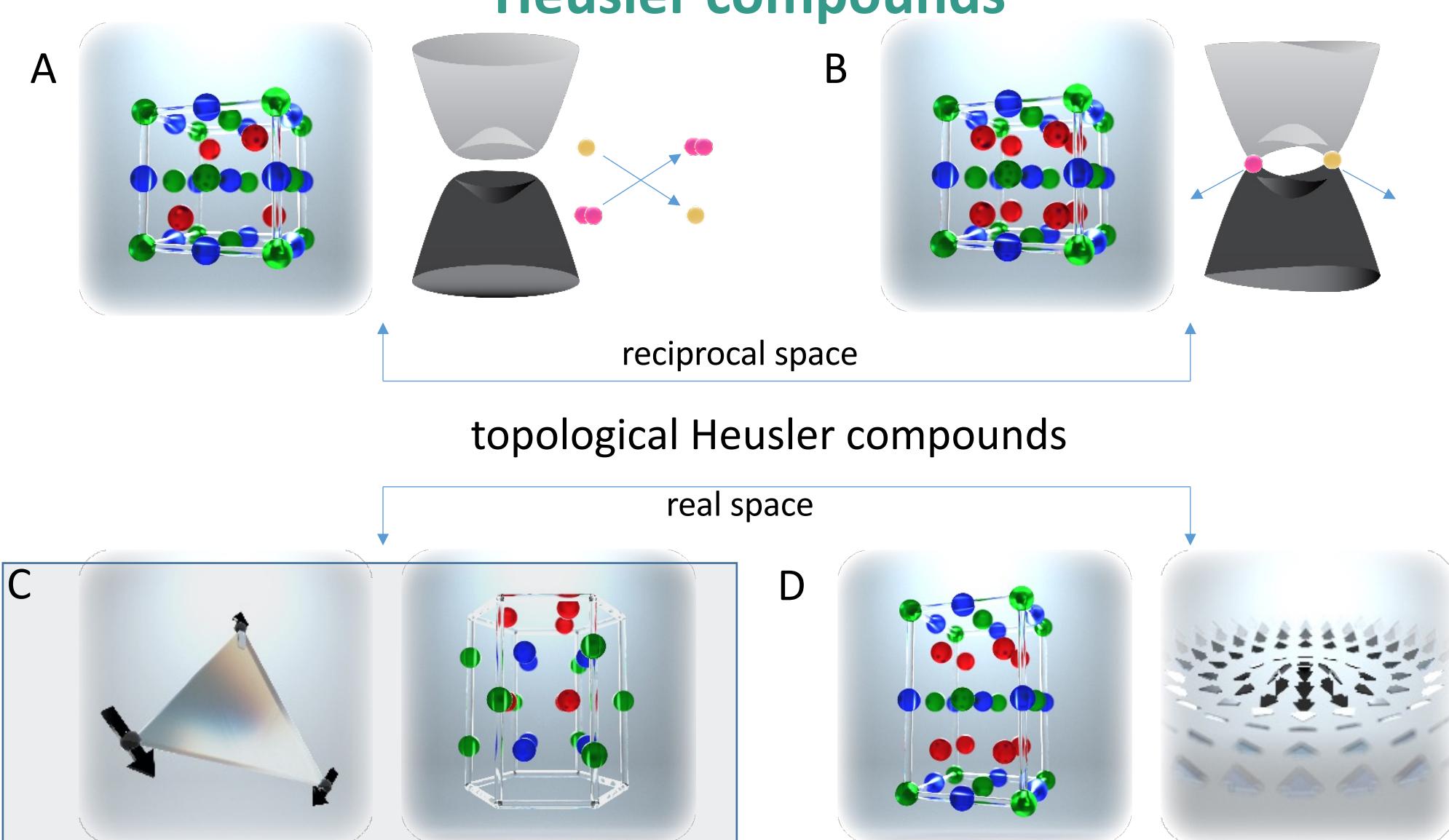


A. K. Nayak, et al., *Nature* 548 (2017) 561

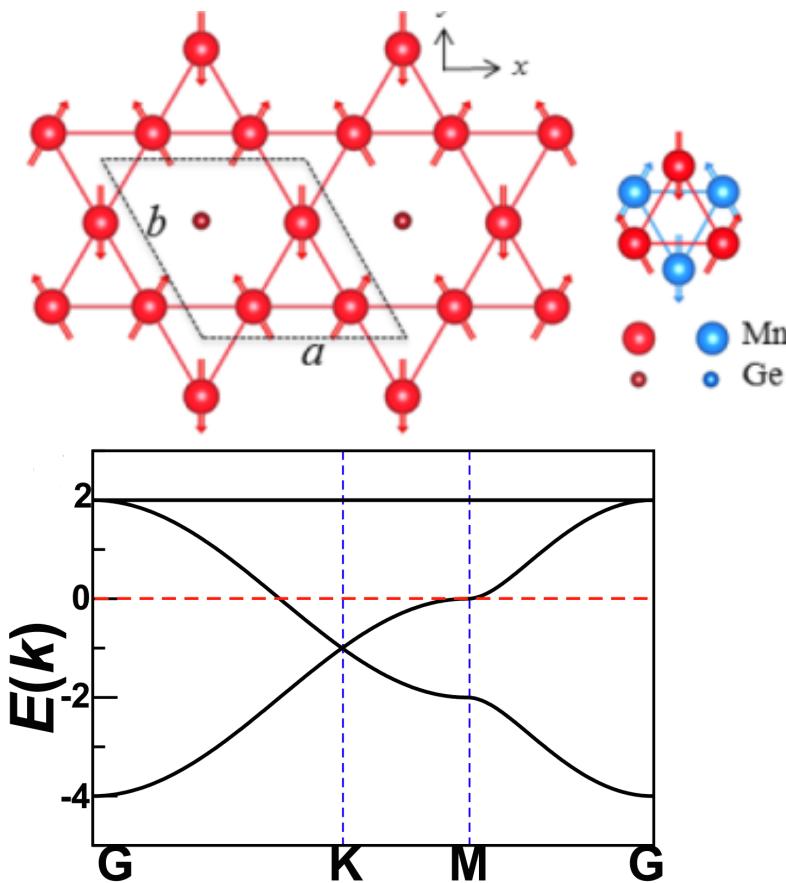


$\text{Mn}_{1.4}\text{PtSn}$

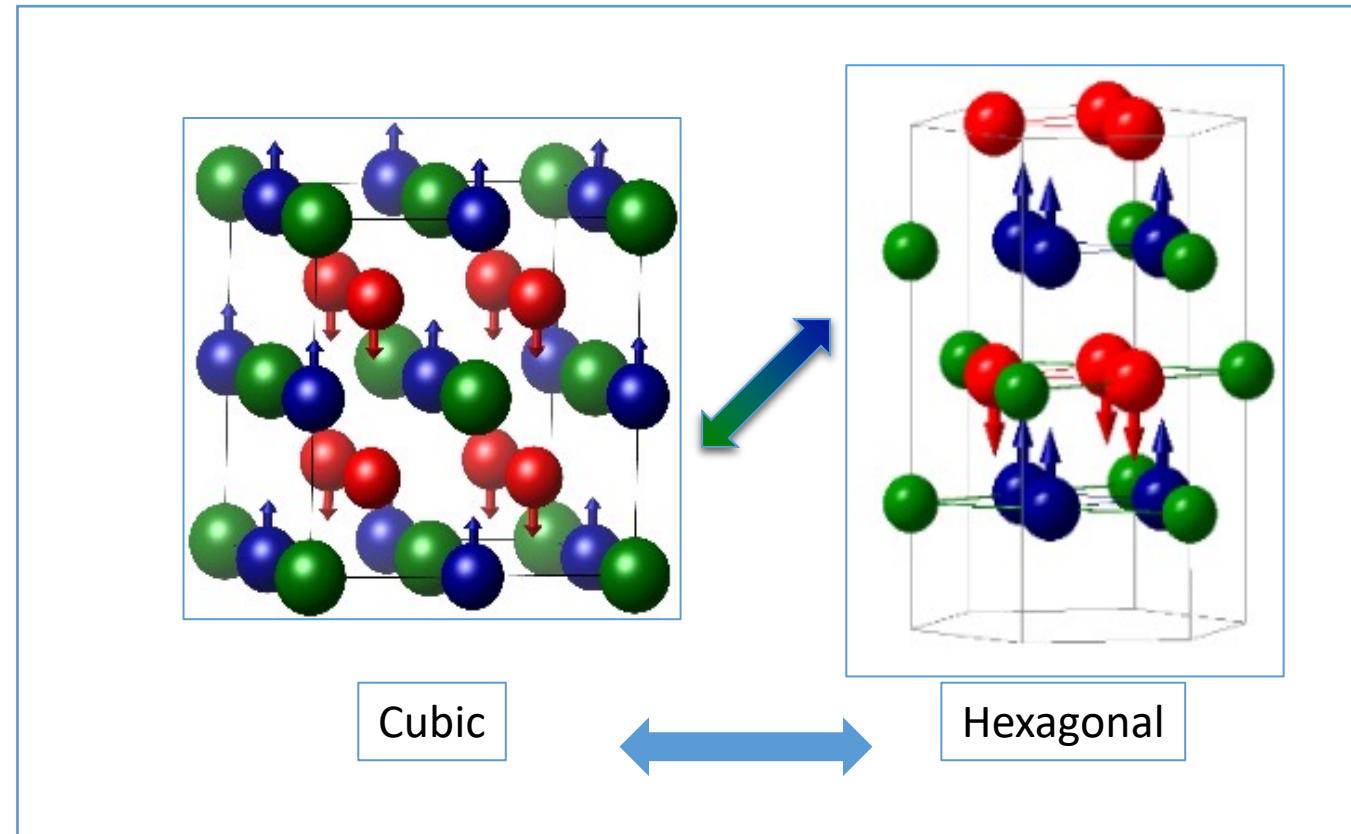
Heusler compounds



Kagome lattice



- Dirac cone,
- flat band, and
- van Hove singularity



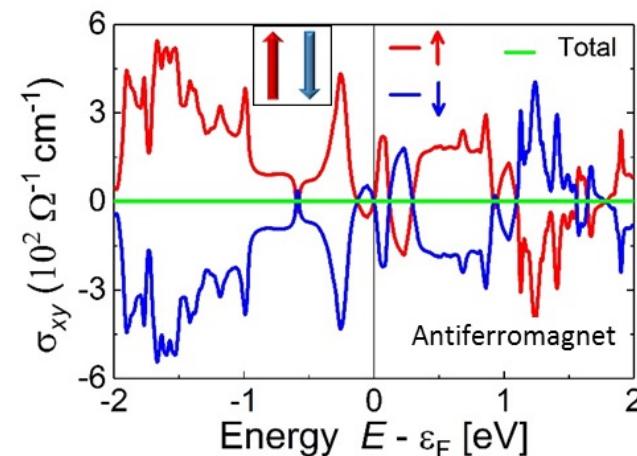
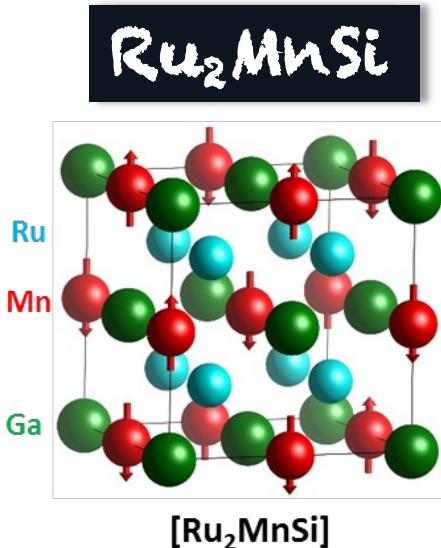
$\text{Mn}_3\text{Ga}, \text{Mn}_3\text{Ge}, \text{Mn}_3\text{Sn}$

antiferromagnetic topological materials

Table 3 | The magnetic topological materials identified in this work

Categories	Properties	Materials
I-A	Non-collinear manganese compounds	Mn ₃ GaC, Mn ₃ ZnC, Mn ₃ CuN, Mn ₃ Sn, Mn ₃ Ge, Mn ₃ Ir, Mn ₃ Pt, Mn ₅ Si ₃
I-B	Actinide intermetallic	UNiGa ₅ , UPtGa ₅ , NpRhGa ₅ , NpNiGa ₅
I-C	Rare-earth intermetallic	NdCo ₂ , TbCo ₂ , NpCo ₂ , PrAg DyCu, NdZn, TbMg, NdMg, Nd ₅ Si ₄ , Nd ₅ Ge ₄ , Ho ₂ RhIn ₈ , Er ₂ CoGa ₈ , Nd ₂ RhIn ₈ , Tm ₂ CoGa ₈ , Ho ₂ RhIn ₈ , DyCo ₂ Ga ₈ , TbCo ₂ Ga ₈ , Er ₂ Ni ₂ In, CeRu ₂ Al ₁₀ , Nd ₃ Ru ₄ Al ₁₂ , Pr ₃ Ru ₄ Al ₁₂ , ScMn ₆ Ge ₆ , YFe ₄ Ge ₄ , LuFe ₄ Ge ₄ , CeCoGe ₃
II-A	Metallic iron pnictides	LaFeAsO, CaFe ₂ As ₂ , EuFe ₂ As ₂ , BaFe ₂ As ₂ , Fe ₂ As, CaFe ₄ As ₃ , LaCrAsO, Cr ₂ As, CrAs, CrN
II-B	Semiconducting manganese pnictides	BaMn ₂ As ₂ BaMn ₂ Bi ₂ , CaMnBi ₂ , SrMnBi ₂ , CaMn ₂ Sb ₂ , CuMnAs, CuMnSb, Mn ₂ As
II-C	Rare-earth intermetallic compounds with the composition 1:2:2	PrNi ₂ Si ₂ , YbCo ₂ Si ₂ , DyCo ₂ Si ₂ , PrCo ₂ P ₂ , CeCo ₂ P ₂ , NdCo ₂ P ₂ , DyCu ₂ Si ₂ , CeRh ₂ Si ₂ , UAu ₂ Si ₂ , U ₂ Pd ₂ Sn, U ₂ Pd ₂ In, U ₂ Ni ₂ Sn, U ₂ Ni ₂ In, U ₂ Rh ₂ Sn
II-D	Rare-earth ternary compounds of the composition 1:1:1	CeMgPb, PrMgPb, NdMgPb, TmMgPb
III-A	Semiconducting actinides/rare-earth pnictides	HoP, UP, UP ₂ , UAs, NpS, NpSe, NpTe, NpSb, NpBi, U ₃ As ₄ , U ₃ P ₄
III-B	Metallic oxides	Ag ₂ NiO ₂ , AgNiO ₂ , Ca ₃ Ru ₂ O ₇ , Double perovskite Sr ₃ CoIrO ₆
III-C	Metal-to-insulator transition compounds	NiS ₂ , Sr ₂ Mn ₃ As ₂ O ₂
III-D	Semiconducting and insulating oxides, borates, hydroxides, silicates and phosphate	LuFeO ₃ , PdNiO ₃ , ErVO ₃ , DyVO ₃ , MnGeO ₃ , Tm ₂ Mn ₂ O ₇ , Yb ₂ Sn ₂ O ₇ , Tb ₂ Sn ₂ O ₇ , Ho ₂ Ru ₂ O ₇ , Er ₂ Ti ₂ O ₇ , Tb ₂ Ti ₂ O ₇ , Cd ₂ Os ₂ O ₇ , Ho ₂ Ru ₂ O ₇ , Cr ₂ ReO ₆ , NiCr ₂ O ₄ , MnV ₂ O ₄ , Co ₂ SiO ₄ , Fe ₂ SiO ₄ , PrFe ₃ (BO ₃) ₄ , KCo ₄ (PO ₄) ₃ , CoPS ₃ , SrMn(VO ₄)(OH), Ba ₅ Co ₅ ClO ₁₃ , FeI ₂

antiferromagnetic materials



$$\sigma_{xy}^A(\mu) = ie^2 \left(\frac{1}{2\pi}\right)^3 \int_k dk \sum_{E(n,k) < \mu} f(n, k, \mu) \Omega_{n,xy}(k),$$

The anomalous Hall conductivity in an antiferromagnetic metal is zero

Manna et al., Phys. Rev. X 8 (2018) 041045, arXiv:1712.10174

Manna et al., Nature Review Materials, 3 (2018) 244 arXiv:1802.02838v1

REVIEWS OF MODERN PHYSICS, VOLUME 82, APRIL–JUNE 2010

Anomalous Hall effect

Naoto Nagaosa

Department of Applied Physics, University of Tokyo, Tokyo 113-8656, Japan
and Cross-Correlated Materials Research Group (CMRG), and Correlated Electron Research Group (CERG), ASI, RIKEN, Wako, 351-0198 Saitama, Japan

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Department of Physics, Texas A&M University, College Station, Texas 77843-4242, USA
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Shigeki Onoda

Condensed Matter Theory Laboratory, ASI, RIKEN, Wako, 351-0198 Saitama, Japan

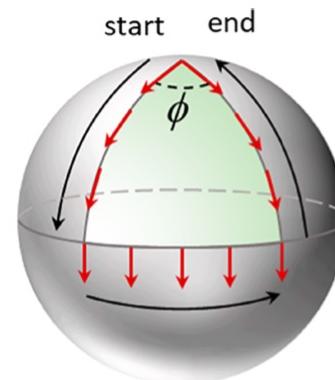
A. H. MacDonald

Department of Physics, University of Texas at Austin, Austin, Texas 78712-1081, USA

N. P. Ong

Department of Physics, Princeton University, Princeton, New Jersey 08544, USA

(Published 13 May 2010)



Berry curvature

antiferromagnetic topological materials

Table 3 | The magnetic topological materials identified in this work

Categories	Properties	Materials	
I-A	Non-collinear manganese compounds	Mn_3GaC , Mn_3ZnC , Mn_3CuN , Mn_3Sn , Mn_3Ge , Mn_3Ir , Mn_3Pt , Mn_5Si_3	
	<p>Mn₃Ge, Mn₃Sn</p> <p>c</p> <p>High</p> <p>0</p> <p>AHC (σ_{II}^k) (arbitrary units)</p> <p>D</p>	<p>A</p> <p>B</p> <p>2 K 100 K 200 K 300 K</p> <p>$I \parallel [0001]$ $H \parallel [01-10]$</p> <p>σ_{xz}</p> <p>C</p> <p>$I \parallel [01-10]$ $H \parallel [2-1-10]$</p> <p>ρ_H ($\mu\text{ohm}\cdot\text{cm}$)</p> <p>$\sigma_H$ ($\text{ohm}^{-1}\cdot\text{cm}^{-1}$)</p> <p>D</p> <p>σ_{yz}</p> <p>E</p> <p>$I \parallel [2-1-10]$ $H \parallel [0001]$</p> <p>σ_{xy}</p> <p>F</p> <p>$B \parallel [0110], I \parallel [0001]$ $B \parallel [2110], I \parallel [0110]$ $B \parallel [0001], I \parallel [0110]$</p> <p>$\rho_H^{\text{AF}}$ ($\mu\Omega\cdot\text{cm}$)</p>	<p>d</p> <p>$B \parallel [0110]$ $I \parallel [0001]$</p> <p>$B \parallel [2110]$ $I \parallel [0110]$</p> <p>$B \parallel [0001]$ $I \parallel [0110]$</p> <p>ρ_H ($\Omega^{-1}\cdot\text{cm}^{-1}$)</p> <p>$B$ (T)</p> <p>f</p> <p>$B \parallel [0110], I \parallel [0001]$ $B \parallel [2110], I \parallel [0110]$ $B \parallel [0001], I \parallel [0110]$</p> <p>$\rho_H^{\text{AF}}$ ($\mu\Omega\cdot\text{cm}$)</p> <p>B (T)</p>

Giant anomalous Nernst effect

PRL 119, 056601 (2017)

PHYSICAL REVIEW LETTERS

week ending
4 AUGUST 2017

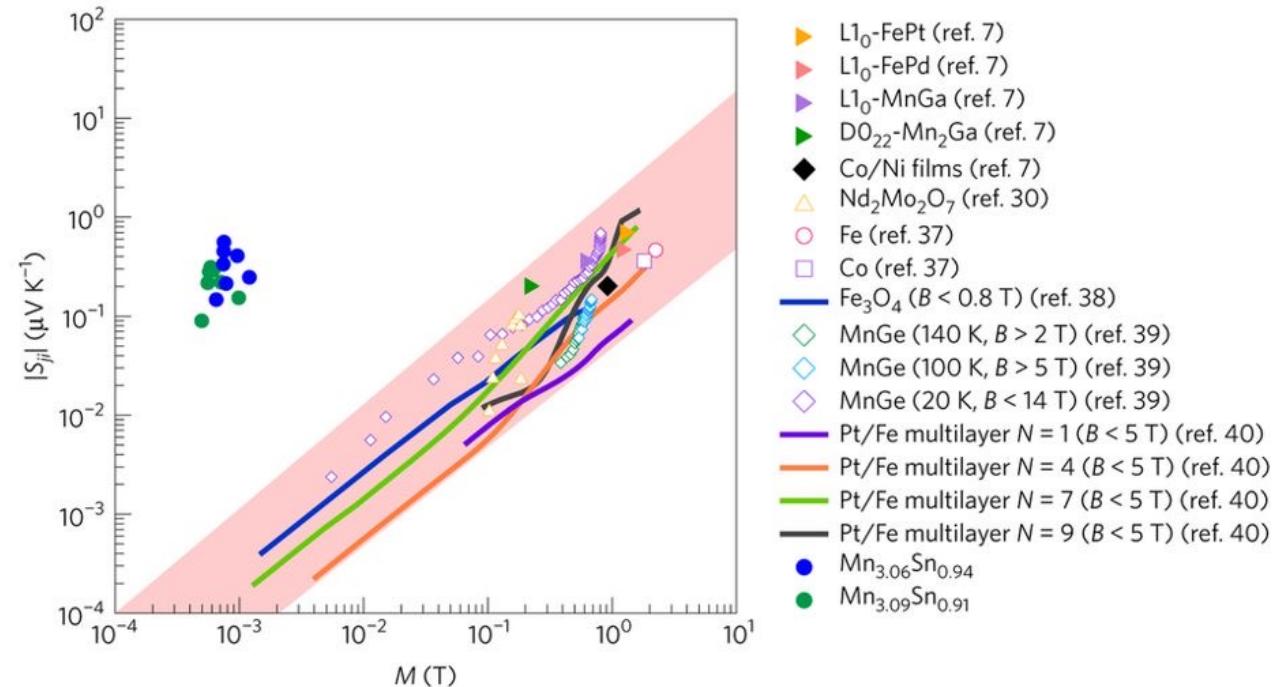
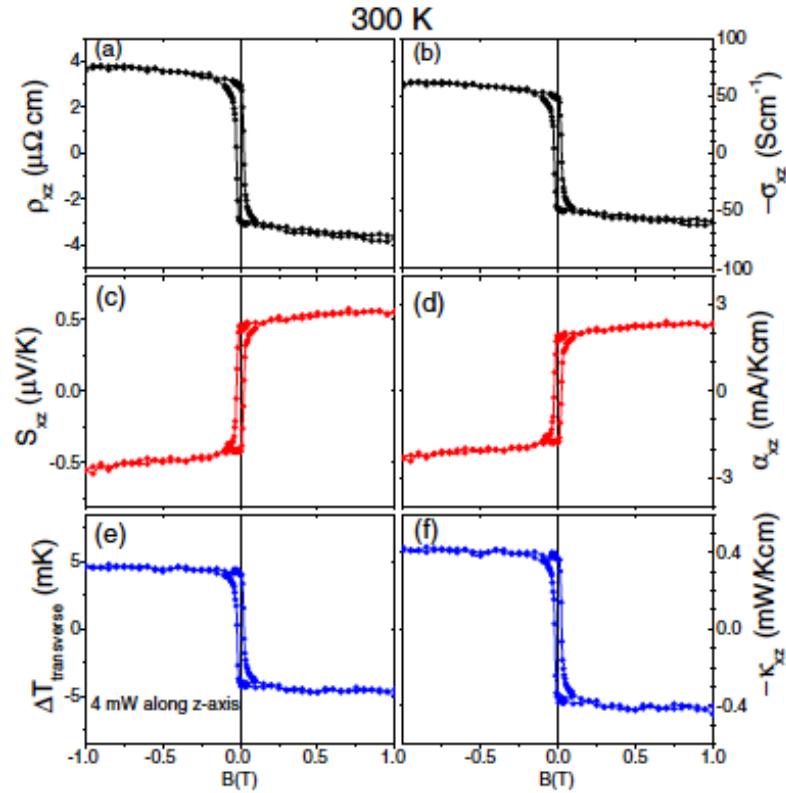
nature
physics

LETTERS

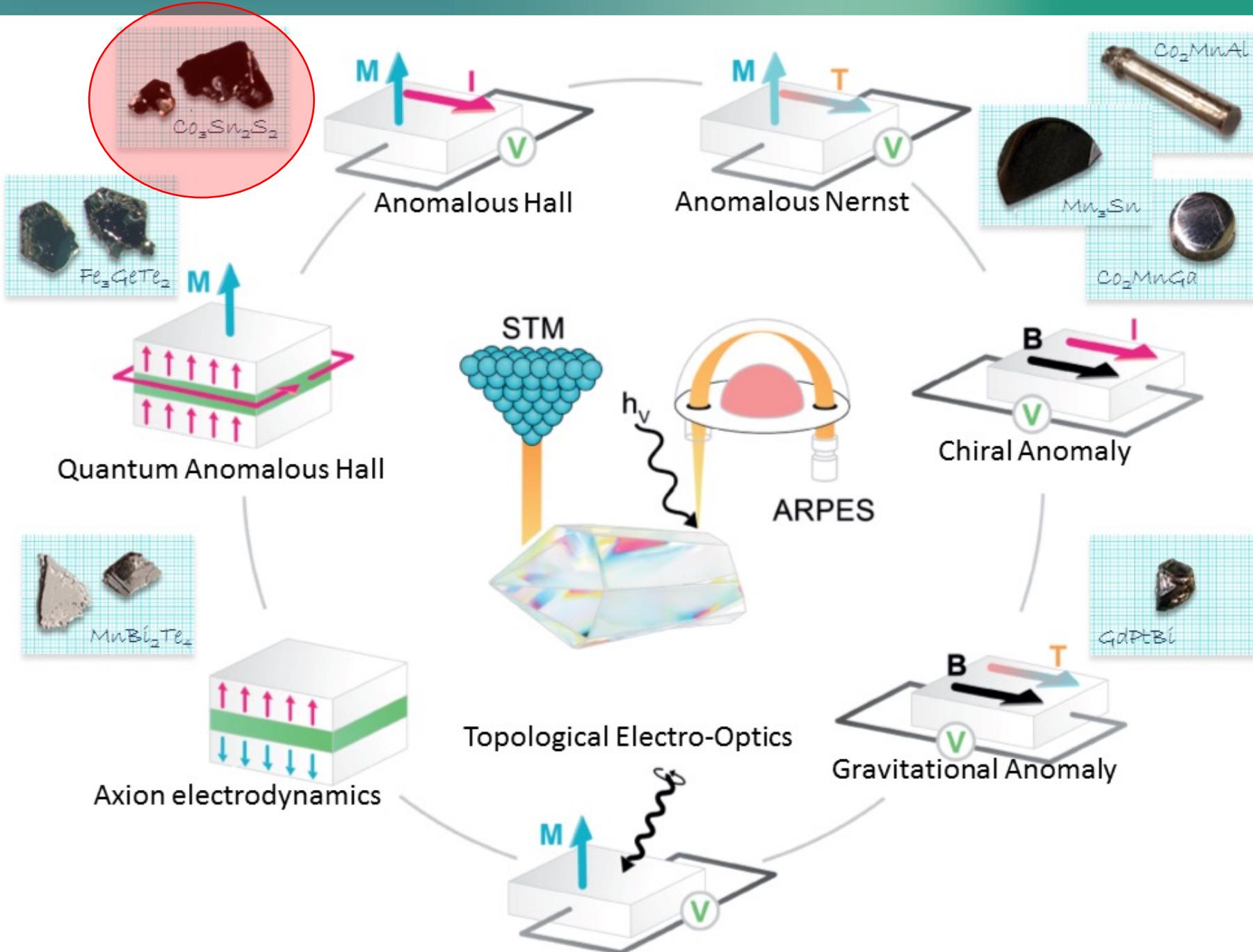
PUBLISHED ONLINE: 24 JULY 2017 | DOI: 10.1038/NPHYS4181

Anomalous Nernst and Righi-Leduc Effects in Mn_3Sn : Berry Curvature and Entropy Flow

Xiaokang Li,¹ Liangcai Xu,¹ Linchao Ding,¹ Jinhua Wang,¹ Mingsong Shen,¹
Xiufang Lu,¹ Zengwei Zhu,^{1,*} and Kamran Behnia^{1,2,†}



Magnetization dependence of the spontaneous Nernst effect for ferromagnetic metals and Mn_3Sn



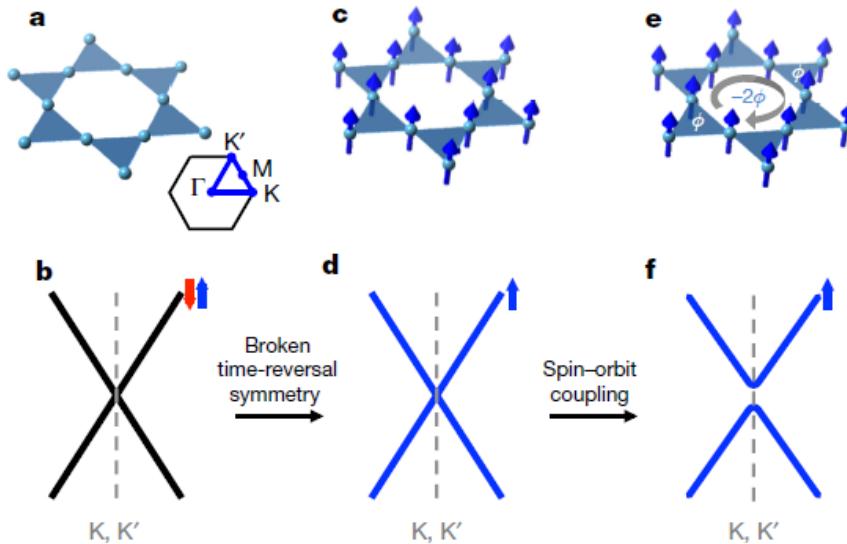
Kagome lattice

Fe₂Sn₃

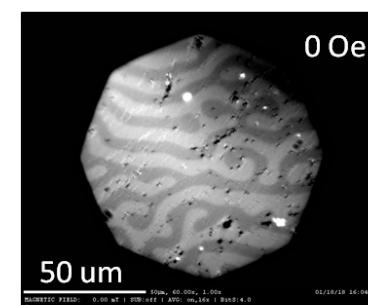
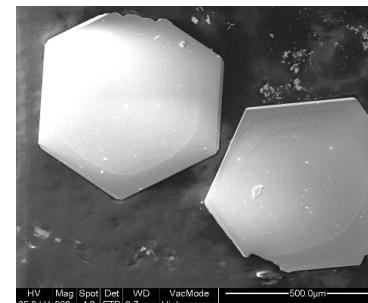
LETTER

doi:10.1038/nature25987

Massive Dirac fermions in a ferromagnetic kagome metal

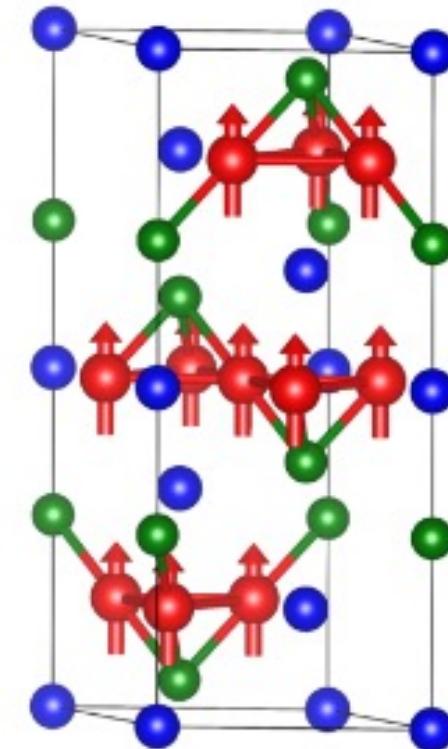


Nature, 2018, doi:10.1038/nature25987

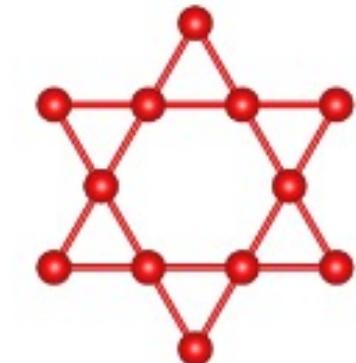


Co₃Sn₂S₂

Looking for Weyl fermions on a ferromagnetic Kagomé lattice with out of plane magnetisation.



● Co
● Sn
● S

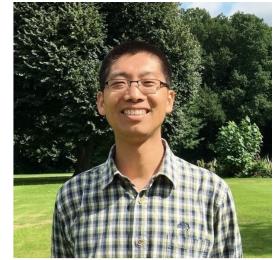


Enke Liu, et al. Nature Physics 14 (2018) 1125 , preprint arXiv:1712.06722

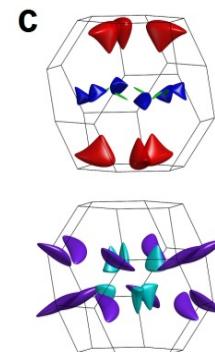
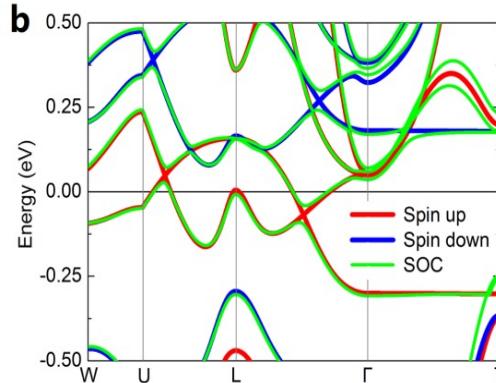


Kagome lattice

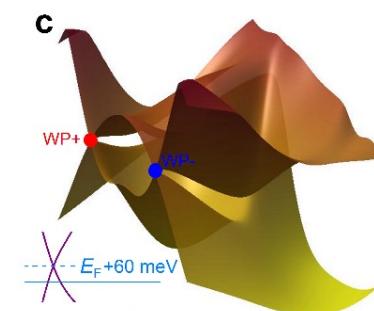
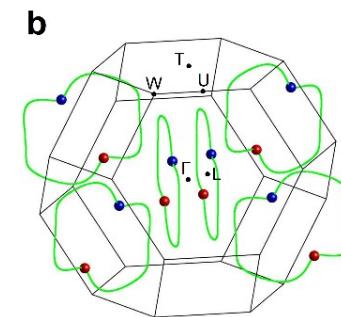
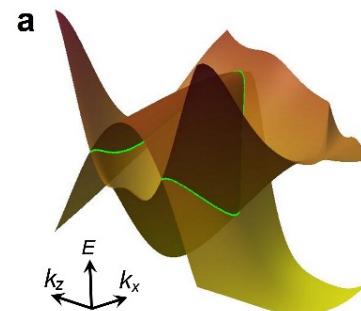
$\text{Co}_3\text{Sn}_2\text{S}_2$



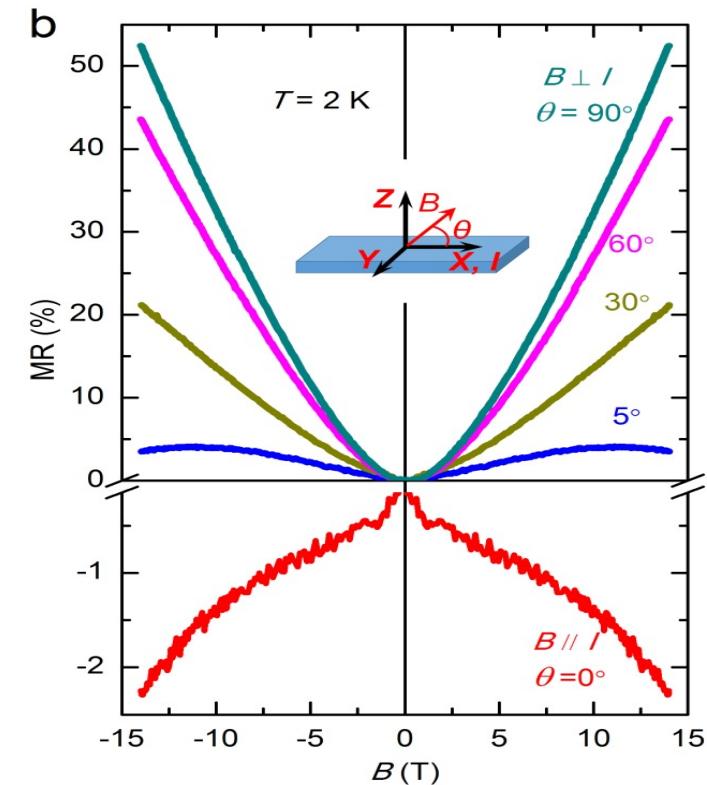
half metallic



$\text{Co}_3\text{Sn}_2\text{S}_2$



chiral anomaly



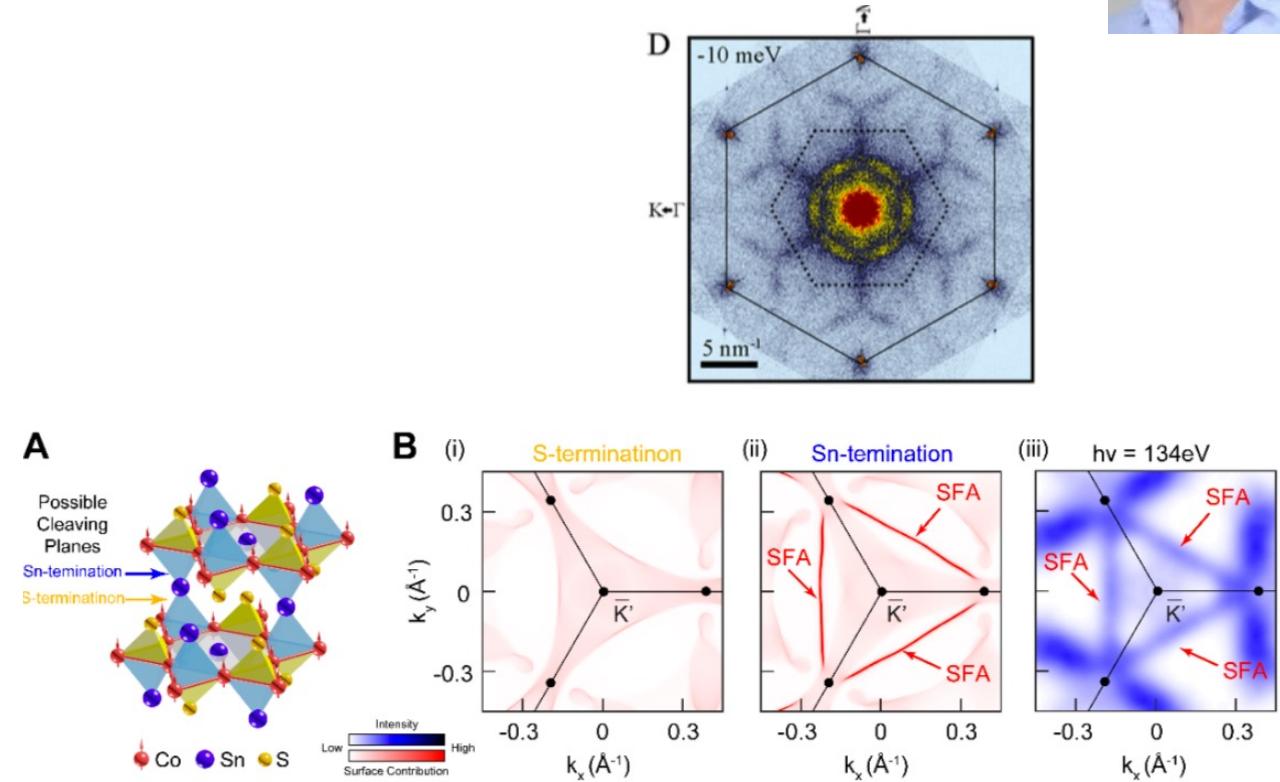
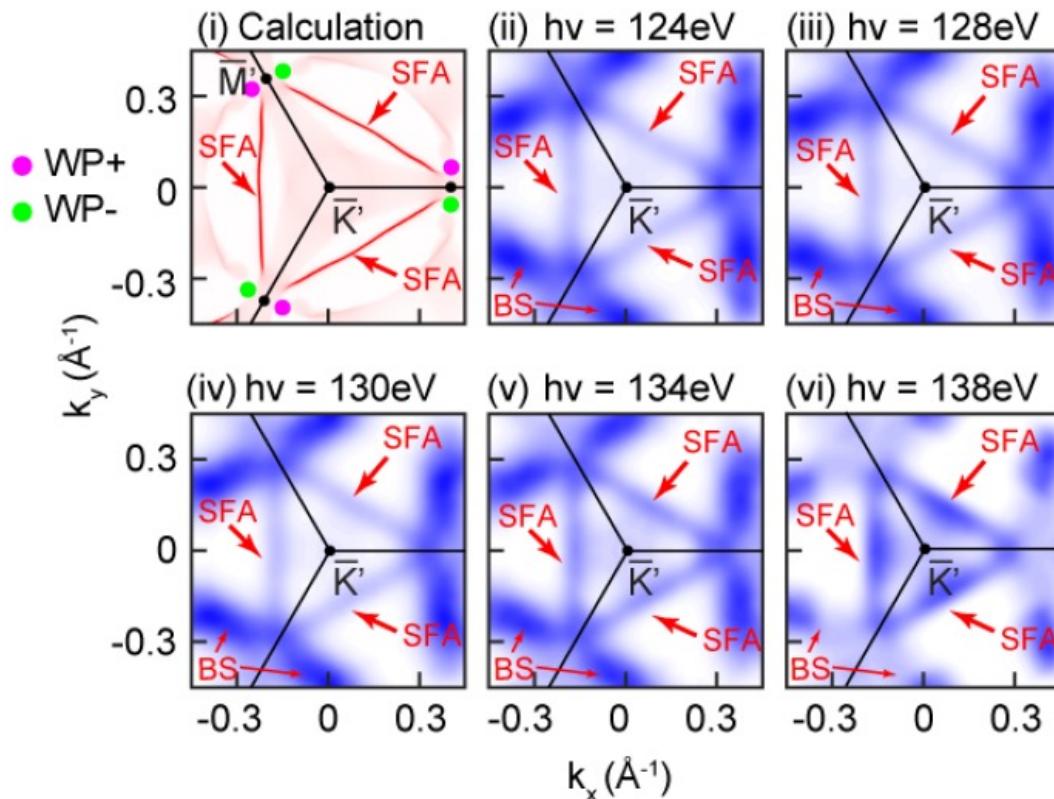
Kagome lattice



$\text{Co}_3\text{Sn}_2\text{S}_2$



STM and ARPES confirms Weyl and Fermiarcs

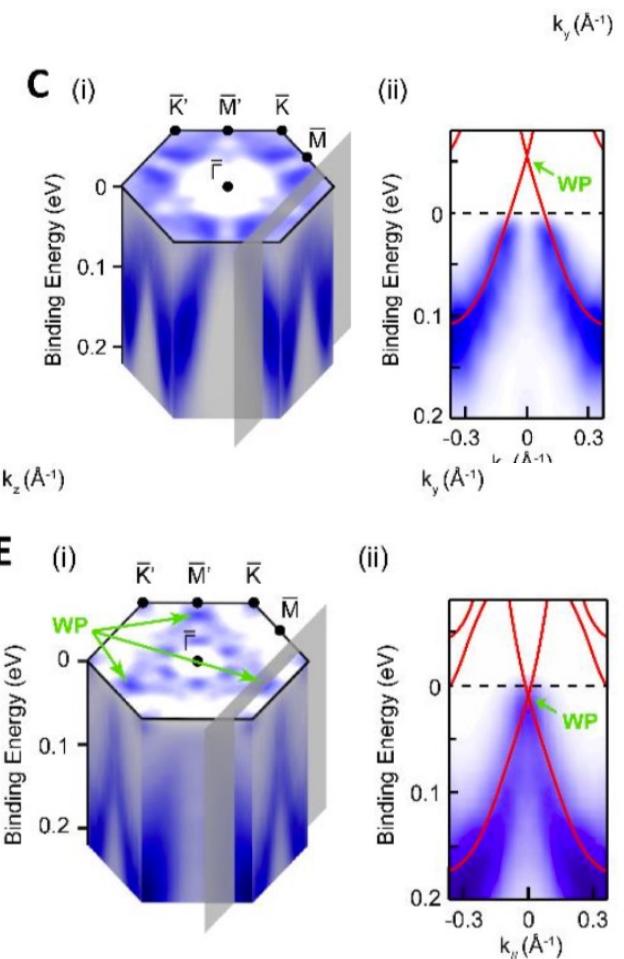
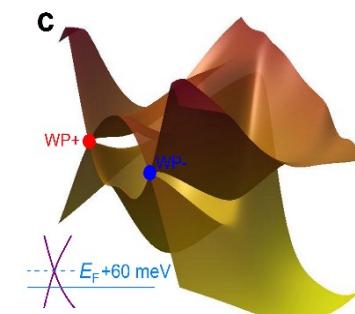
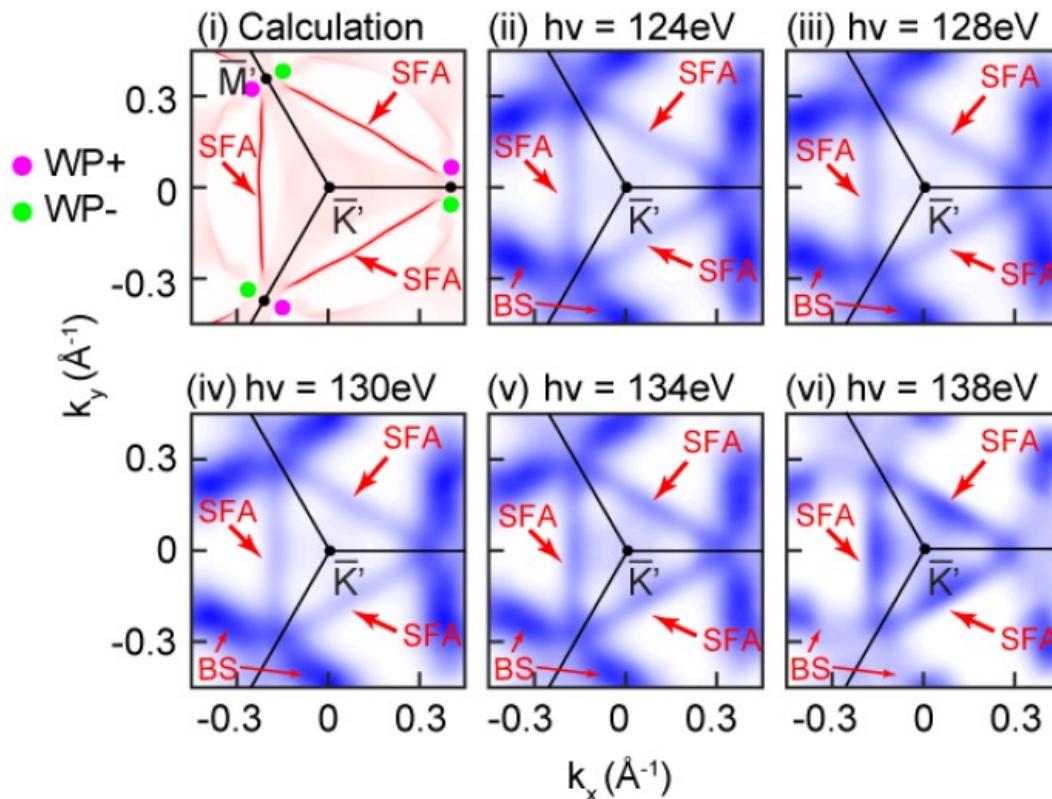


Kagome lattice

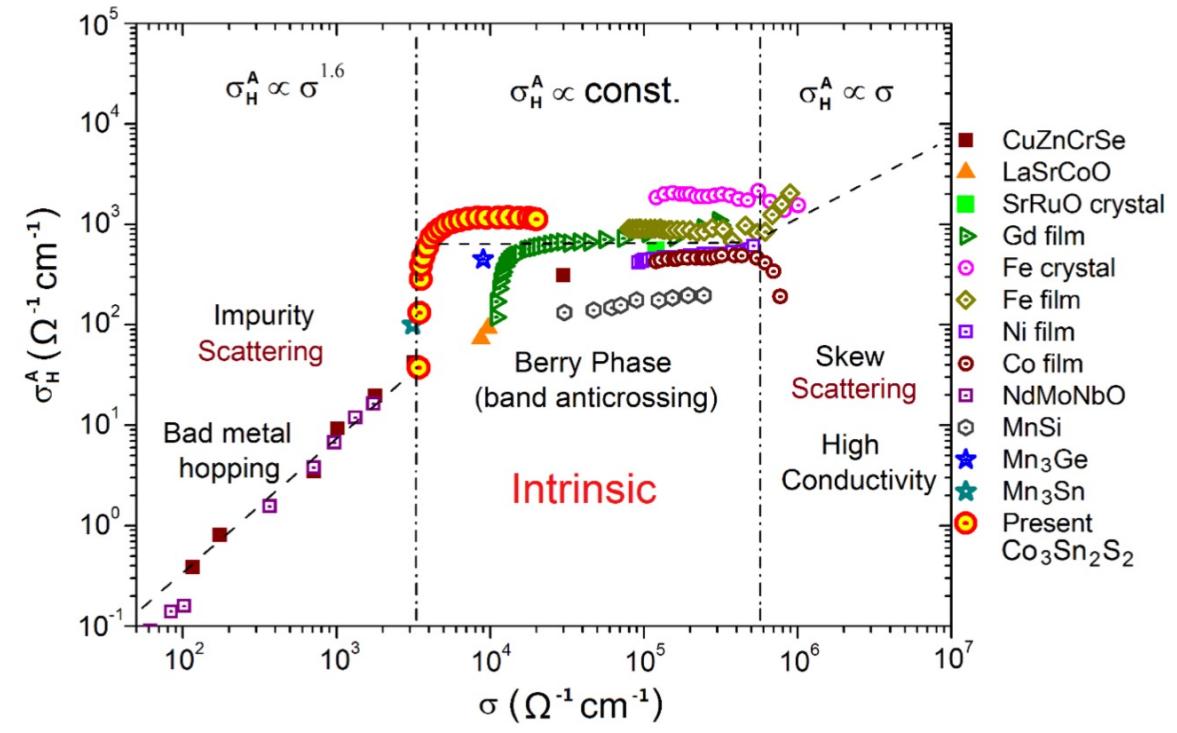
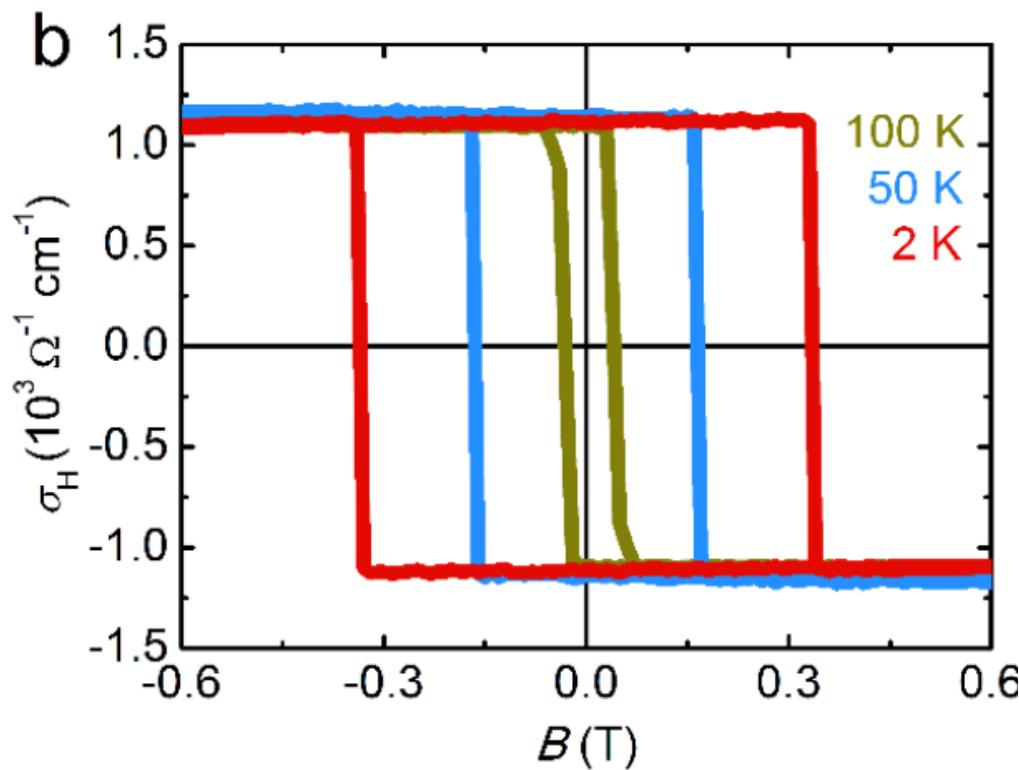
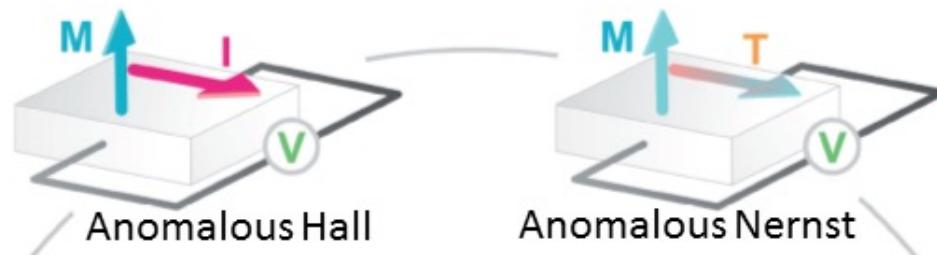


$\text{Co}_3\text{Sn}_2\text{S}_2$

ARPES confirms Weyl and Fermiarcs

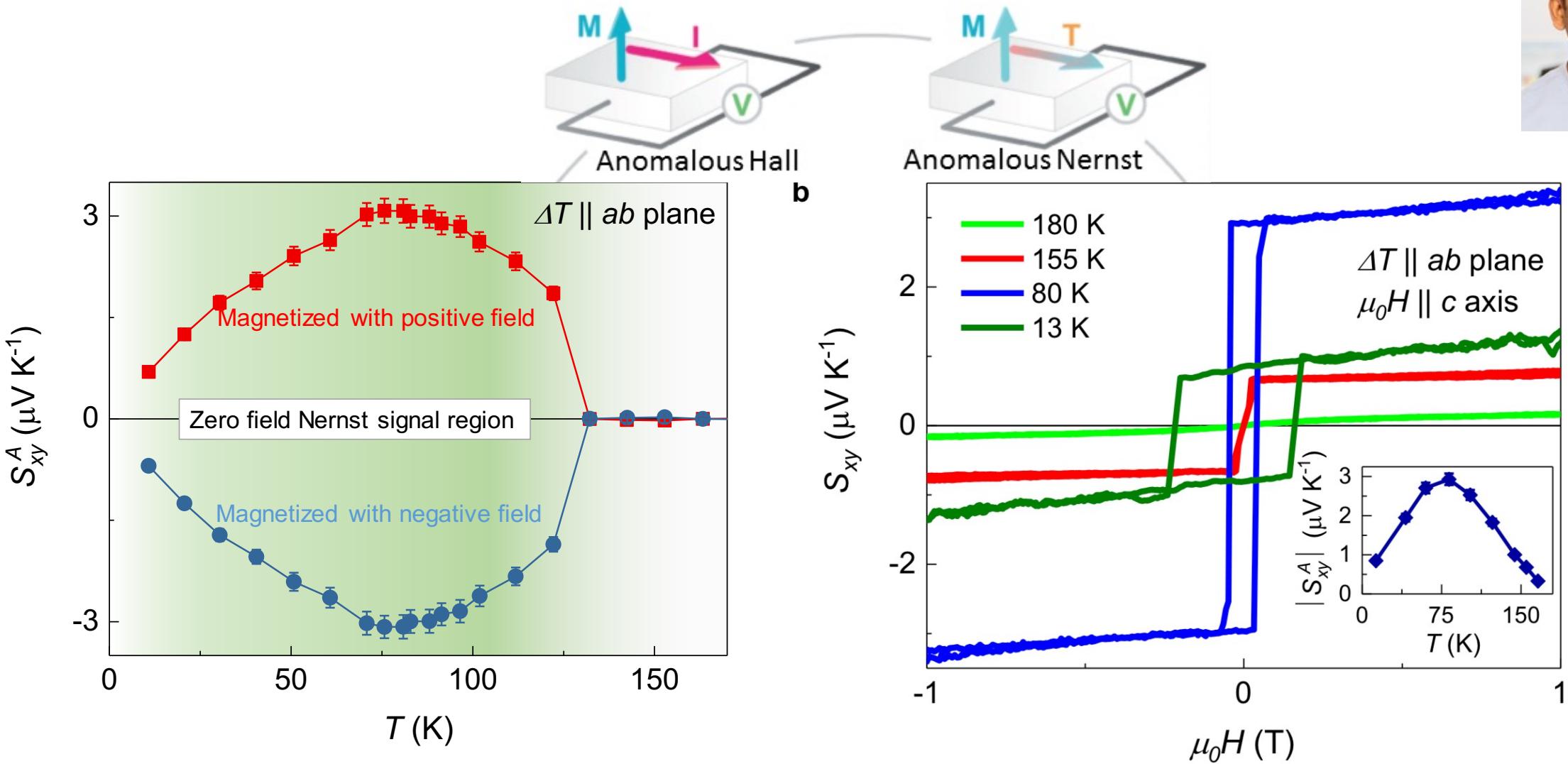
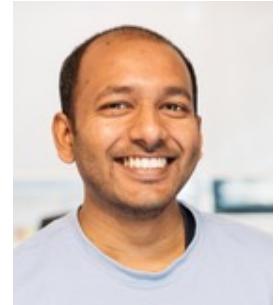


Giant anomalous Hall effect



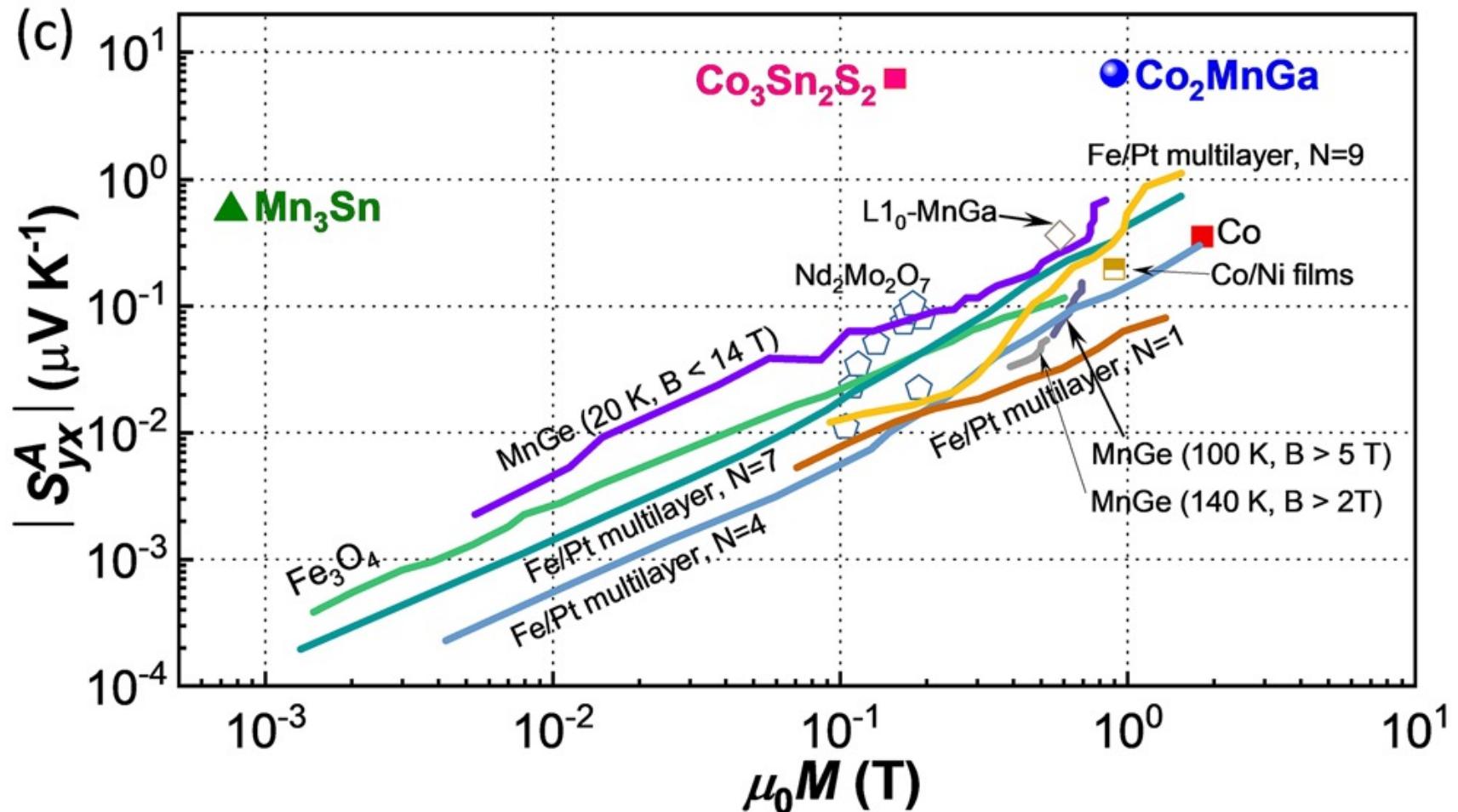
Naoto Nagaosa and Yoshinori Tokura 2012 Phys. Scr. 2012 014020

Giant anomalous Nernst effect

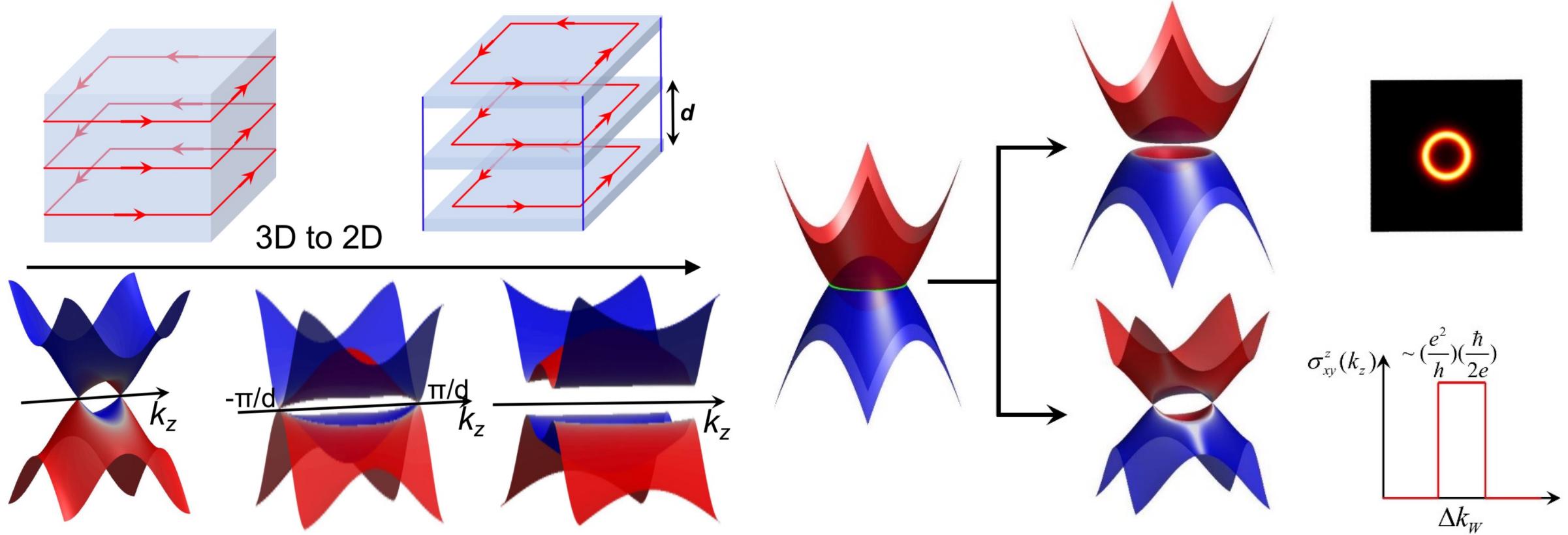


Giant anomalous Nernst effect

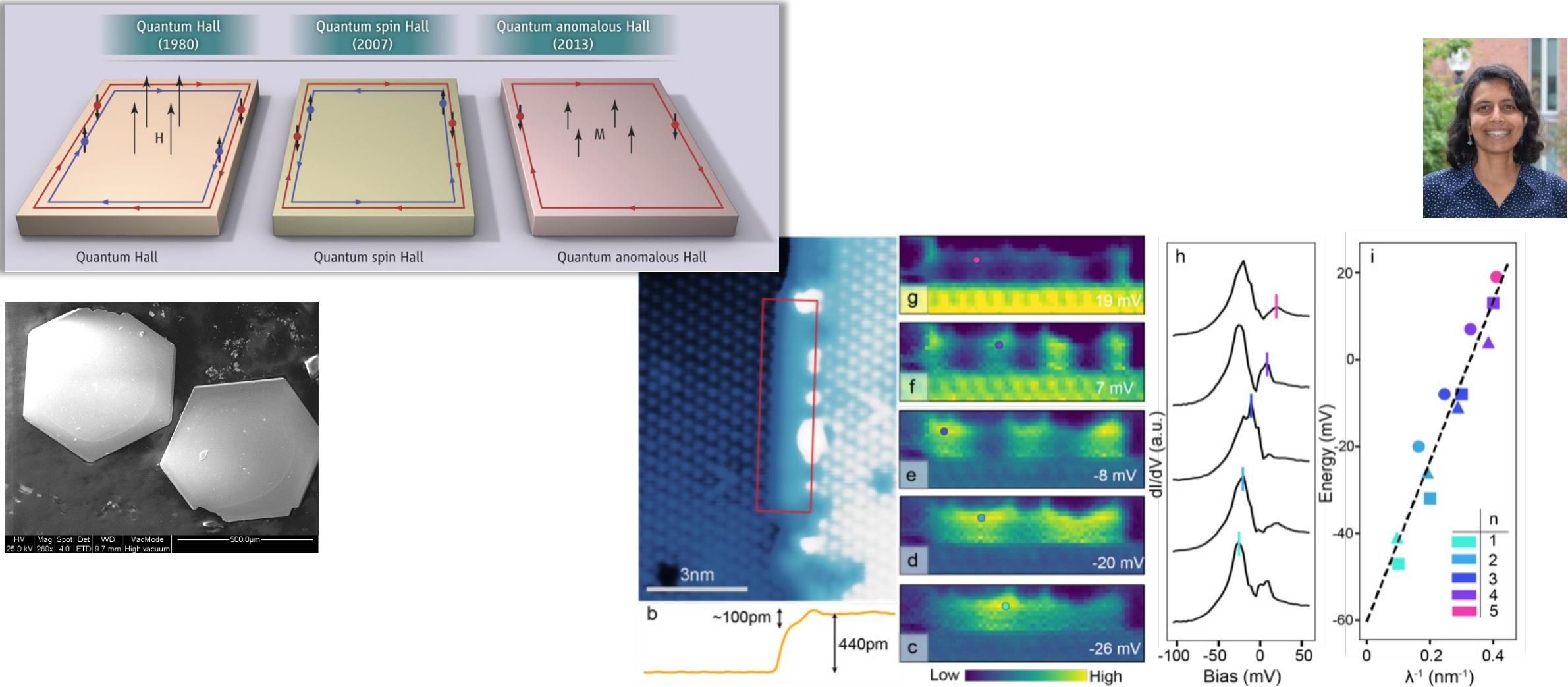
magnetic Weyl
semimetals as a recipe
for large Nernst effects



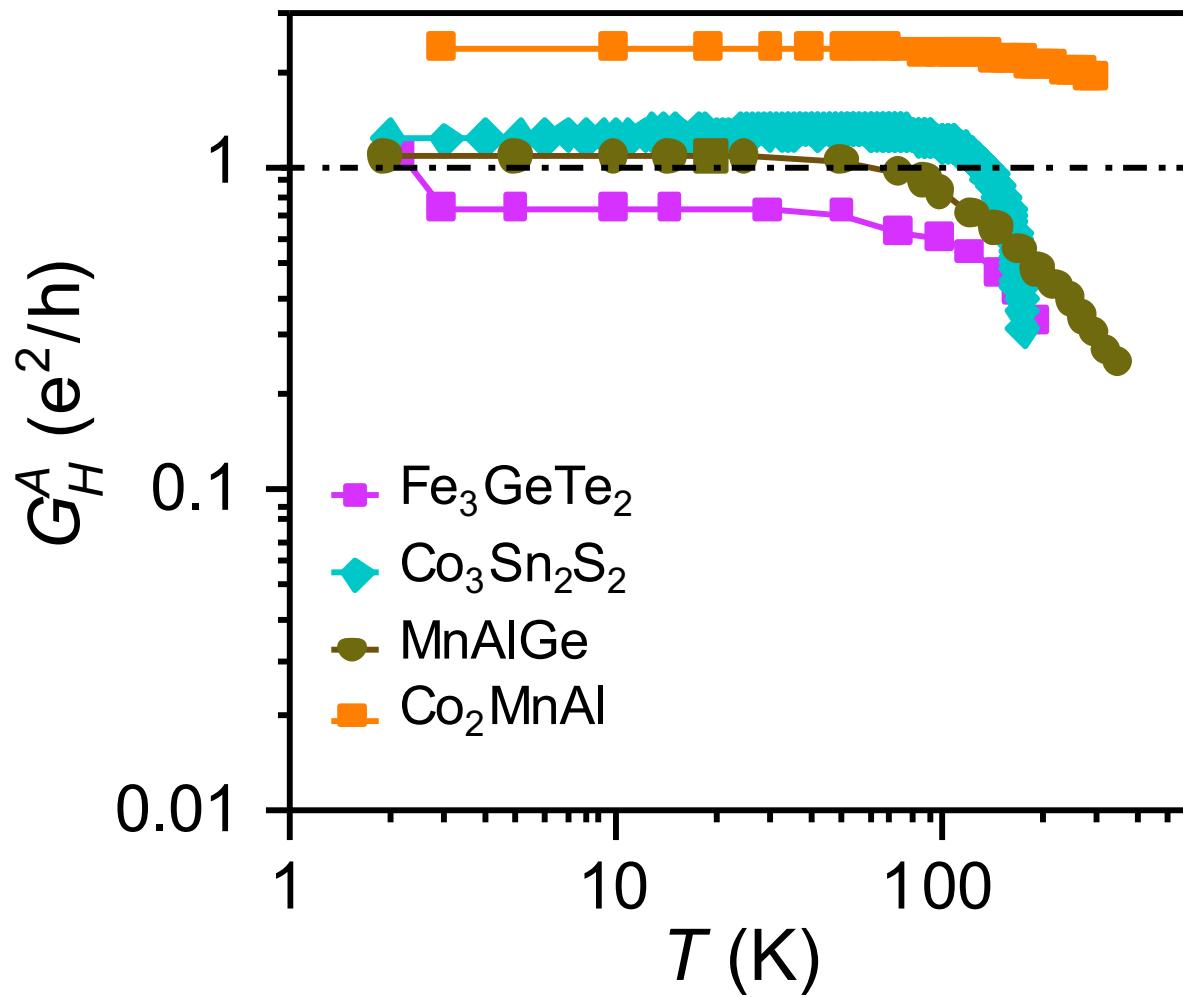
from 3D to 2D quantum effects



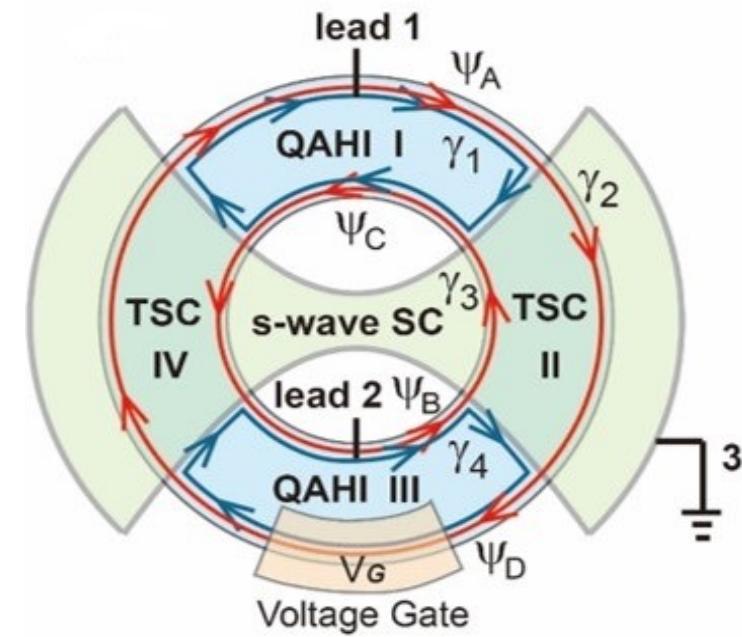
Quantum anomalous Hall effect?



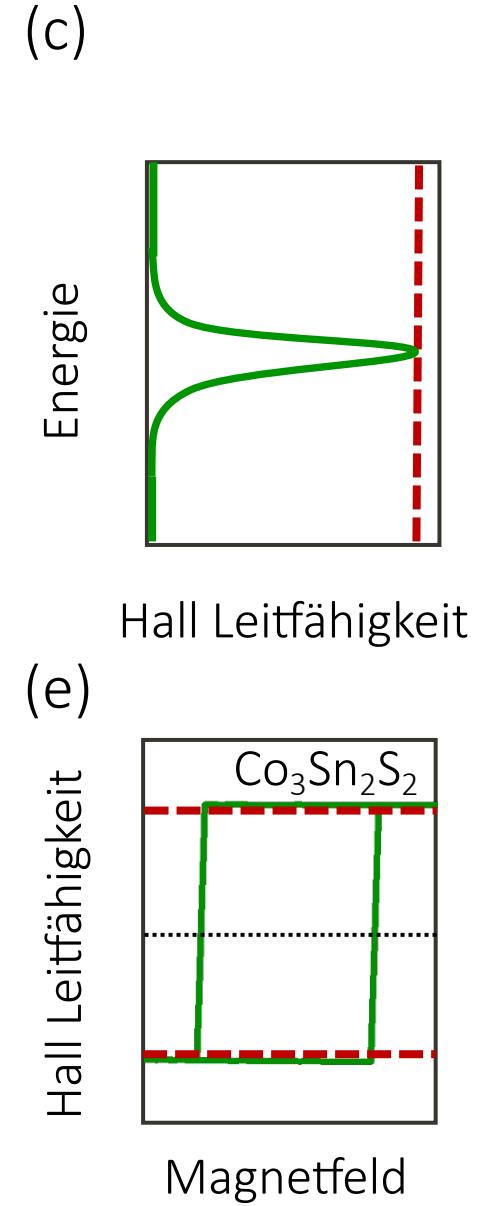
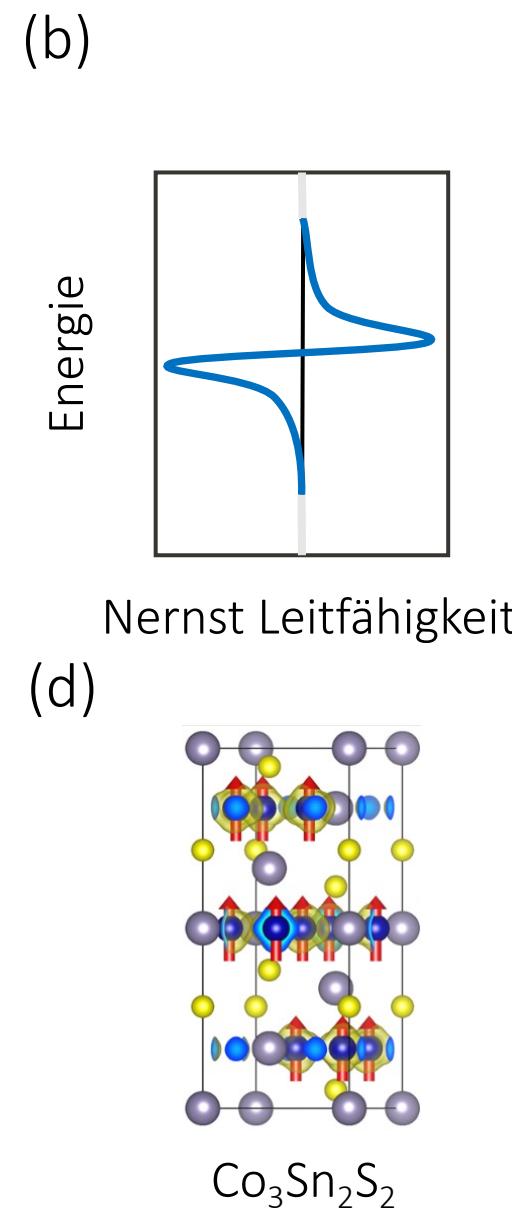
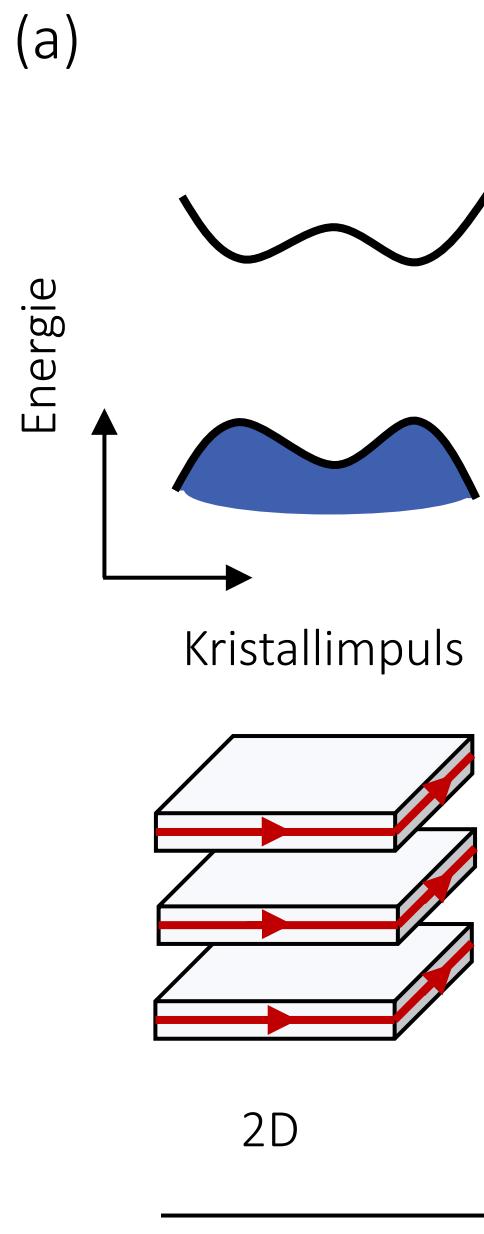
Berry curvature design



Qubit based on topology



Lian, et al., PNAS 2018

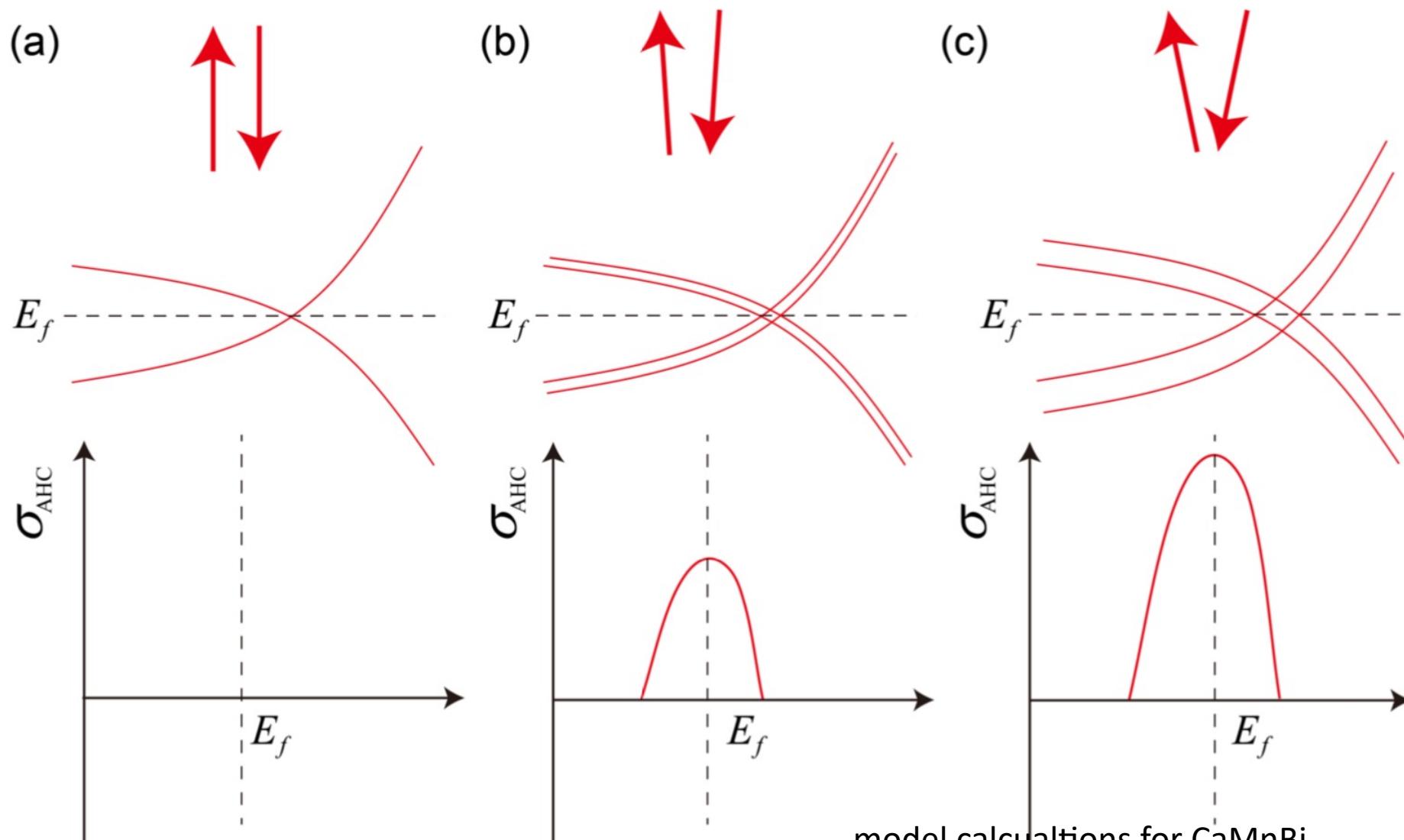


antiferromagnetic topological materials

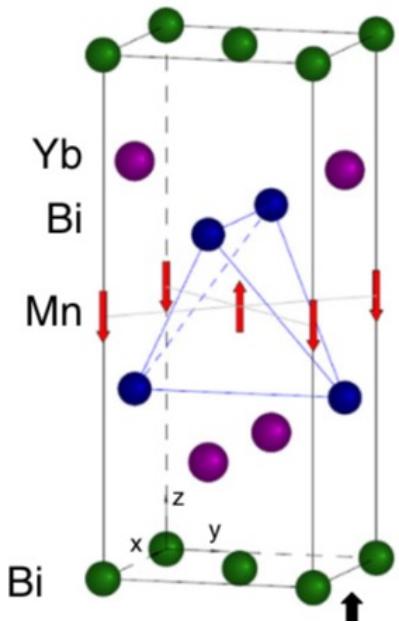
Table 3 | The magnetic topological materials identified in this work

Categories	Properties	Materials
I-A	Non-collinear manganese compounds	Mn ₃ GaC, Mn ₃ ZnC, Mn ₃ CuN, Mn ₃ Sn, Mn ₃ Ge, Mn ₃ Ir, Mn ₃ Pt, Mn ₅ Si ₃
I-B	Actinide intermetallic	UNiGa ₅ , UPtGa ₅ , NpRhGa ₅ , NpNiGa ₅
I-C	Rare-earth intermetallic	NdCo ₂ , TbCo ₂ , NpCo ₂ , PrAg DyCu, NdZn, TbMg, NdMg, Nd ₅ Si ₄ , Nd ₅ Ge ₄ , Ho ₂ RhIn ₈ , Er ₂ CoGa ₈ , Nd ₂ RhIn ₈ , Tm ₂ CoGa ₈ , Ho ₂ RhIn ₈ , DyCo ₂ Ga ₈ , TbCo ₂ Ga ₈ , Er ₂ Ni ₂ In, CeRu ₂ Al ₁₀ , Nd ₃ Ru ₄ Al ₁₂ , Pr ₃ Ru ₄ Al ₁₂ , ScMn ₆ Ge ₆ , YFe ₄ Ge ₄ , LuFe ₄ Ge ₄ , CeCoGe ₃
II-A	Metallic iron pnictides	LaFeAsO, CaFe ₂ As ₂ , EuFe ₂ As ₂ , BaFe ₂ As ₂ , Fe ₂ As, CaFe ₄ As ₃ , LaCrAsO, Cr ₂ As, CrAs, CrN
II-B	Semiconducting manganese pnictides	BaMn ₂ As ₂ BaMn ₂ Bi ₂ , CaMnBi ₂ , SrMnBi ₂ , CaMn ₂ Sb ₂ , CuMnAs, CuMnSb, Mn ₂ As
II-C	Rare-earth intermetallic compounds with the composition 1:2:2	PrNi ₂ Si ₂ , YbCo ₂ Si ₂ , DyCo ₂ Si ₂ , PrCo ₂ P ₂ , CeCo ₂ P ₂ , NdCo ₂ P ₂ , DyCu ₂ Si ₂ , CeRh ₂ Si ₂ , UAu ₂ Si ₂ , U ₂ Pd ₂ Sn, U ₂ Pd ₂ In, U ₂ Ni ₂ Sn, U ₂ Ni ₂ In, U ₂ Rh ₂ Sn
II-D	Rare-earth ternary compounds of the composition 1:1:1	CeMgPb, PrMgPb, NdMgPb, TmMgPb
III-A	Semiconducting actinides/rare-earth pnictides	HoP, UP, UP ₂ , UAs, NpS, NpSe, NpTe, NpSb, NpBi, U ₃ As ₄ , U ₃ P ₄
III-B	Metallic oxides	Ag ₂ NiO ₂ , AgNiO ₂ , Ca ₃ Ru ₂ O ₇ , Double perovskite Sr ₃ CoIrO ₆
III-C	Metal-to-insulator transition compounds	NiS ₂ , Sr ₂ Mn ₃ As ₂ O ₂
III-D	Semiconducting and insulating oxides, borates, hydroxides, silicates and phosphate	LuFeO ₃ , PdNiO ₃ , ErVO ₃ , DyVO ₃ , MnGeO ₃ , Tm ₂ Mn ₂ O ₇ , Yb ₂ Sn ₂ O ₇ , Tb ₂ Sn ₂ O ₇ , Ho ₂ Ru ₂ O ₇ , Er ₂ Ti ₂ O ₇ , Tb ₂ Ti ₂ O ₇ , Cd ₂ Os ₂ O ₇ , Ho ₂ Ru ₂ O ₇ , Cr ₂ ReO ₆ , NiCr ₂ O ₄ , MnV ₂ O ₄ , Co ₂ SiO ₄ , Fe ₂ SiO ₄ , PrFe ₃ (BO ₃) ₄ , KCo ₄ (PO ₄) ₃ , CoPS ₃ , SrMn(VO ₄)(OH), Ba ₅ Co ₅ ClO ₁₃ , FeI ₂

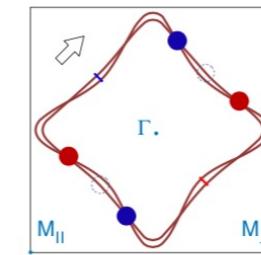
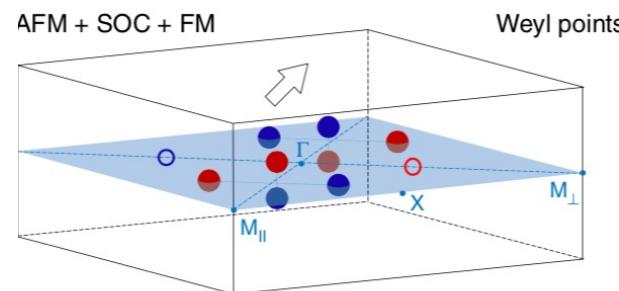
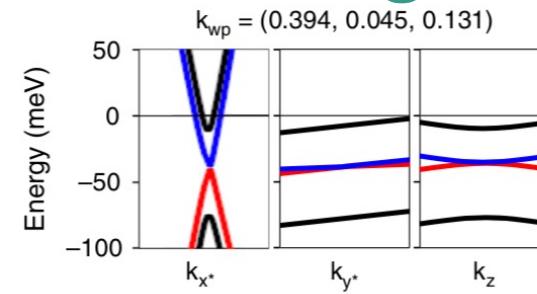
canting and anomalous Hall



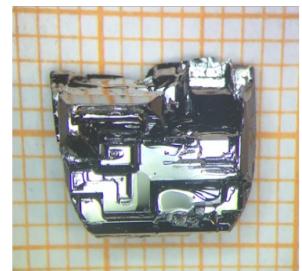
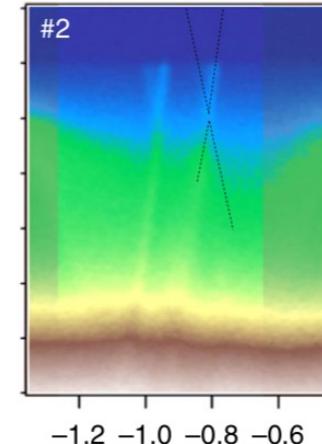
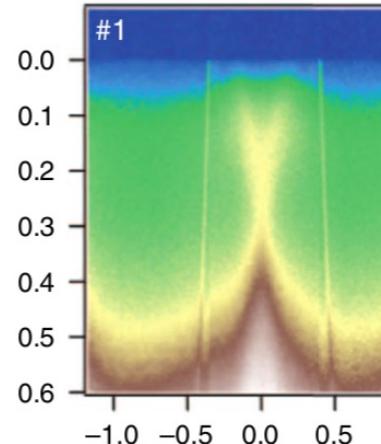
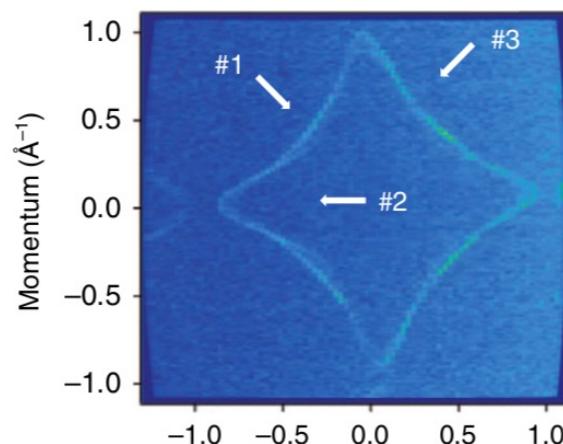
non-collinear antiferromagnet YbMnBi_2



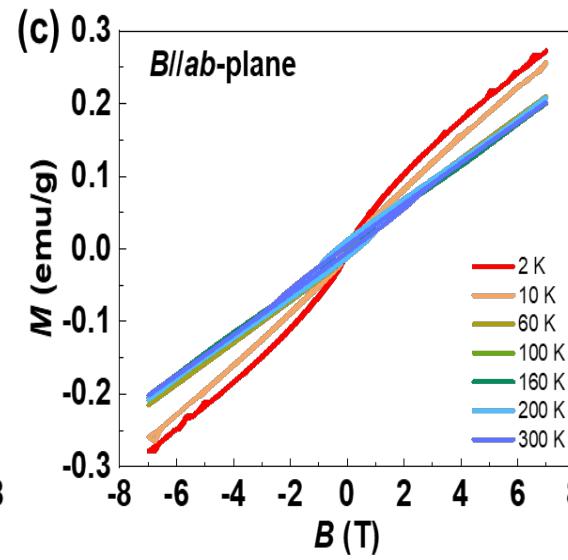
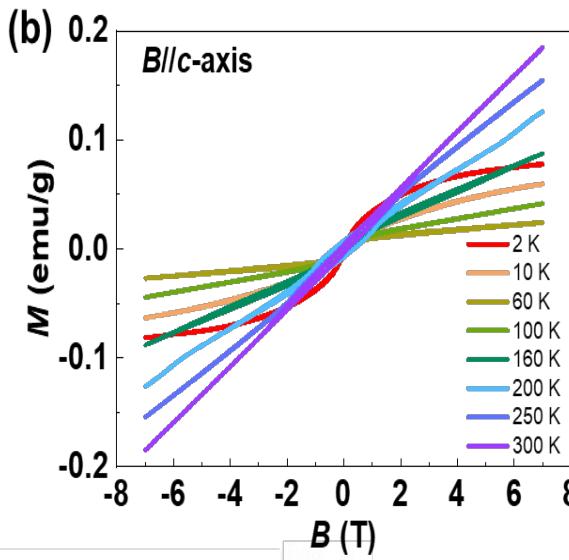
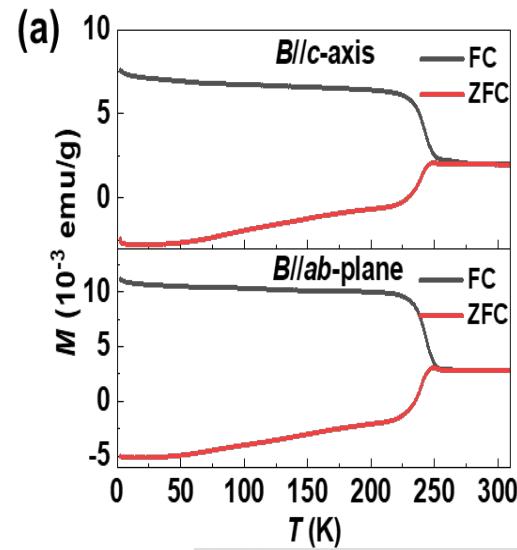
Canted AFM in YbMnBi_2
Mn
In-plane FM component



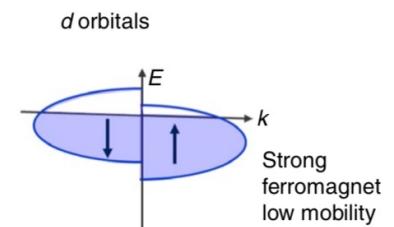
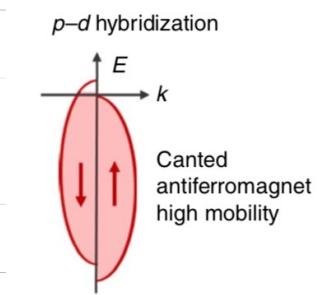
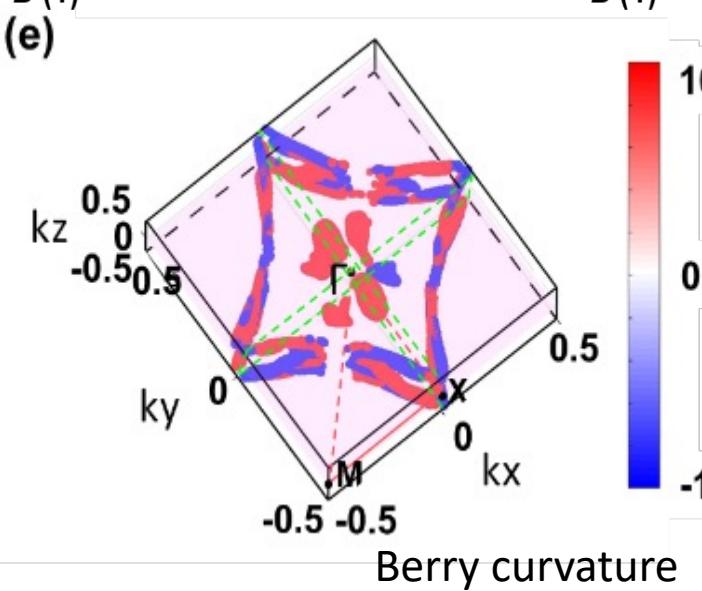
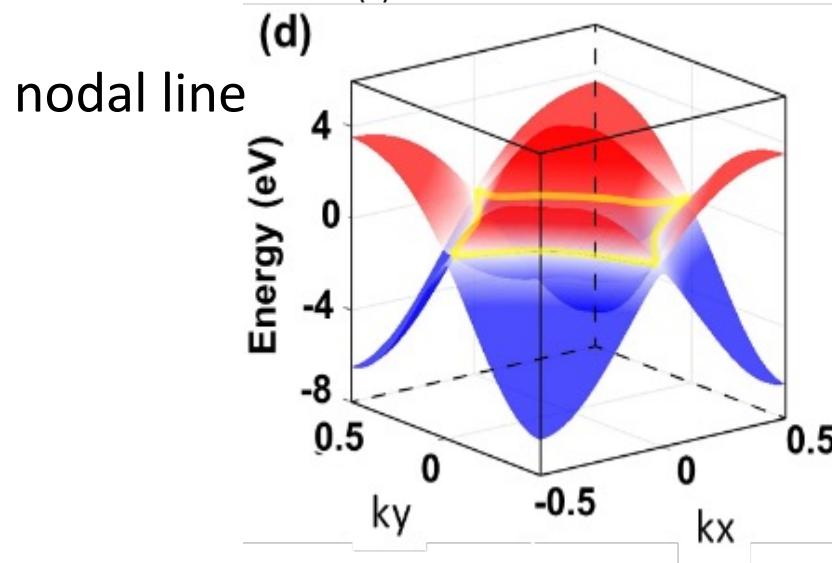
magnetic ordering below 290K
canted Mn atoms



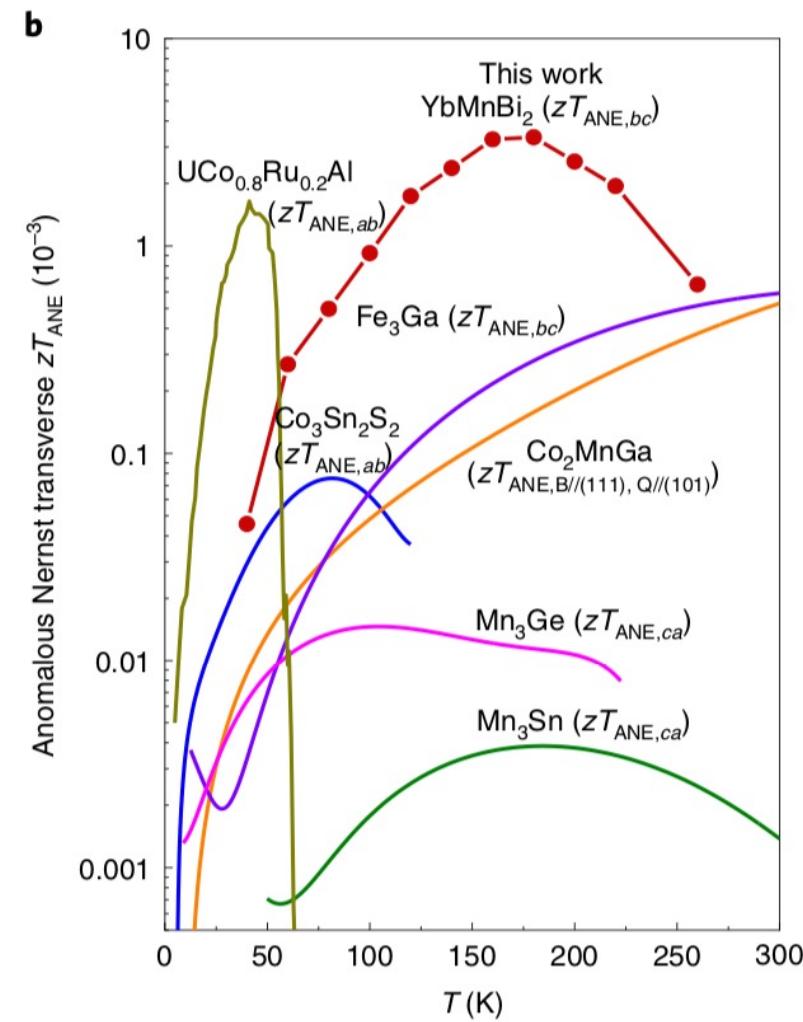
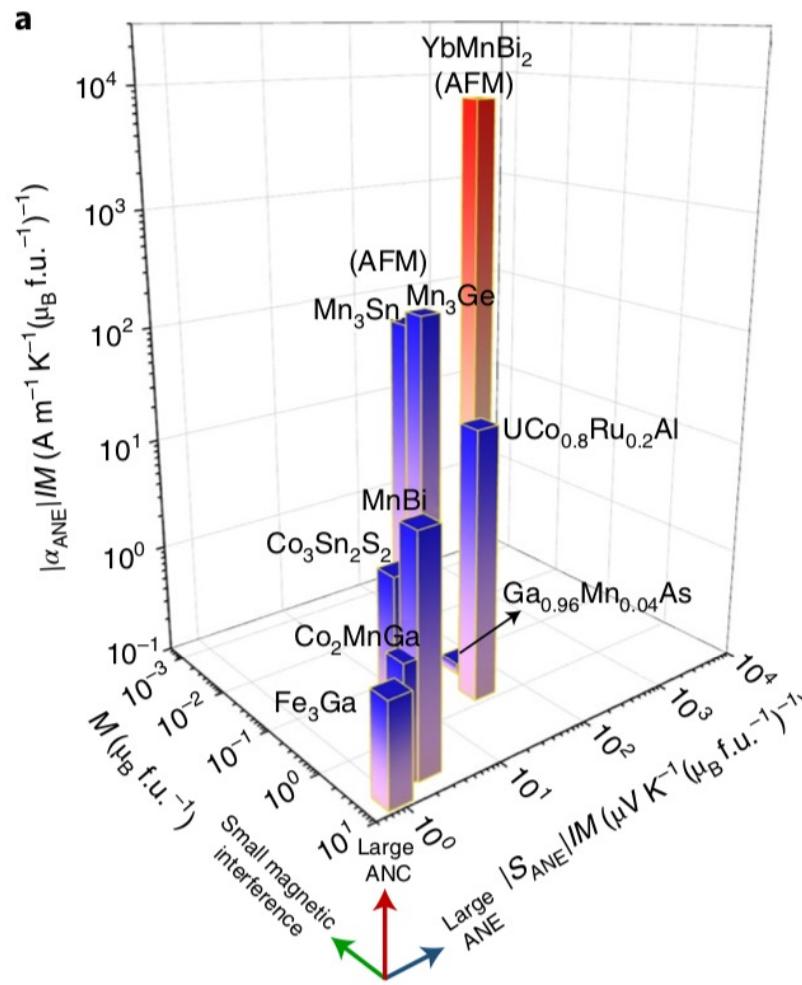
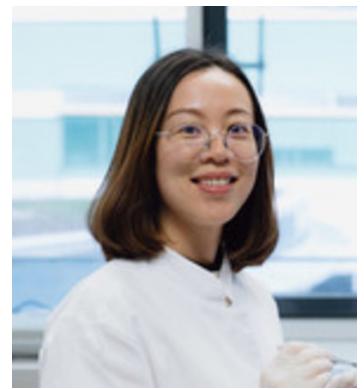
non-collinear antiferromagnet YbMnBi_2



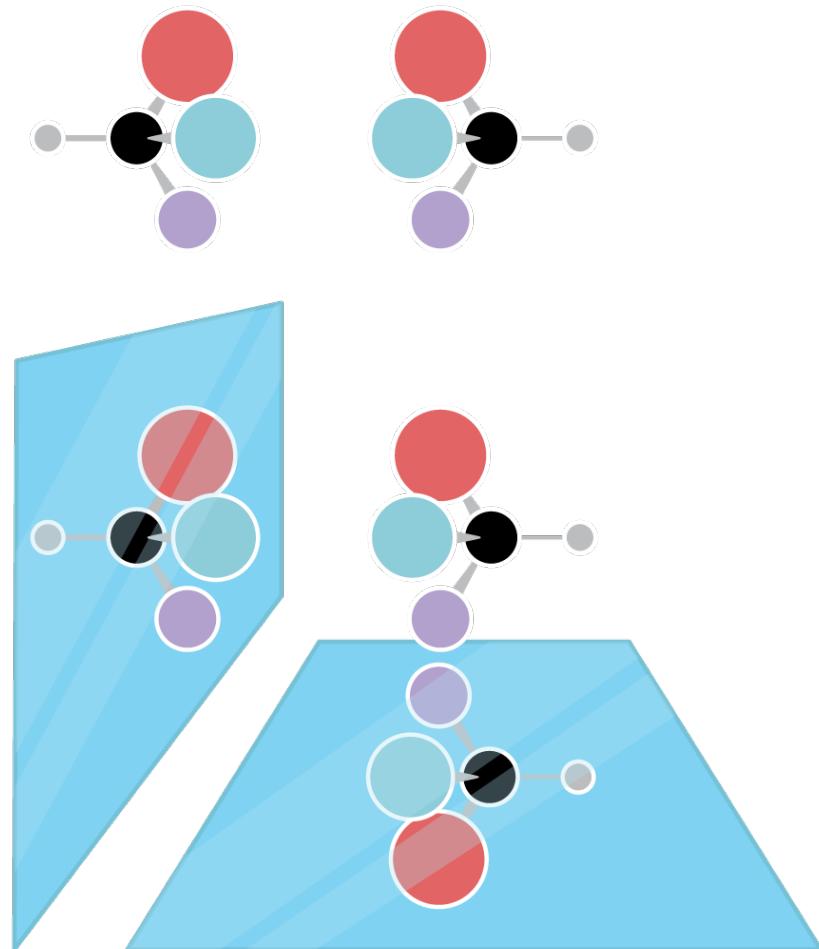
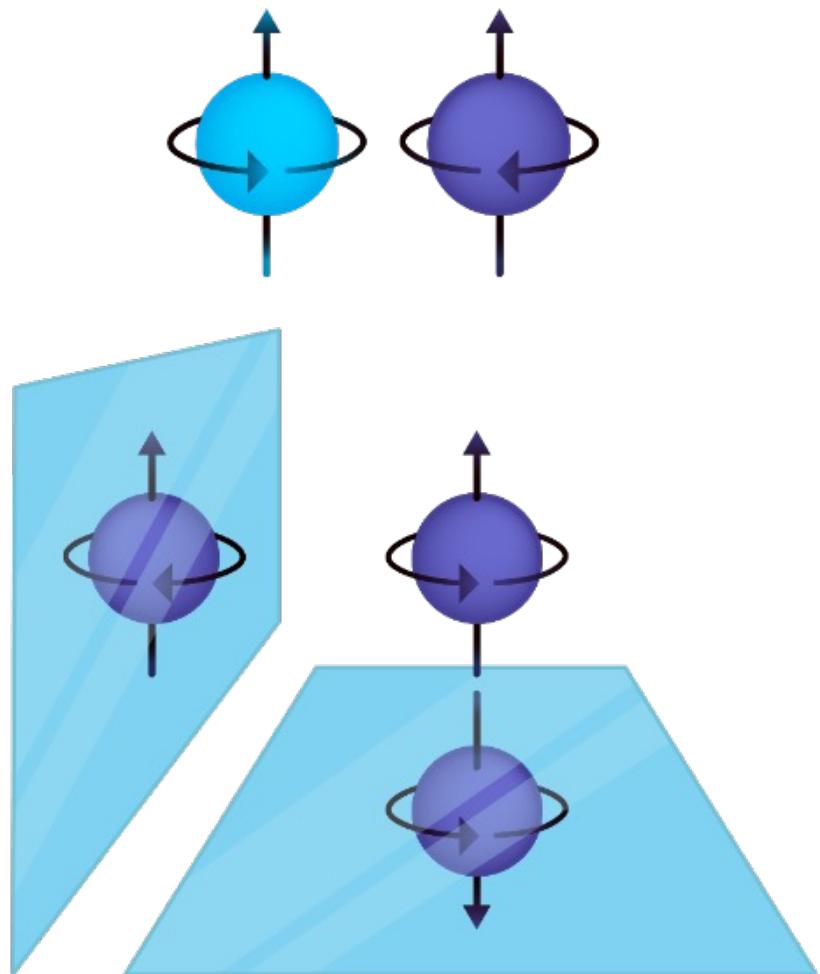
evidence for a small moment $0.0015 \mu_B$ per Mn



non-collinear antiferromagnet YbMnBi_2



chirality



new fermions

RESEARCH

RESEARCH ARTICLE SUMMARY

TOPOLOGICAL MATTER

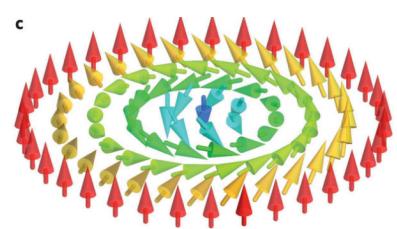
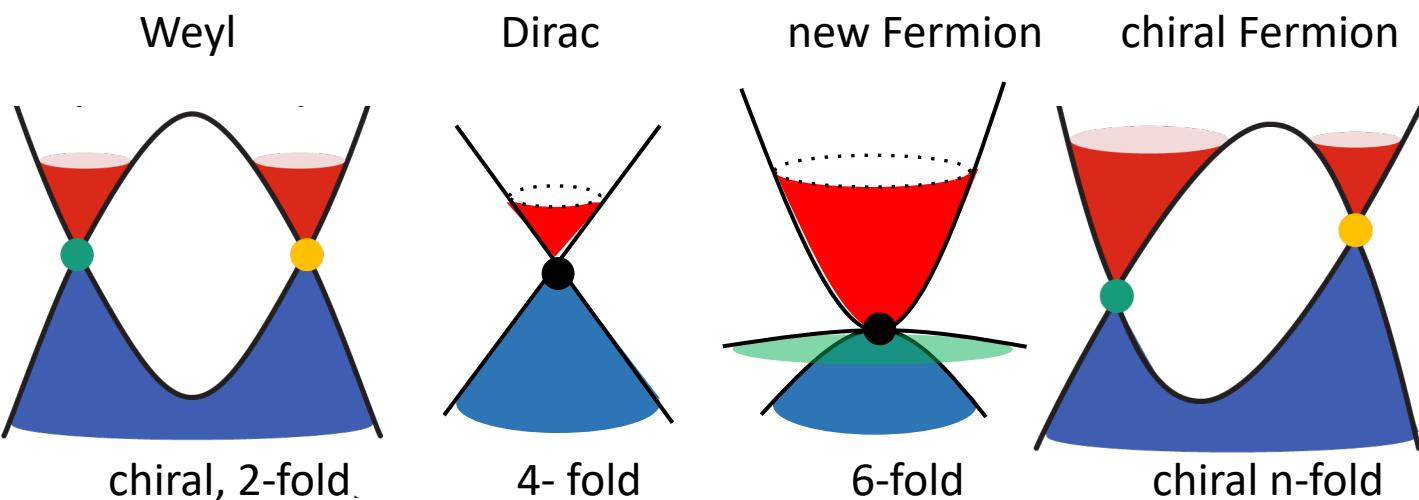
Beyond Dirac and Weyl fermions: Unconventional quasiparticles in conventional crystals

Barry Bradlyn,* Jennifer Cano,* Zhijun Wang,* M. G. Vergniory, C. Felser,
R. J. Cava, B. Andrei Bernevig†

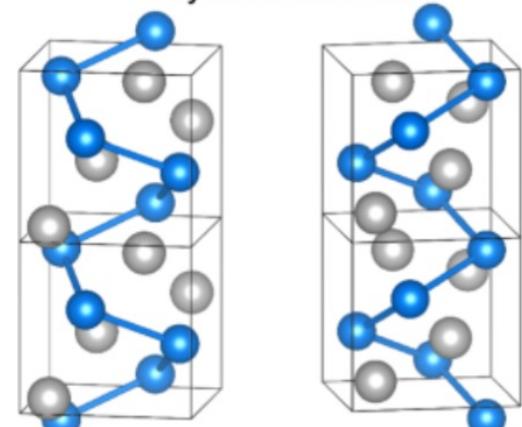
free fermionic excitations in solid-state systems that have
no high-energy counterparts.

Some of these new Fermions are even chiral

- Chiral Crystals
 - B20, Skyrmions, CoSi, MnSi, PdGa, RhSi
- Superconductors
 - A15 superconductors: Nb₃Sn, Li₂Pd₃B



Enantiomer A Enantiomer B
Crystal structure



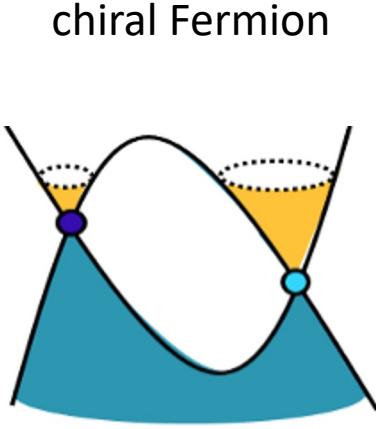
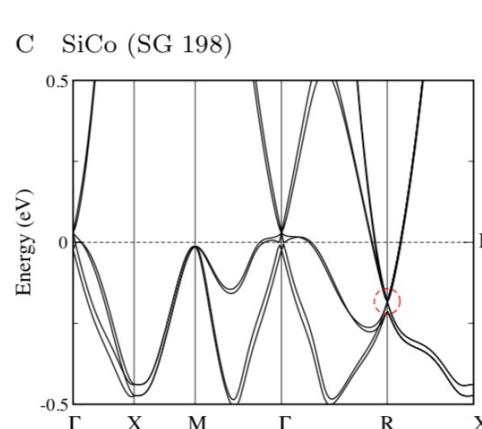


E SUMMARY

TOPOLOGICAL MATTER

Beyond Dirac and Weyl fermions: Unconventional quasiparticles in conventional crystals

Barry Bradlyn,* Jennifer Cano,* Zhijun Wang,* M. G. Vergniory, C. Felser,
R. J. Cava, B. Andrei Bernevig†



chiral fermions

ARTICLES

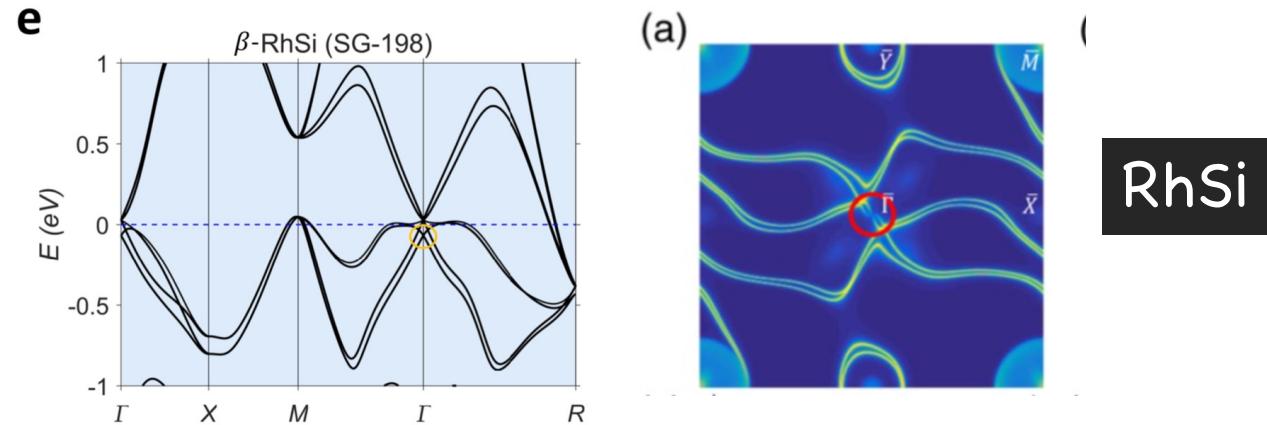
<https://doi.org/10.1038/s41563-018-0169-3>



nature
mat

Topological quantum properties of chiral crystals

Guoqing Chang^{1,2,3,4,12}, Benjamin J. Wieder^{5,6,7,12}, Frank Schindler^{8,12}, Daniel S. Sanchez¹, Ilya Belopolski¹, Shin-Ming Huang⁹, Bahadur Singh¹⁰, Di Wu^{2,3}, Tay-Rong Chang¹⁰, Titus Neupert¹², Su-Yang Xu^{1*}, Hsin Lin^{2,3,4*} and M. Zahid Hasan^{1,11*}



Unconventional Chiral Fermions and Large Topological Fermi Arcs in RhSi

Guoqing Chang, Su-Yang Xu, Benjamin J. Wieder, Daniel S. Sanchez, Shin-Ming Huang, Ilya Belopolski, Tay-Rong Chang, Songtian Zhang, Arun Bansil, Hsin Lin, and M. Zahid Hasan
Phys. Rev. Lett. **119**, 206401 – Published 17 November 2017

Bringing order to the expanding fermion zoo

Carlo Beenakker Commentary

Heisenberg (1930): We observe space as a continuum, but we might entertain the thought that there is an underlying lattice and that space is actually a crystal. Which particles would inhabit such a lattice world? Werner Heisenberg *Gitterwelt* (lattice world) **hosted electrons that could morph into protons, photons that were not massless**, and more peculiarities that compelled him to abandon “this completely crazy idea”

chirality and topology

chiral crystals
optical activity

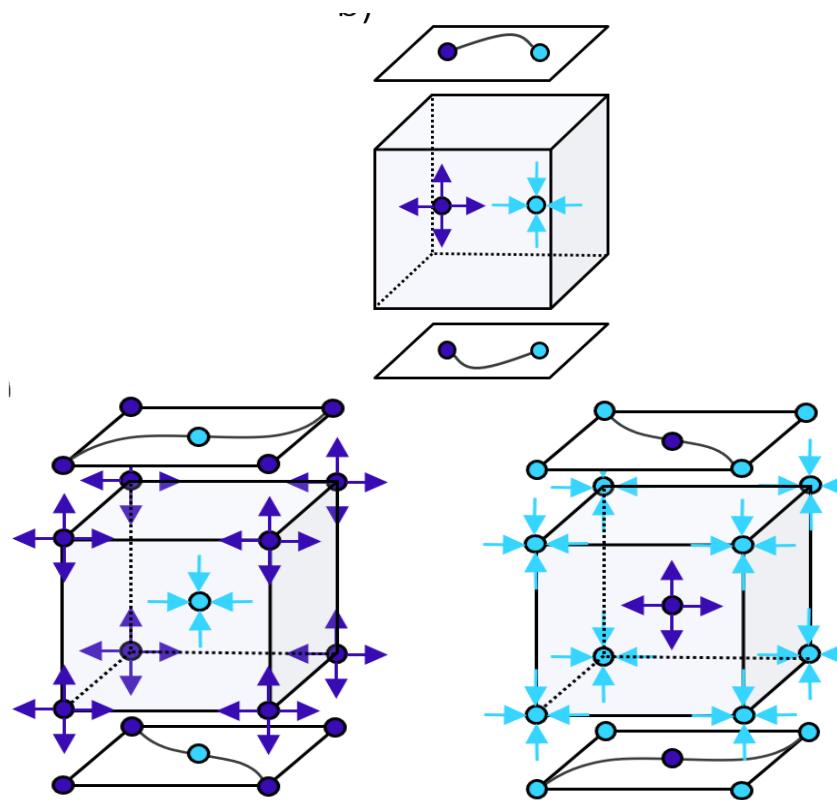
...

topological crystals
unusual surface states
large photogalvanic effect

...

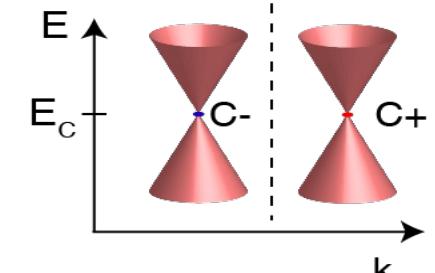
new Fermions
largest Fermi arc
Quantized circular photogalvanic effect

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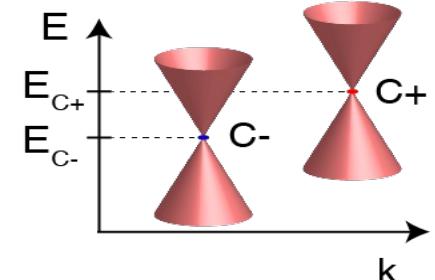


Non-chiral topological semimetal

Mirror



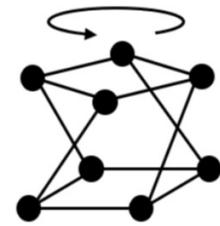
Chiral topological semimetal



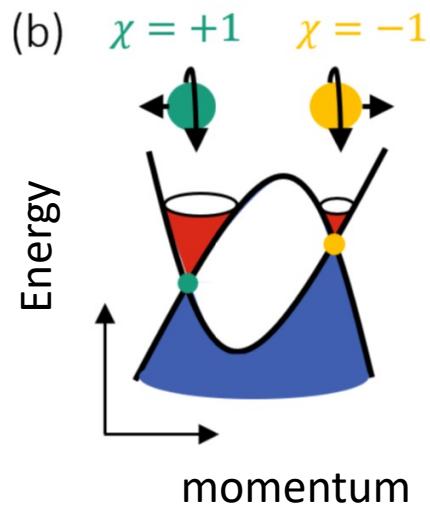
chirality and chiral anomaly

(a)

Enantominer A

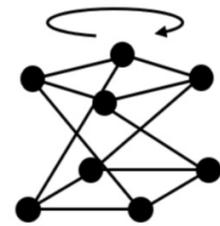


(b)

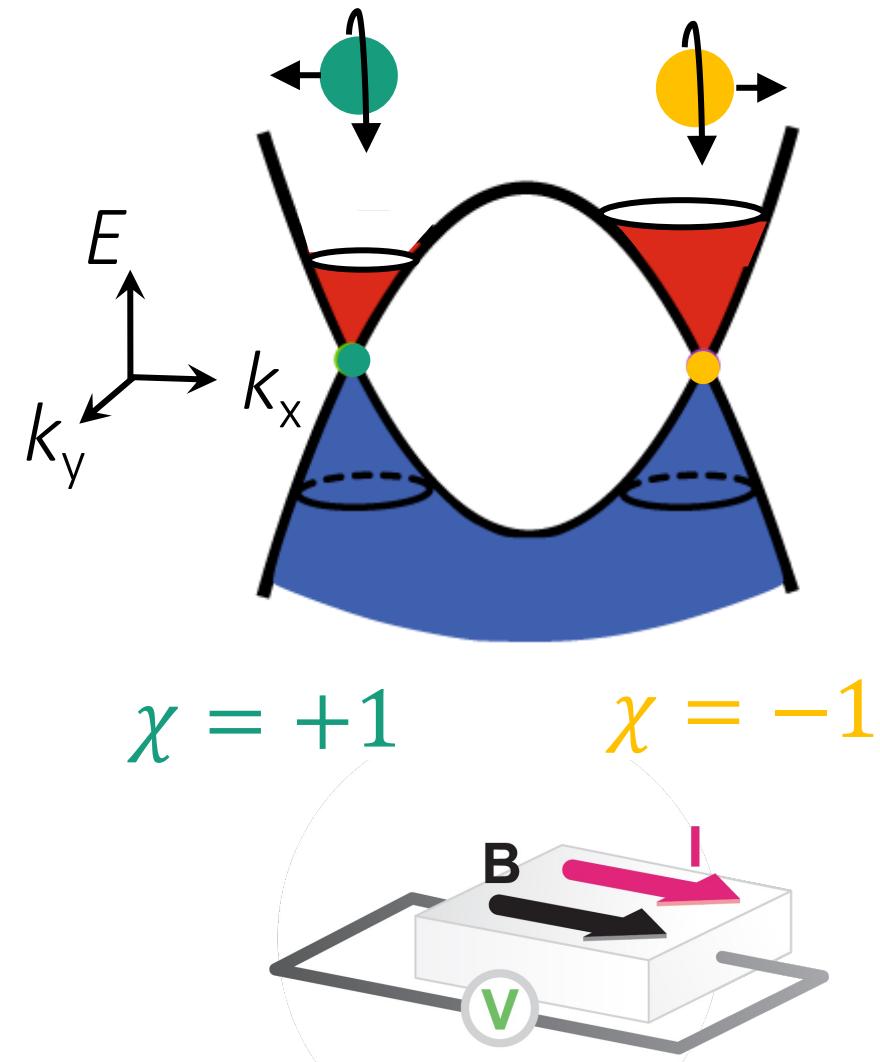
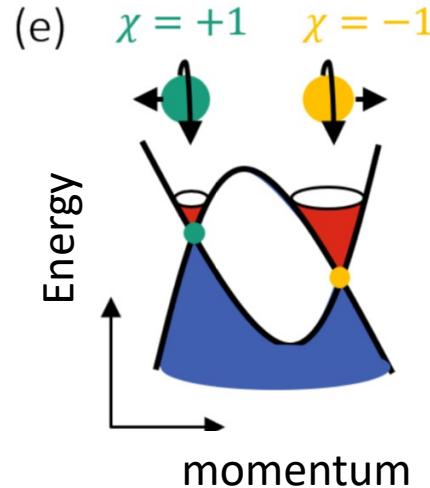


(d)

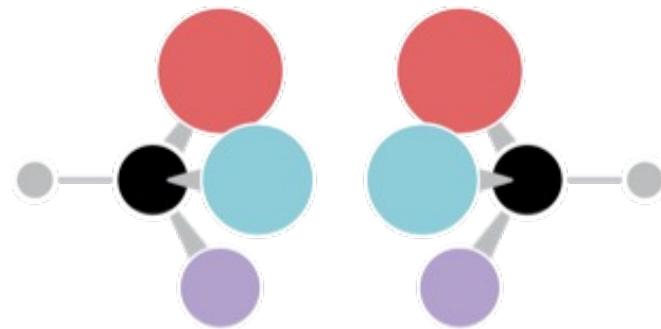
Enantominer B



(e)



chirality in chemistry

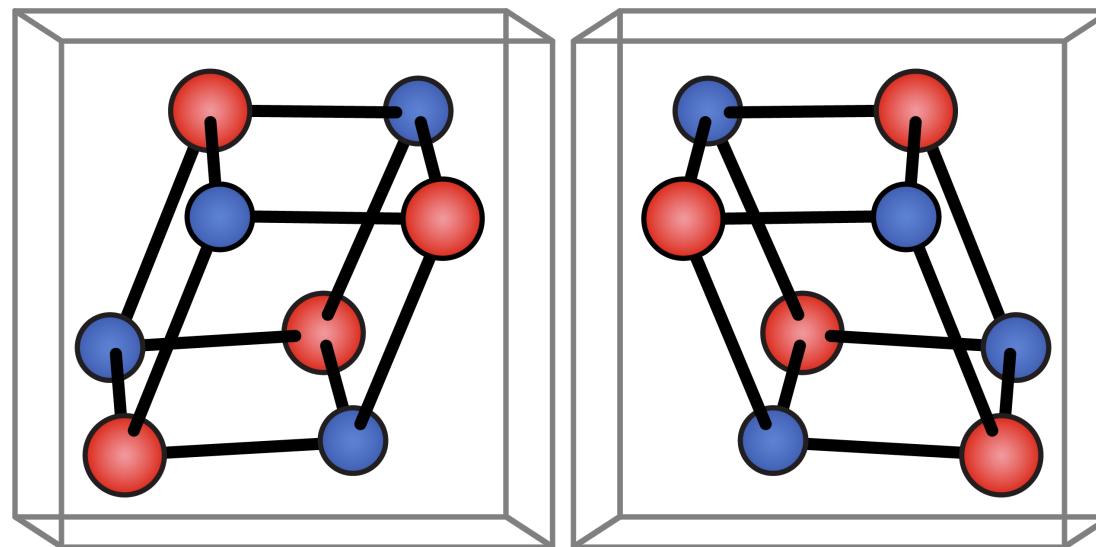


Molecules with different chiralities have different optical and catalytic properties

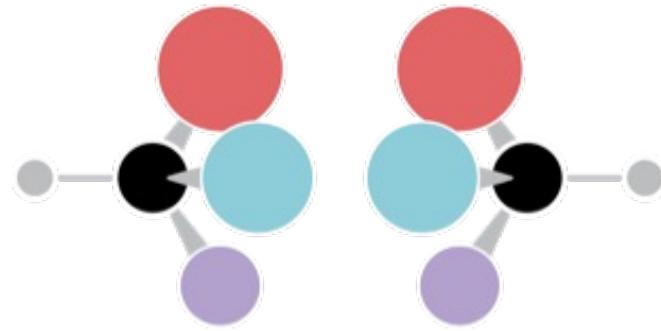


Lord Kelvin: An object or a system is *chiral* if it is distinguishable from its mirror image; that is, it cannot be superposed onto it.

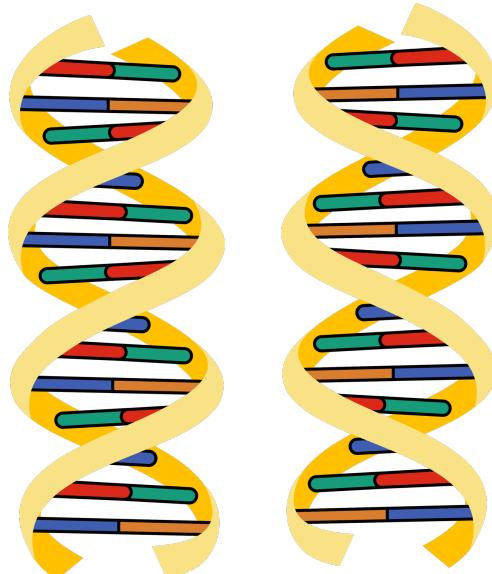
In terms of point groups, all chiral molecules lack an improper axis of rotation (S_n)



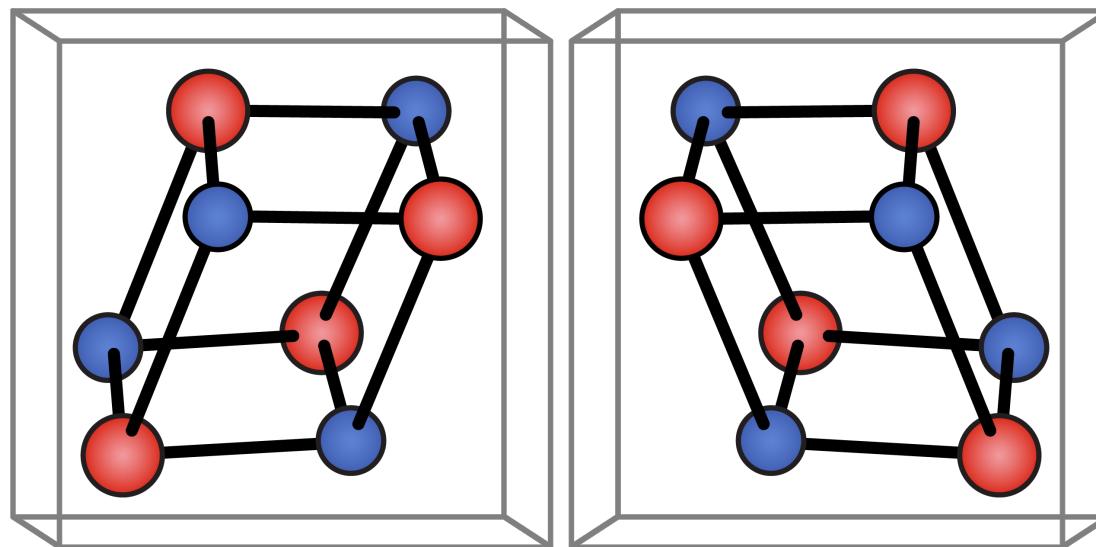
chirality in chemistry



Molecules with different chiralities have different optical and catalytic properties



Chiral space groups contain symmetry operations of the first kind (rotation). There are 11 pairs of enantiomorphous space groups (e.g. P61 and P65) which are chiral. 43 achiral space groups can host a chiral crystal structure.



synthesis

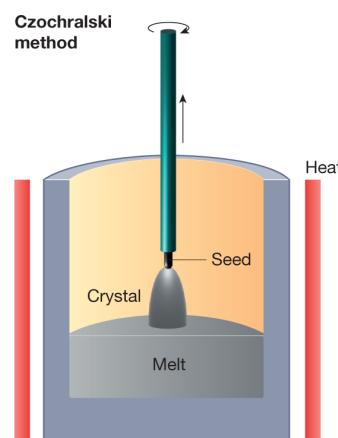
RhSi



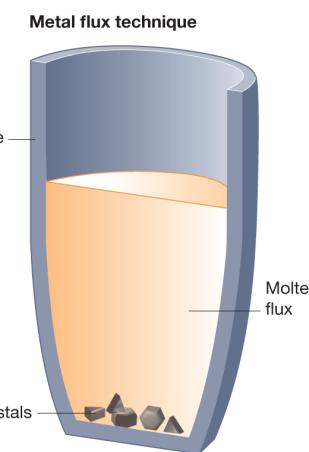
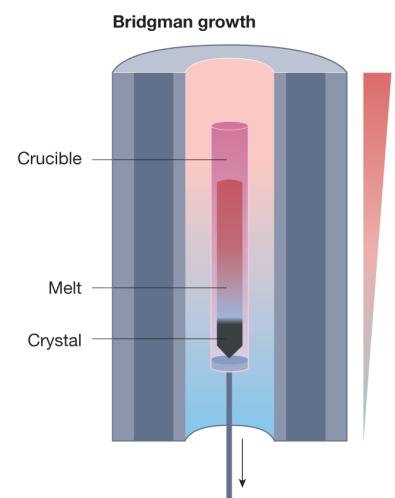
Flux +
Bridgman
CoSi



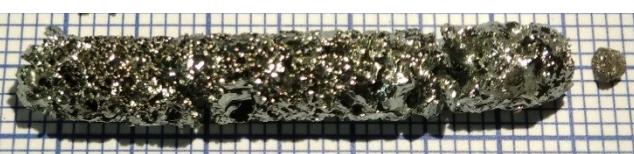
Chemical vapor transport
Bridgman, Flux



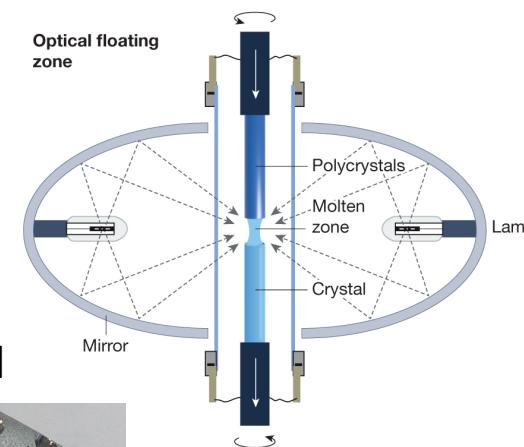
Self-Flux

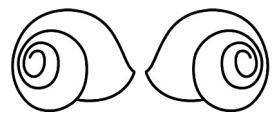


Self-Flux



PdGa: Chiral-1, Right handed

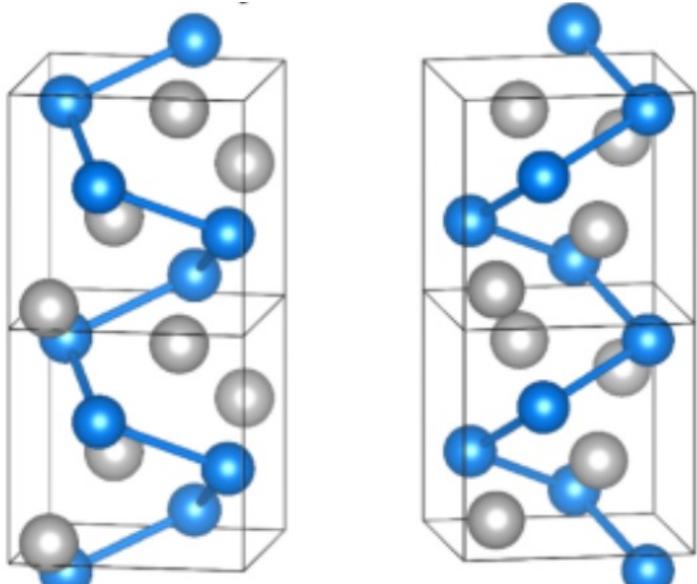




chiral crystals which host new fermions

Compounds with the B20 structure

Enantiomer A and Enantiomer B
Crystal structure



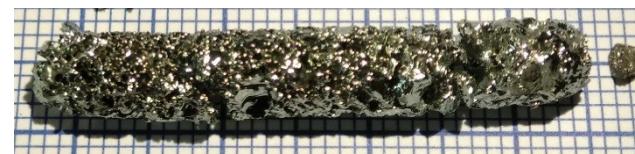
PdGa

Chiral-1, Right handed



PdGa

Chiral-2, Left handed

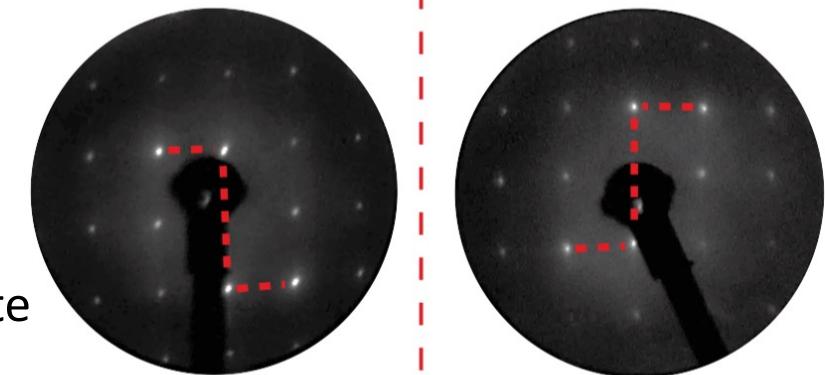
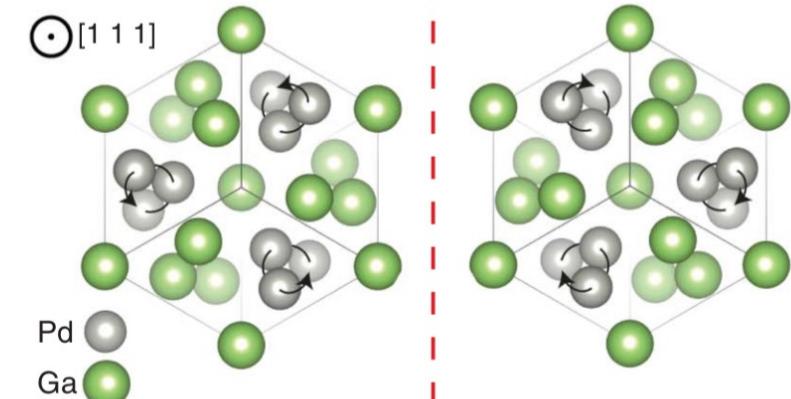


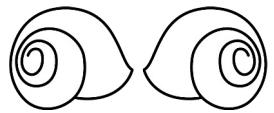
PdGa crystallizes in space group $P2_13$ (No. 198) 4- and 6-fold degenerate Fermions: more examples include CoSi, RhSi, PtAl, CoGe, RhGe, PtGa, PtAl and magnetic MnSi and FeGe

Mirror

Enantiomer A
right-handed
Pd-helix along (111)

Enantiomer B
left-handed
Pd-helix along (111)



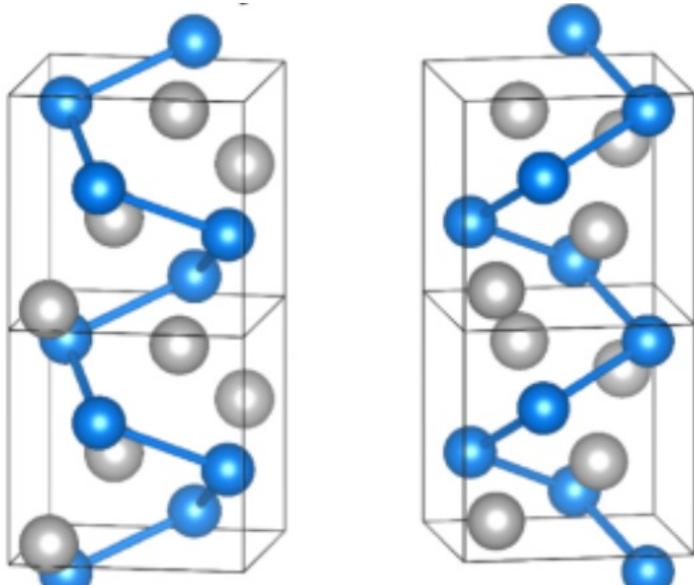


homochiral crystals

Compounds with the B20 structure

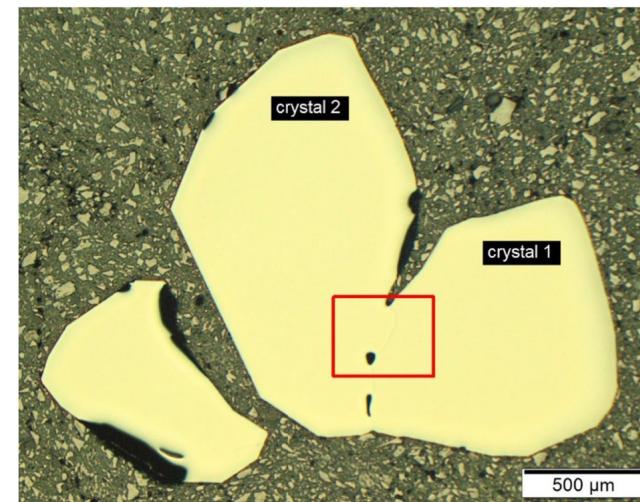
Enantiomer A and Enantiomer B

Crystal structure



CoSi

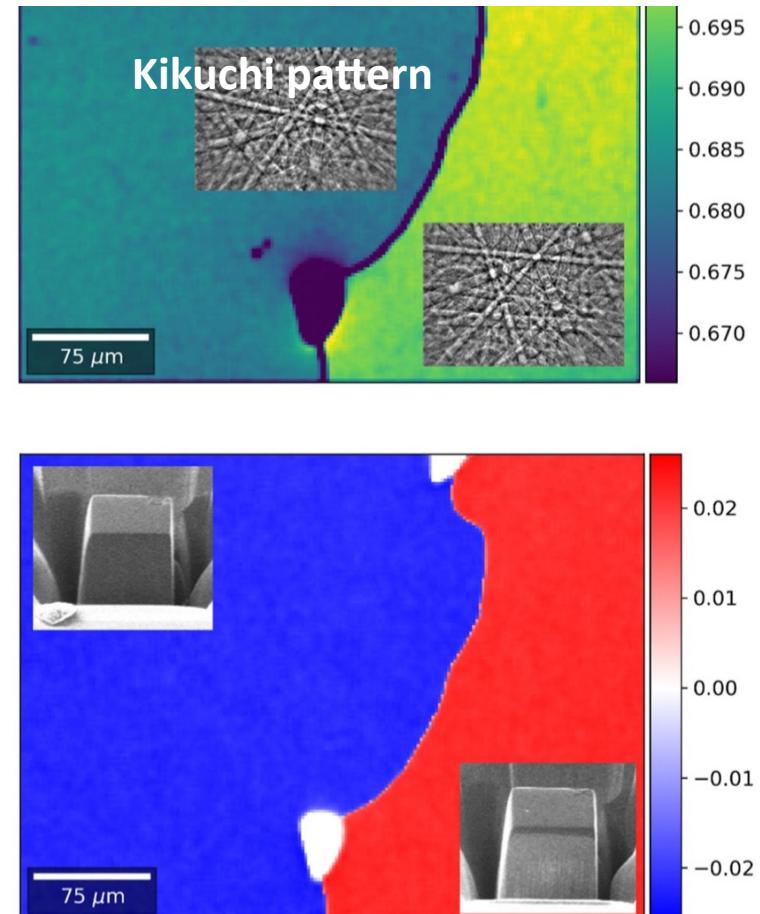
- X-ray (volume sensitive)
- electron backscatter diffraction (EBSD) (surface sensitive)

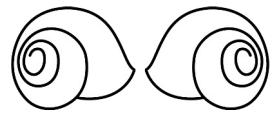


Fermions: more examples include CoSi, RhSi, PtAl, CoGe, RhGe, PtGa, PtAl and magnetic MnSi and FeGe

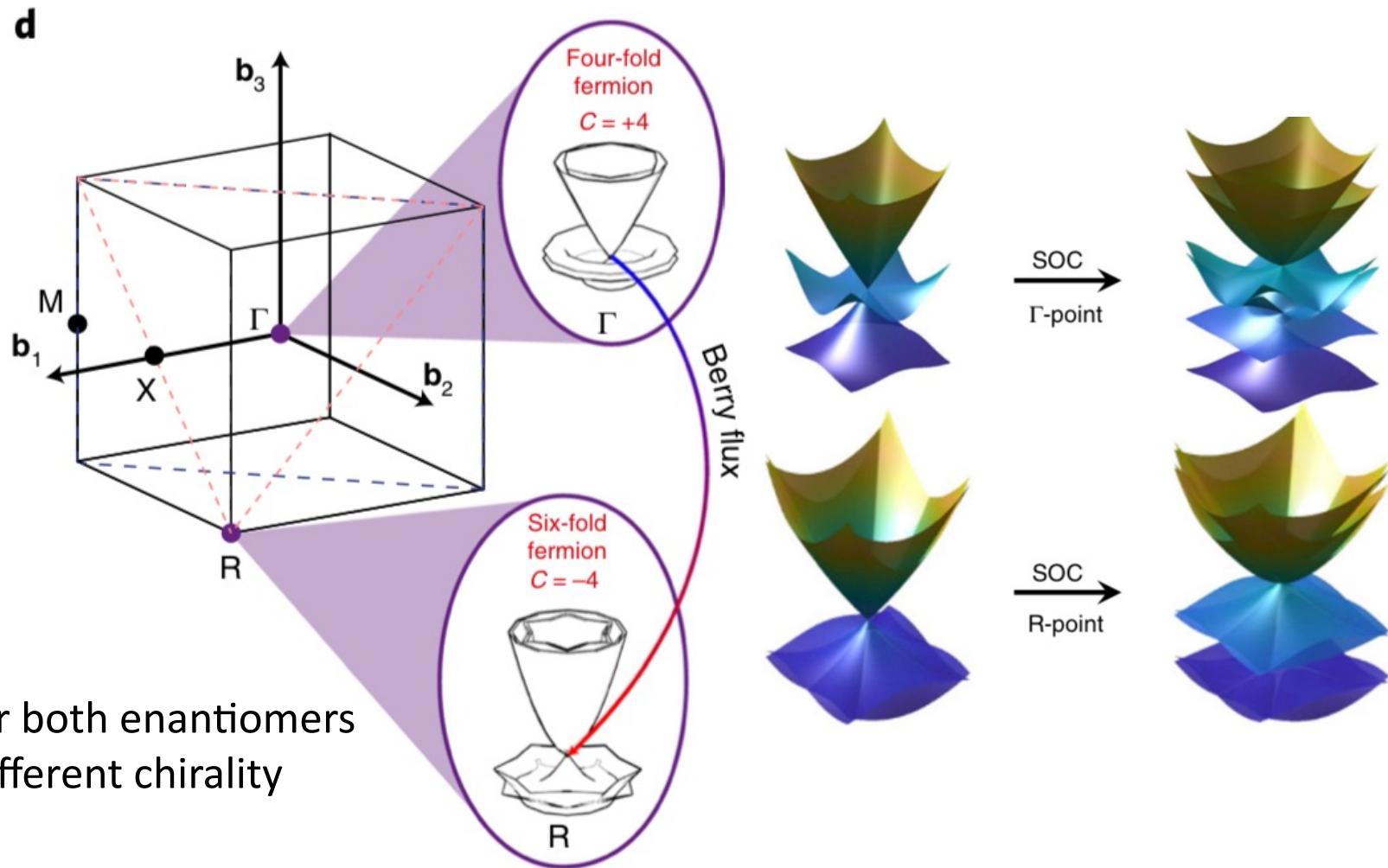
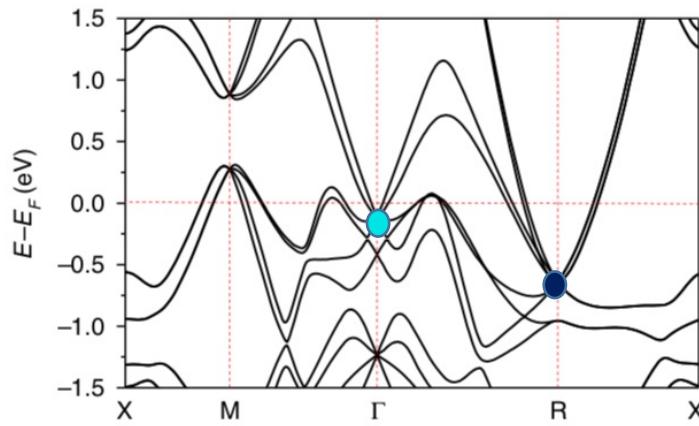
Absolute Structure from Scanning Electron Microscopy

Ulrich Burkhardt^{1*}, Horst Borrmann¹, Philip Moll^{1,2}, Marcus Schmidt¹, Yuri Grin¹ & Aimo Winkelmann^{1,3,4}





new fermions with chiral surface states

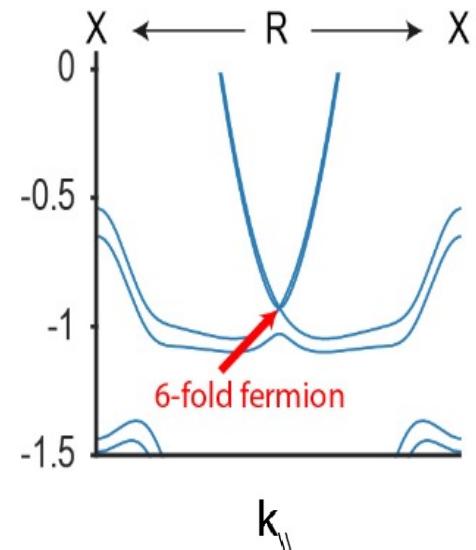
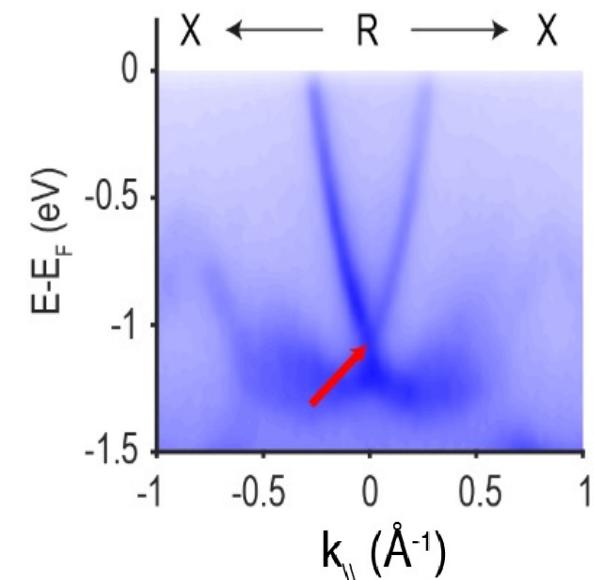
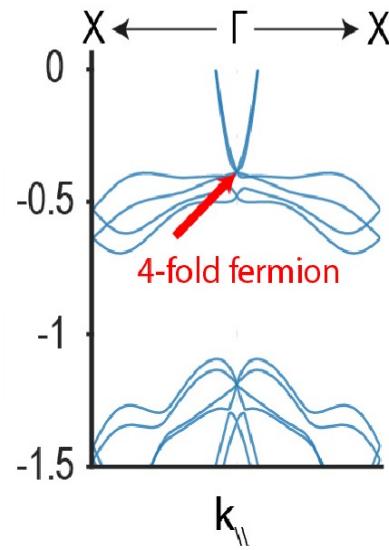
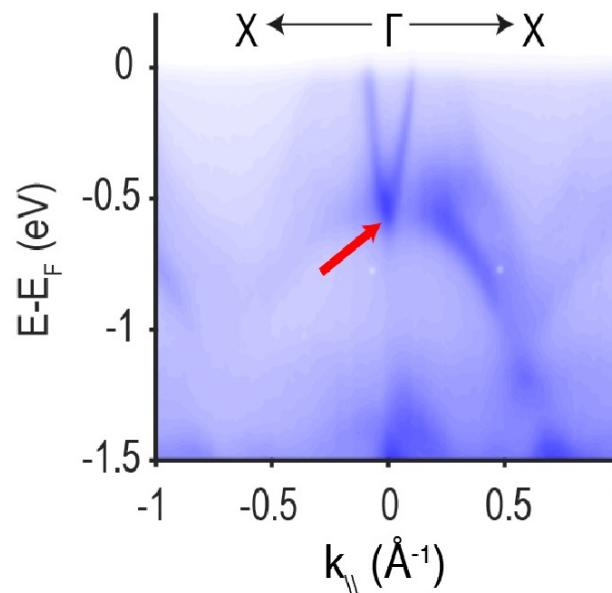
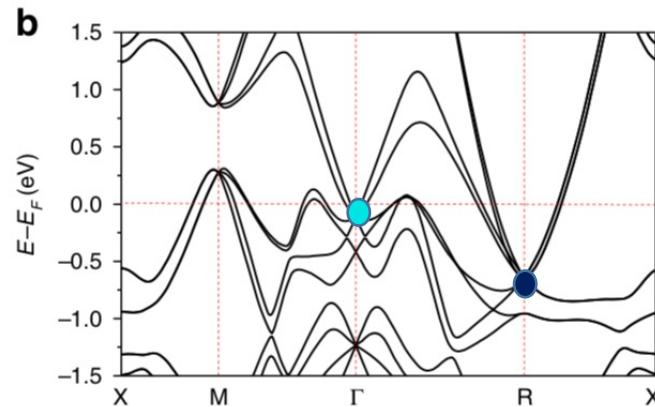
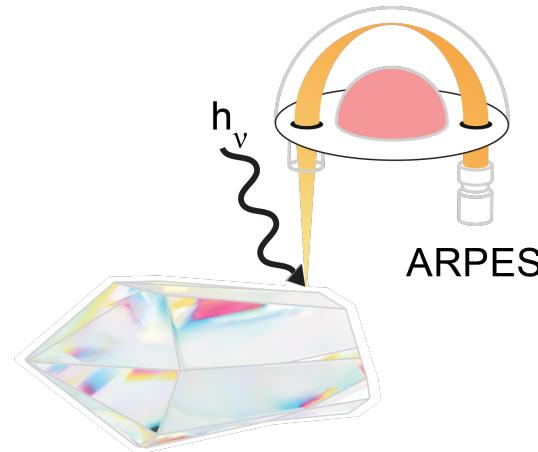


the band structure is the same for both enantiomers
however the Weyl points have different chirality



new fermions with chiral surface states

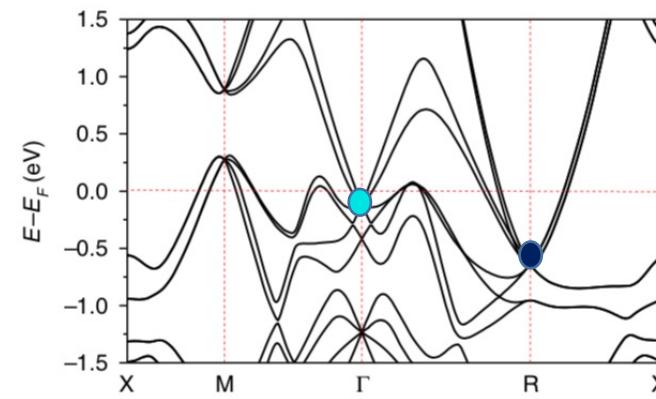
PtAl



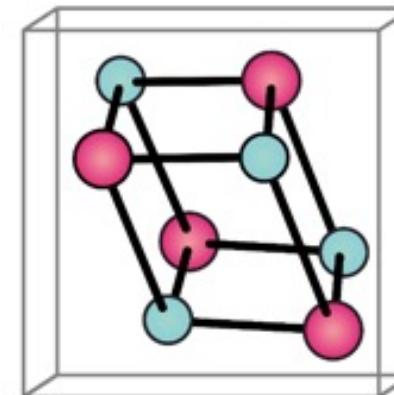
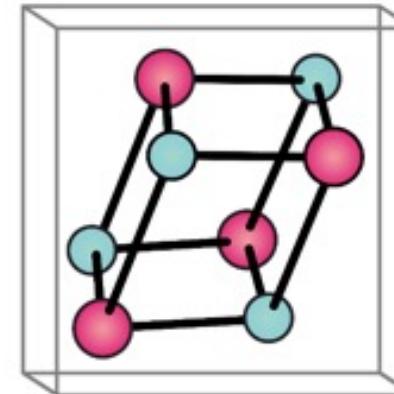
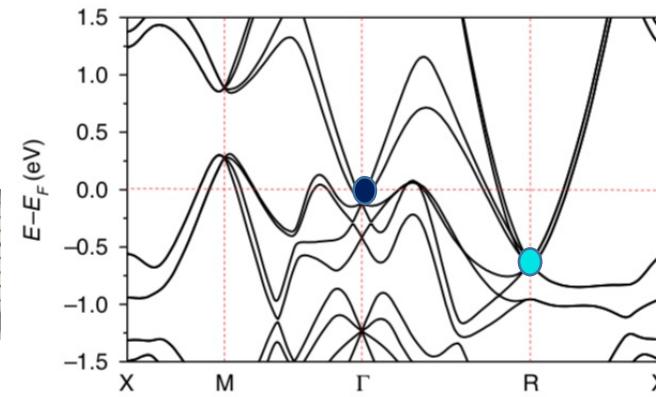
new fermions with chiral surface states



PdGa Chiral-1, Right handed



PdGa Chiral-2, Left handed

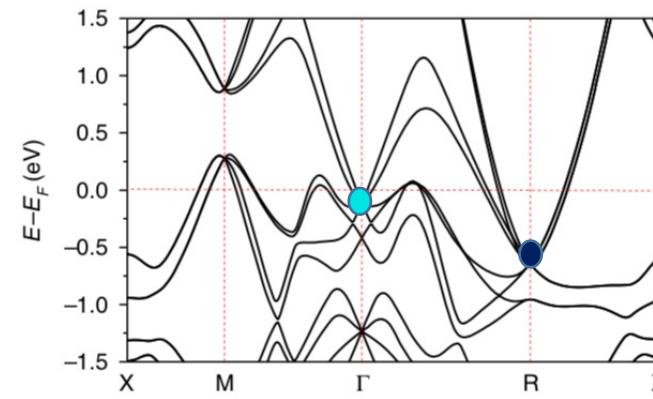


The crystal structure is **chiral**, the electronic structure is **identical**, however, the Weyl points and the Fermi arcs are **chiral**

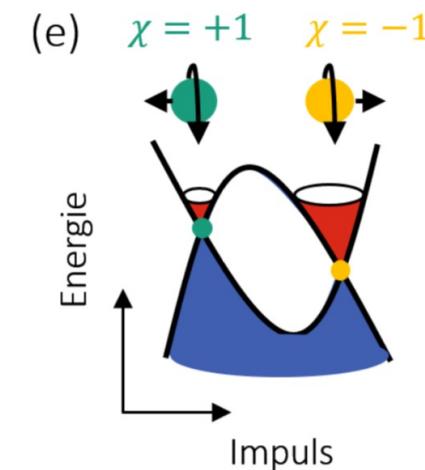
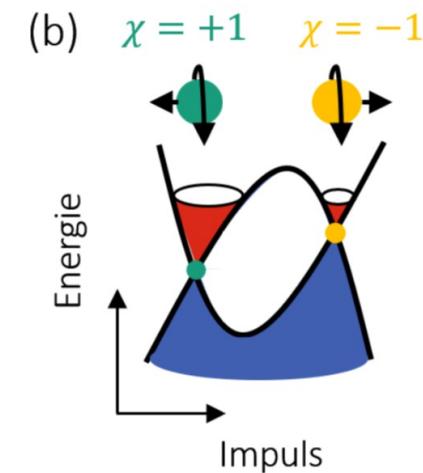
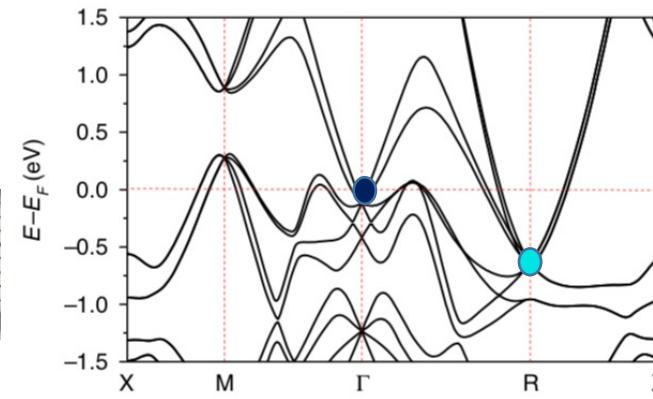
new fermions with chiral surface states



PdGa Chiral-1, Right handed



PdGa Chiral-2, Left handed

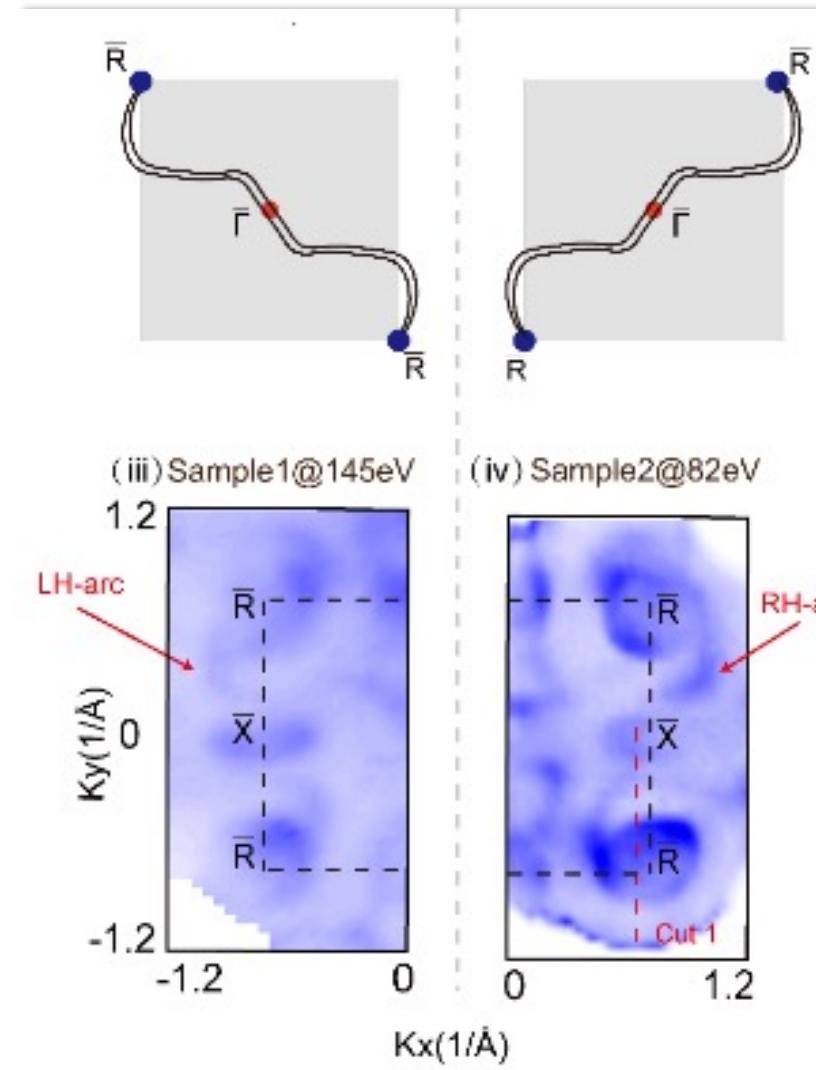
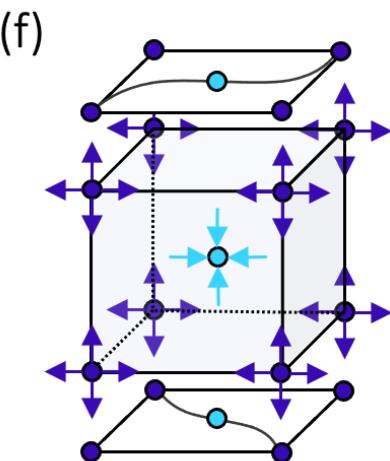
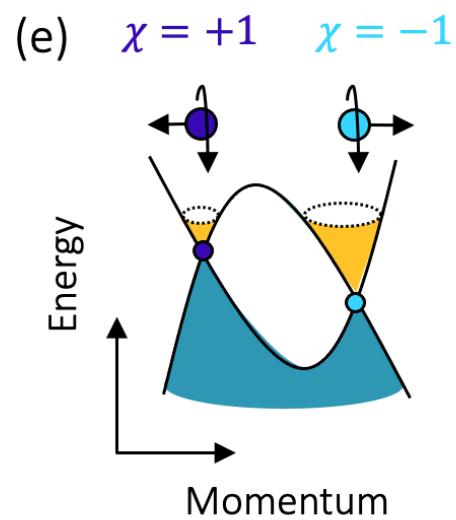
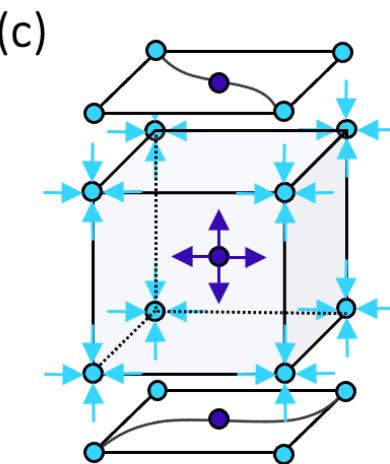
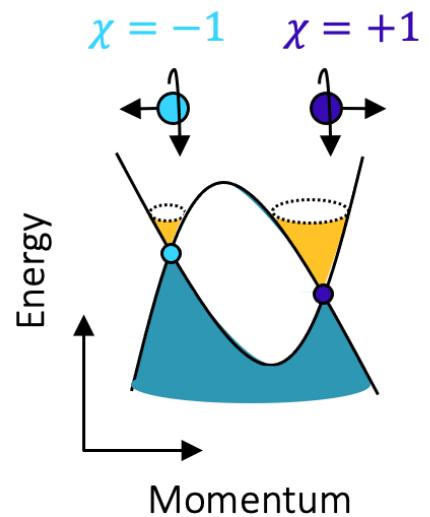


The crystal structure is **chiral**, the electronic structure is **identical**, however, the Weyl points and the Fermi arcs are **chiral**

chiral Fermions



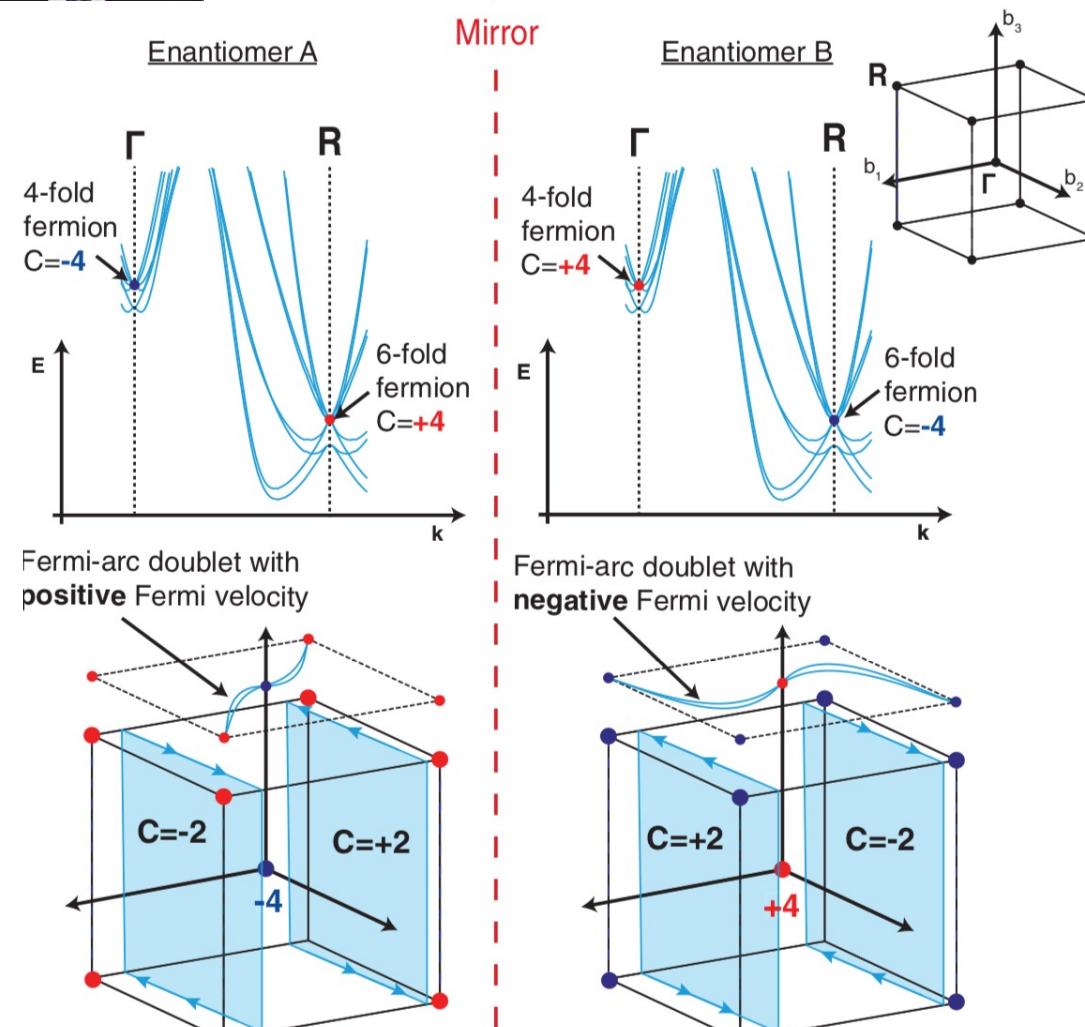
Weyl electrons and the Fermi arc **chiral**



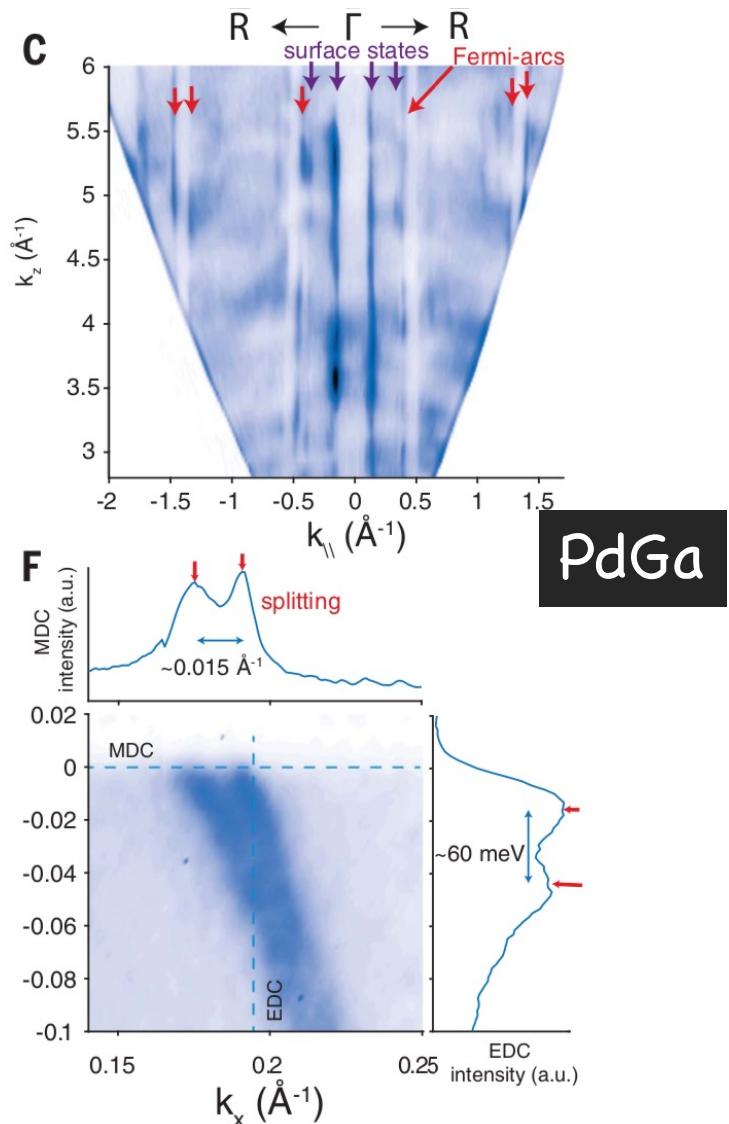
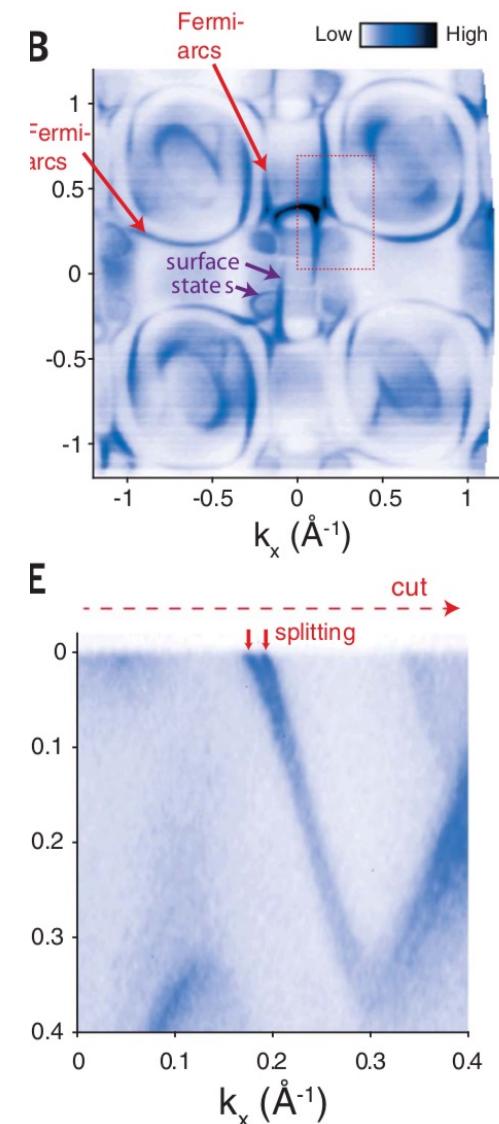
PdGa



new fermions with chiral surface states

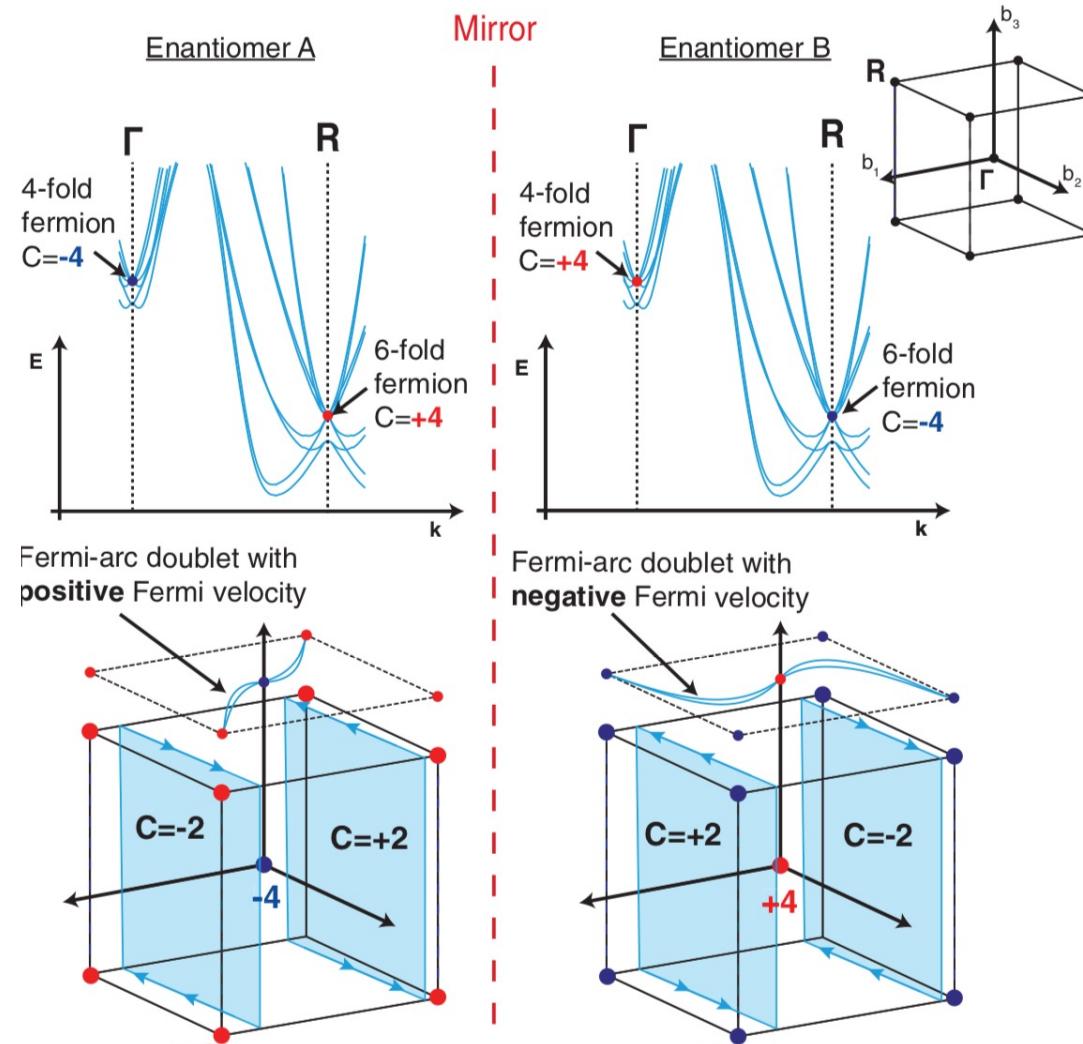


Schröter *et al.*, Science 369, 179–183 (2020)

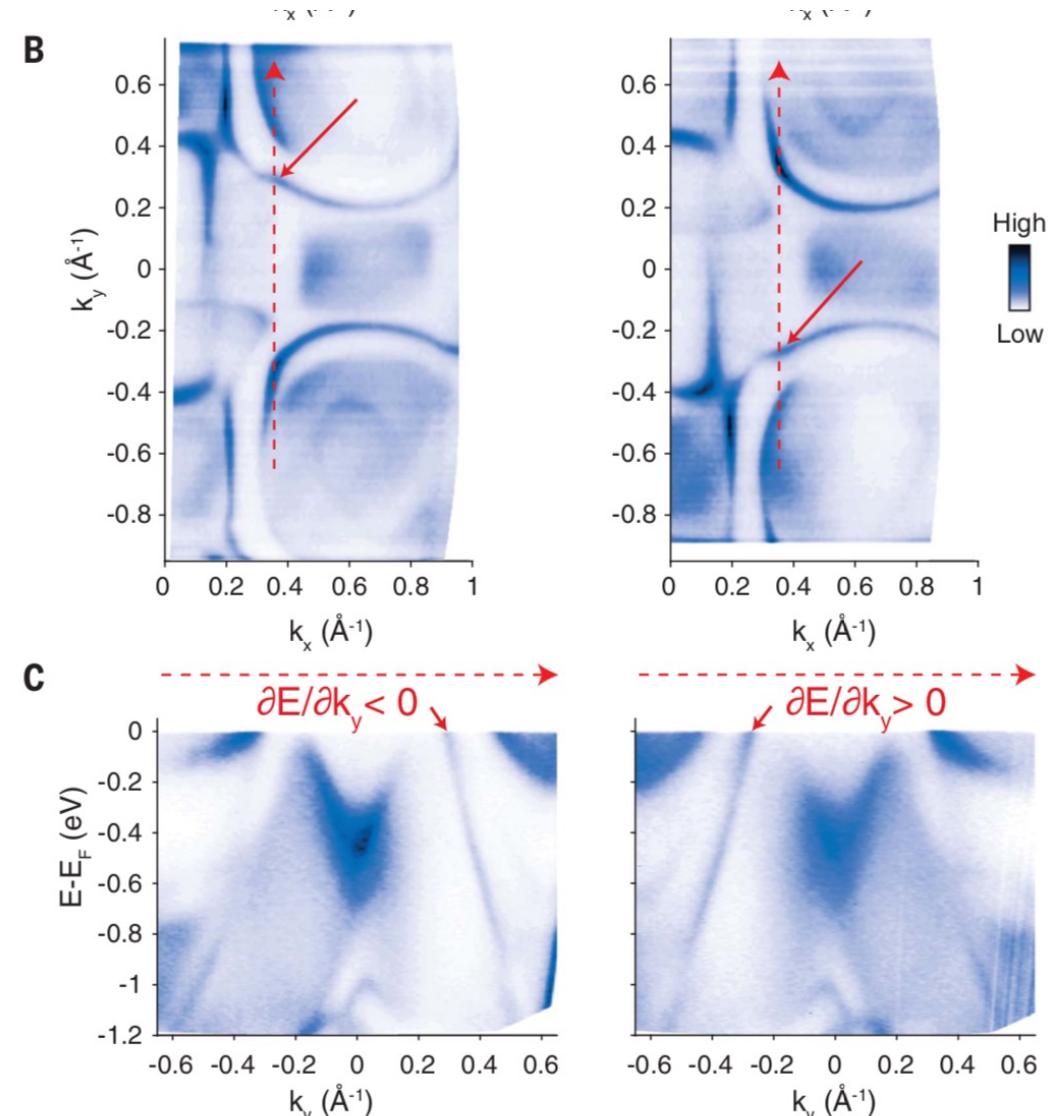




new fermions with chiral surface states



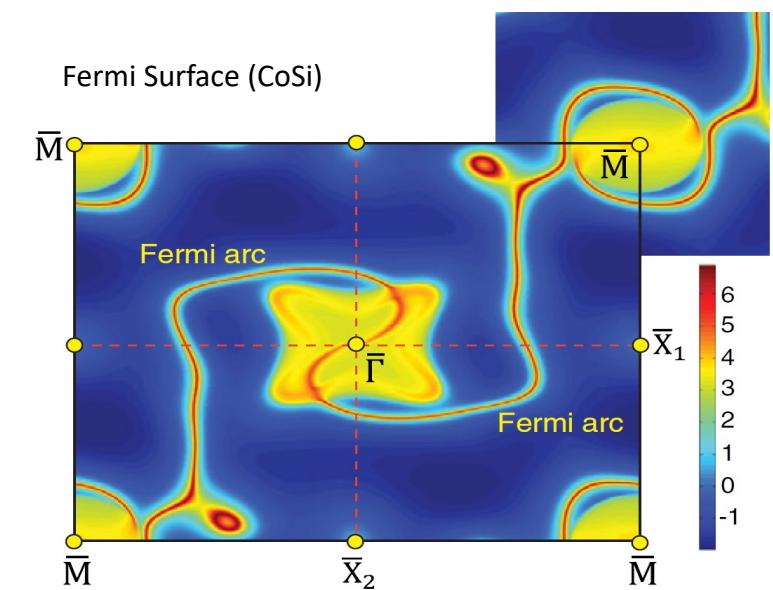
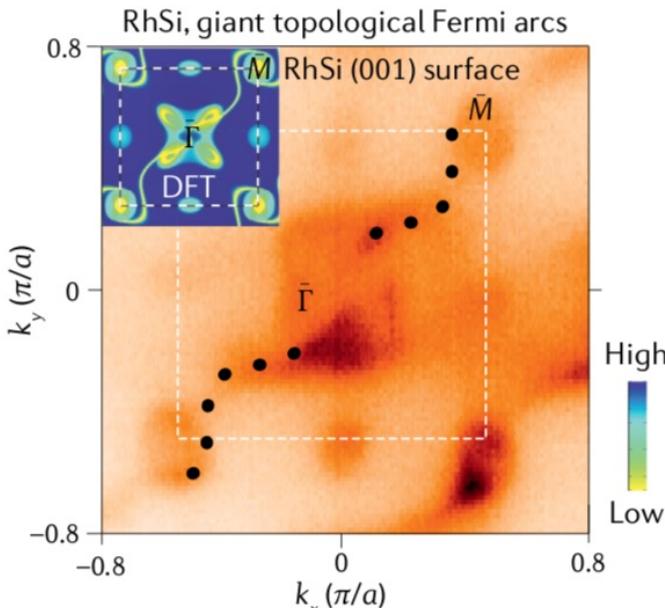
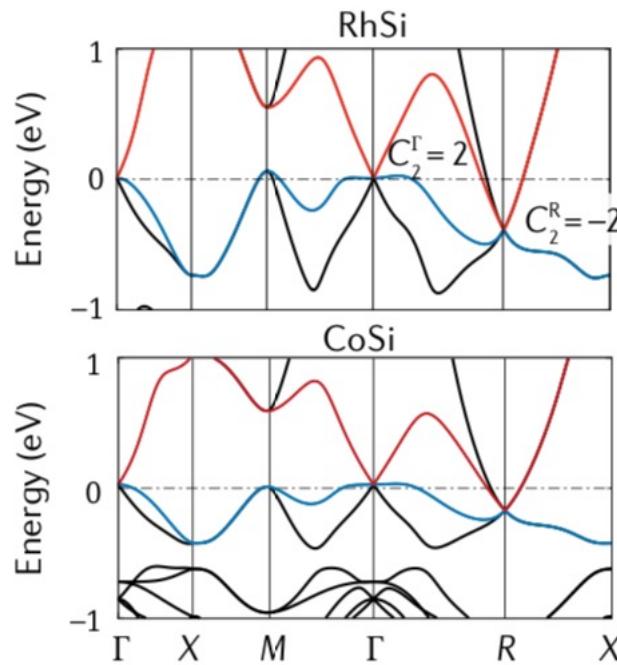
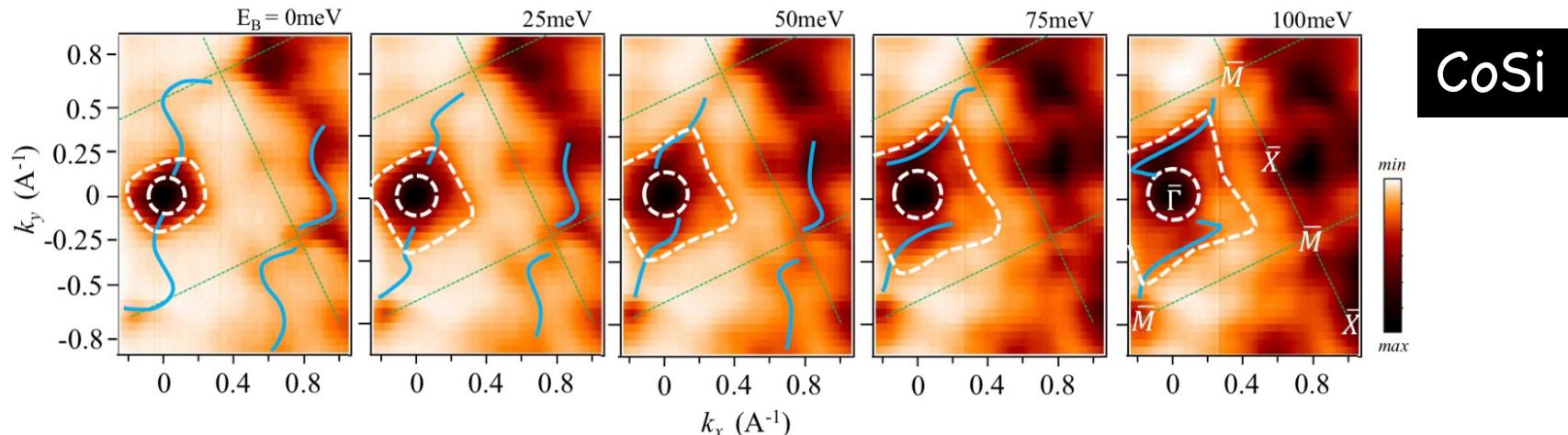
Schröter *et al.*, Science 369, 179–183 (2020)



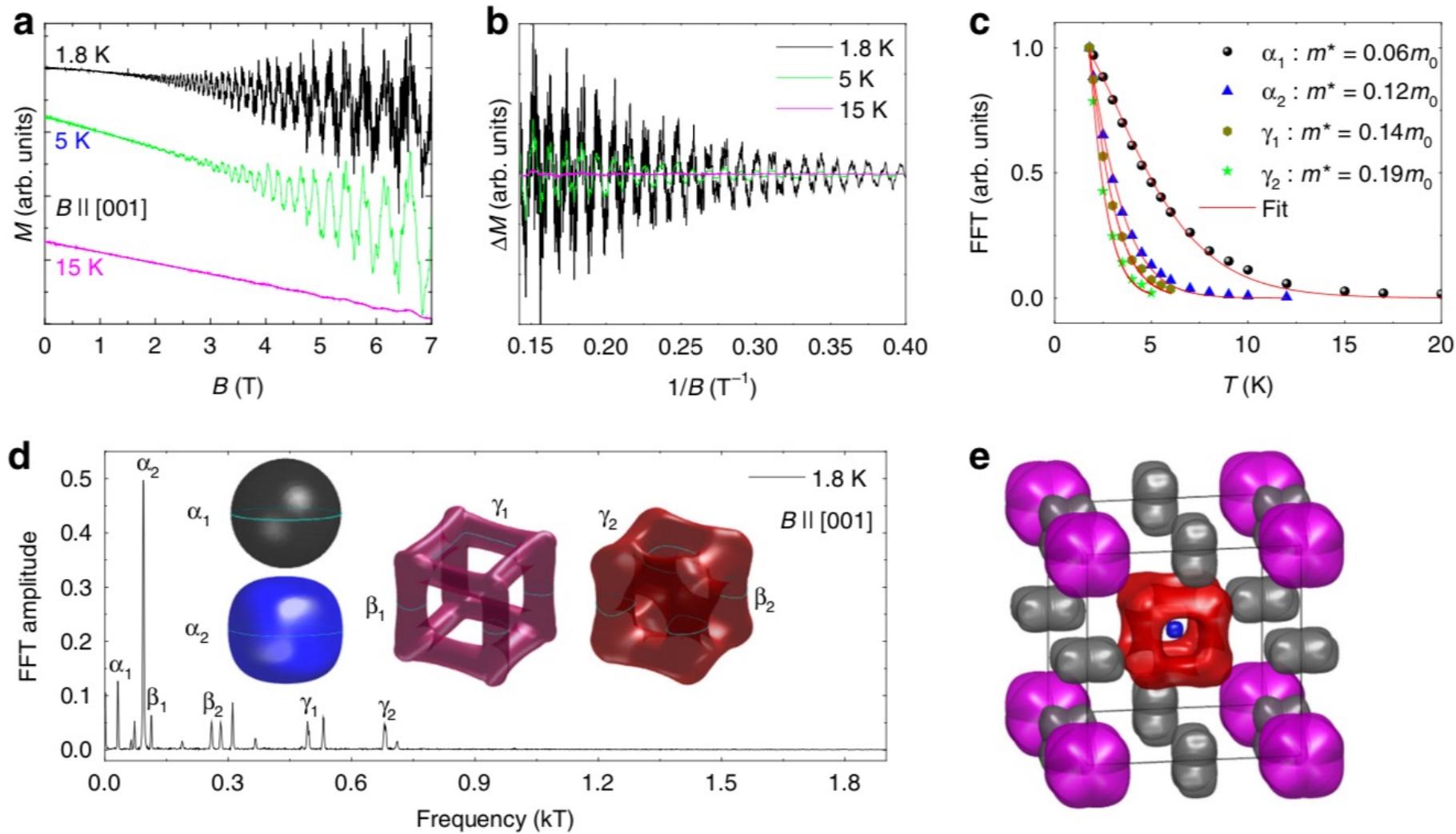


new fermions with chiral surface states

RhSi

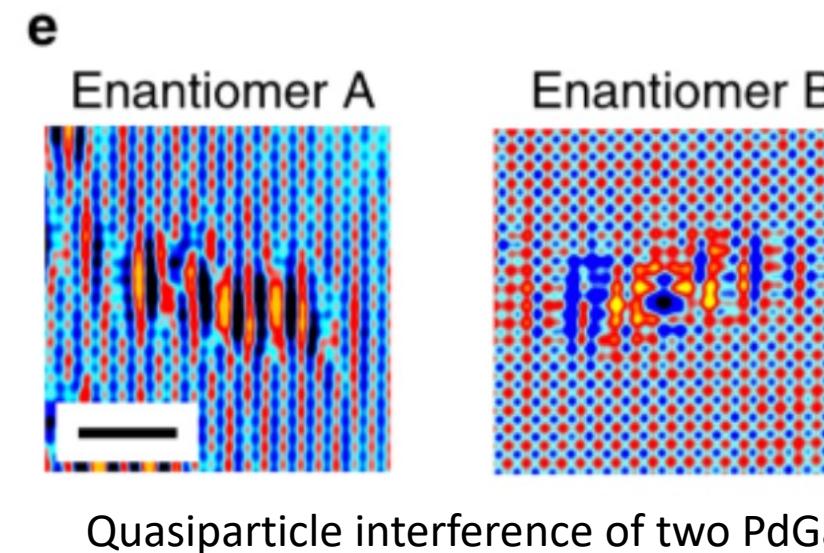
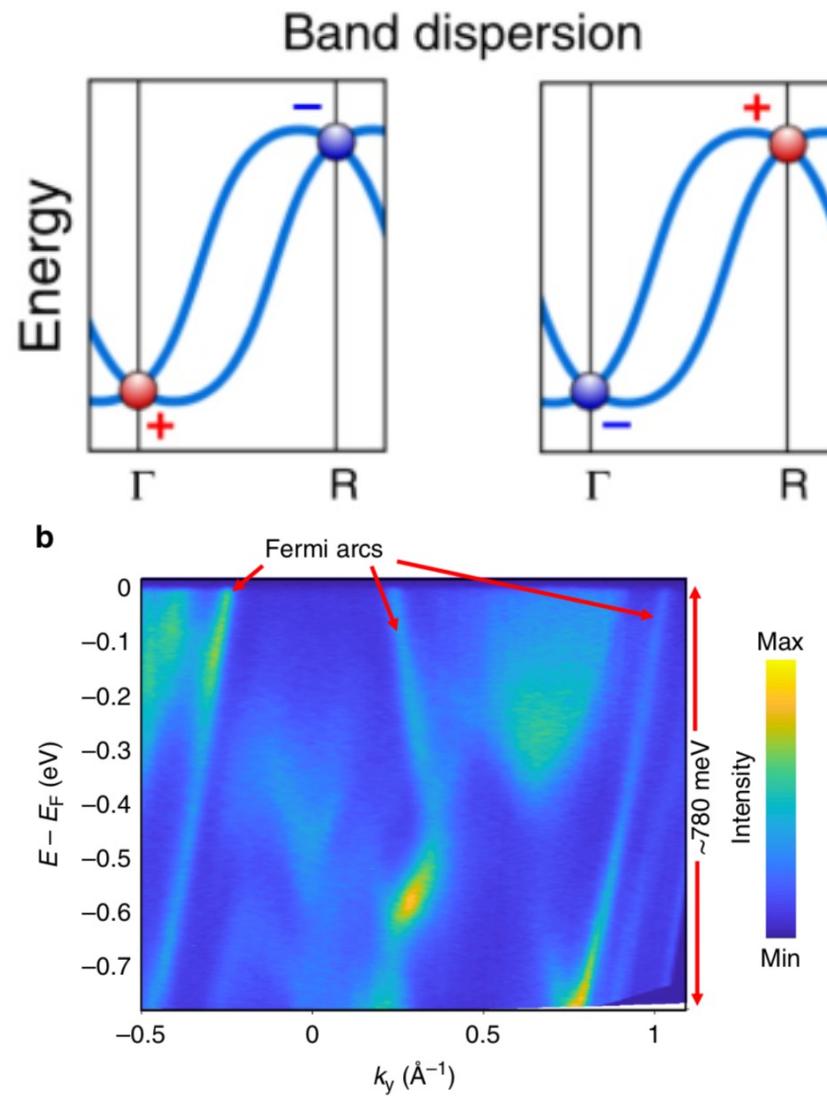


quantum oscillations of PtGa



PtGa

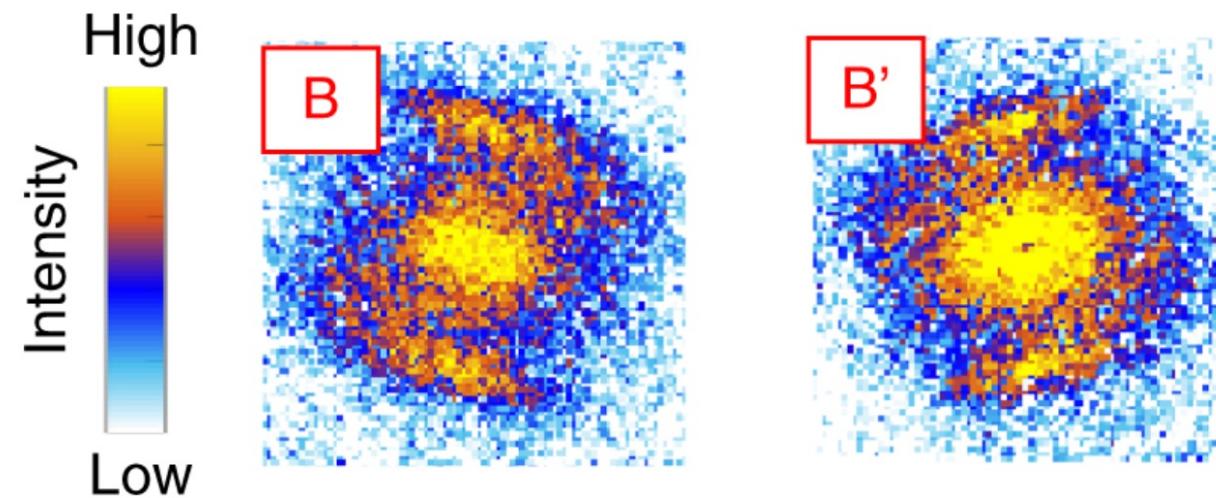
chiral surface states with STM



From STM investigations:

- (i) the perturbation developing around native defects is chiral
- (ii) the scattering vector associated with scattering events between opposite Fermi arcs is also chiral

Quasiparticle interference of two PdGa(001) enantiomers

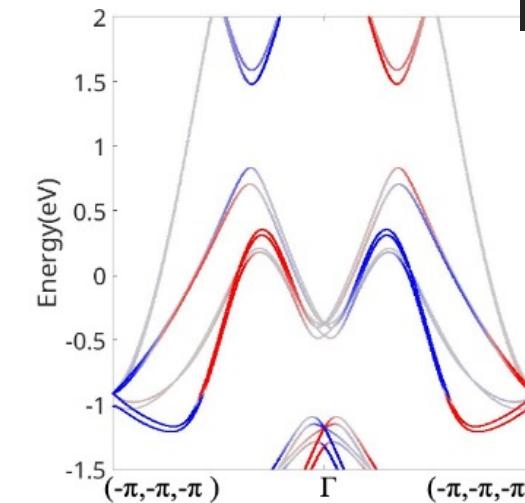
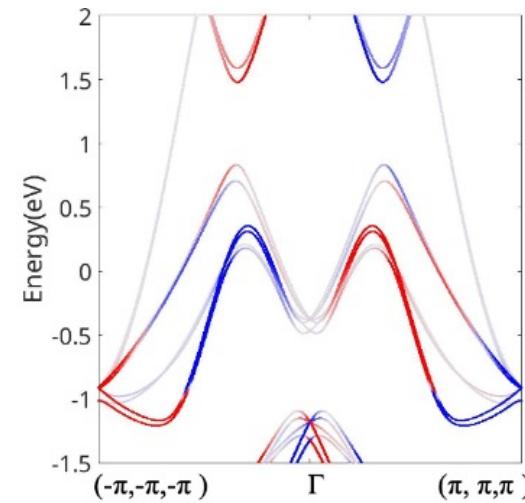
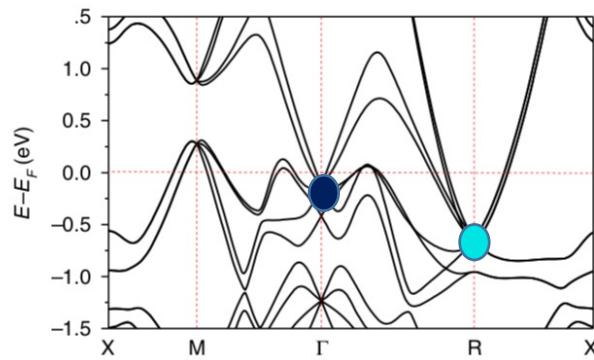
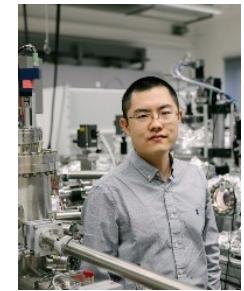
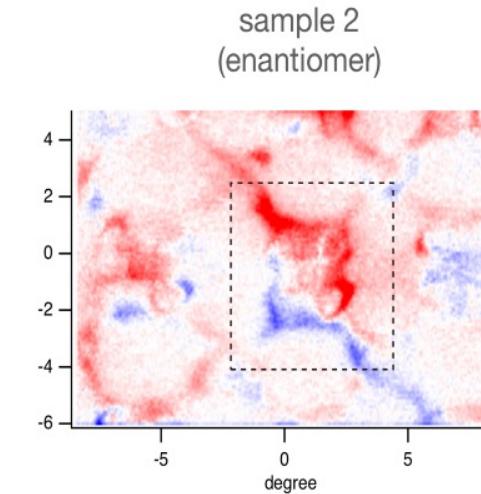
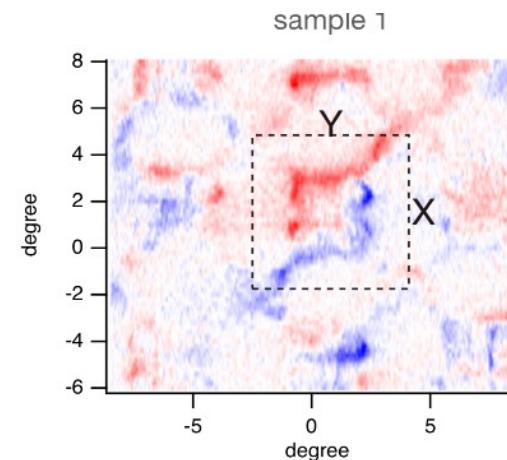
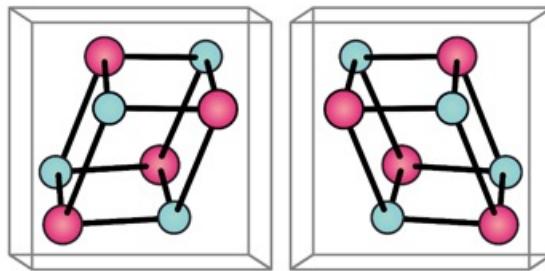


PdGa

Quasiparticle interference of two PdGa(001) enantiomers

chiral electrons in the bulk

With circular polarized light we can visualize the difference in the band structure



PdGa



Orbital angular momentum, <111>



Quantized circular photovoltaic effect

ARTICLE

Received 27 Dec 2016 | Accepted 18 May 2017 | Published 6 Jul 2017

DOI: 10.1038/ncomms15995

OPEN

Quantized circular photovoltaic effect in Weyl semimetals

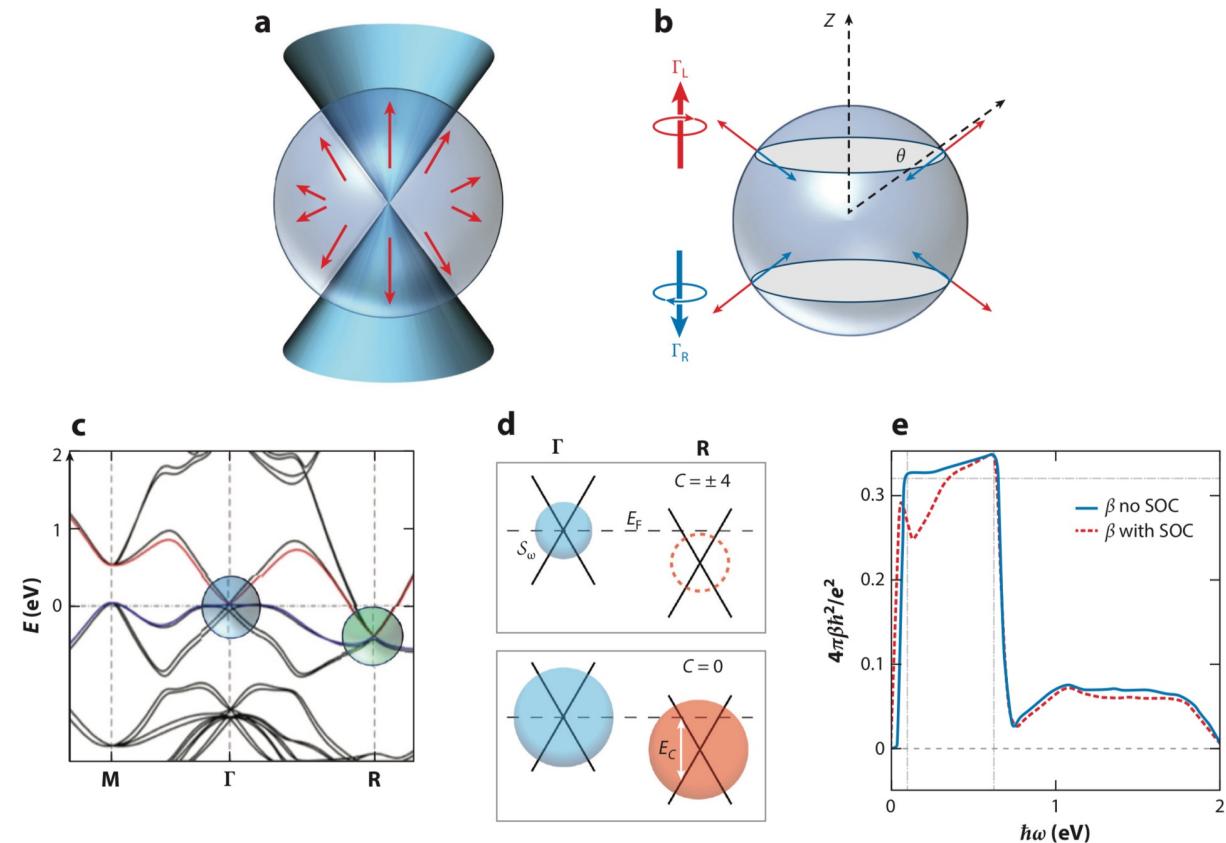
Fernando de Juan^{1,2,3}, Adolfo G. Grushin¹, Takahiro Morimoto¹ & Joel E. Moore^{1,4}

Prediction: Excitation of Weyl fermions –
a current that is quantized in units of material-independent
fundamental constants over a range of photon energies

The rate of change of the difference in photocurrent
generated by left-circularly and right-circularly polarized light,
 dj/dt , is quantized to the **Chern number**

Experiment: Frequency-independent plateau at low photon
energy abruptly falls-off above 0.66 eV

RhSi





Quantized circular photogalvanic effect

ARTICLE

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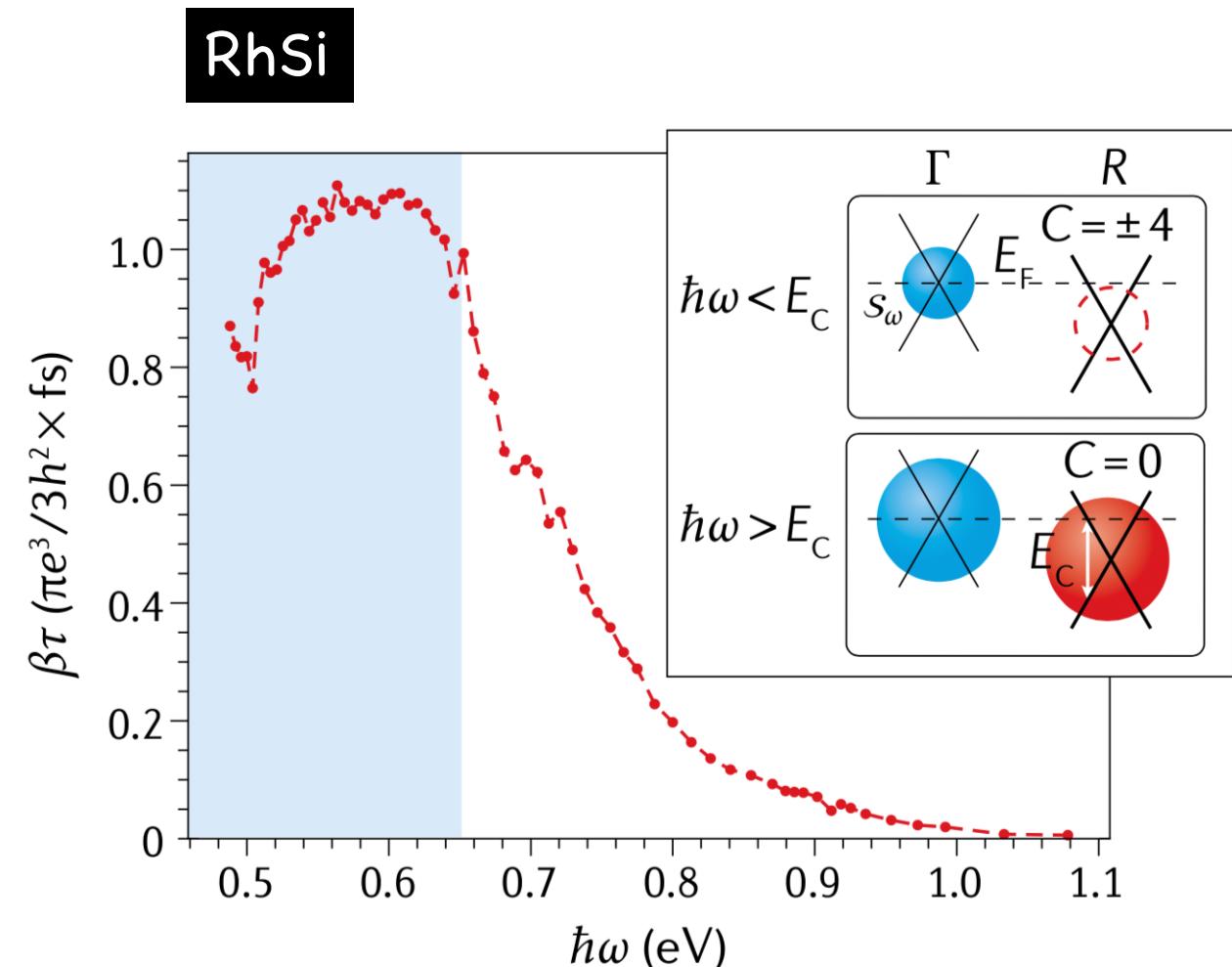
Quantized circular photogalvanic effect in Weyl semimetals

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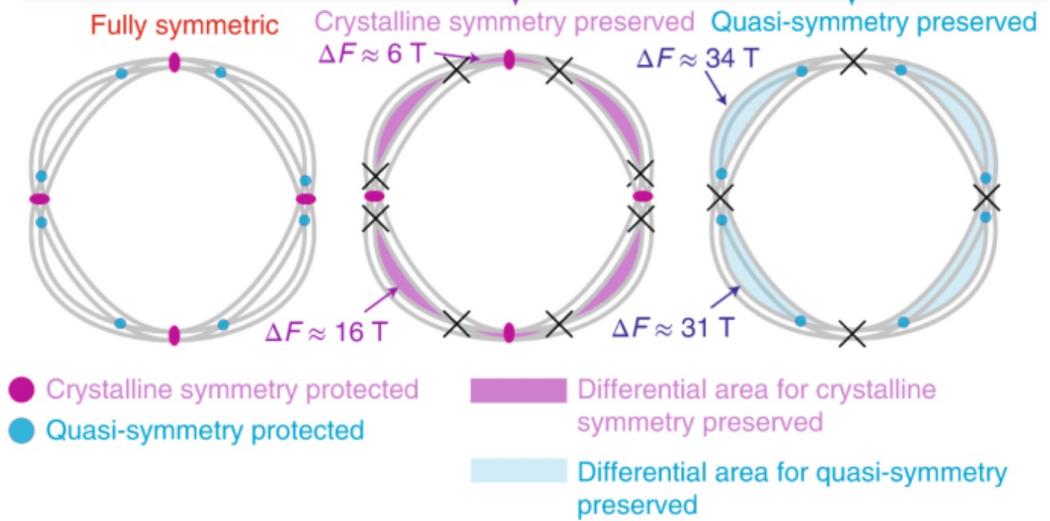
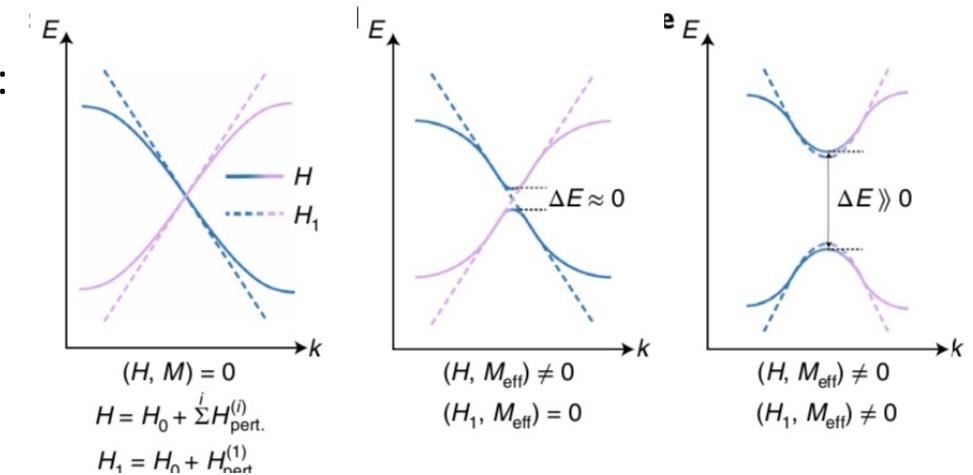
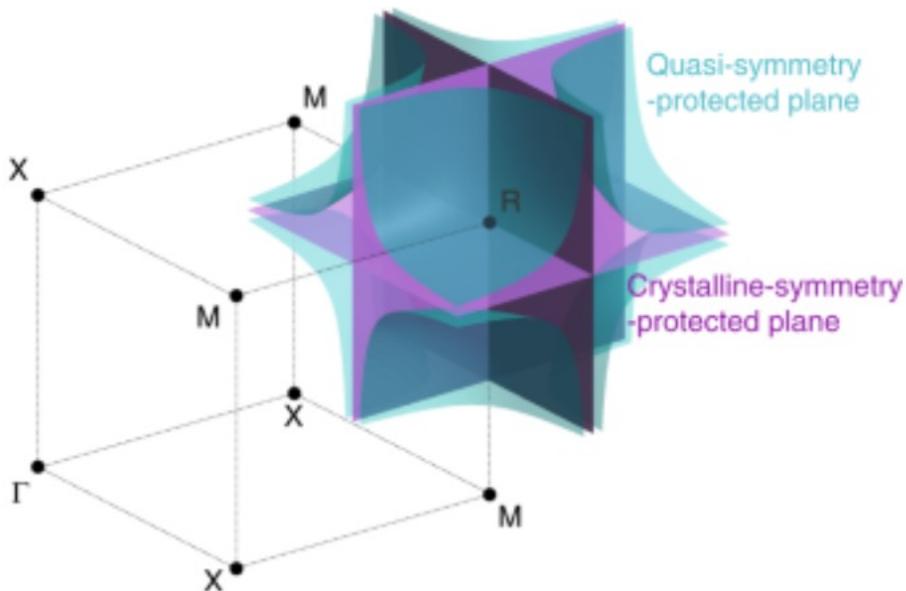


Quasi-symmetry-protected topology



a new concept of hidden or quasi symmetries:
avoided crossings with small gaps but large
Berry curvature in the electronic bands of the
topological semimetal CoSi.

CoSi



chiral surface states in CoSi

RhSi and CoSi: with multiple helicoid arc saddle points: type-I and type-II van Hove singularities.

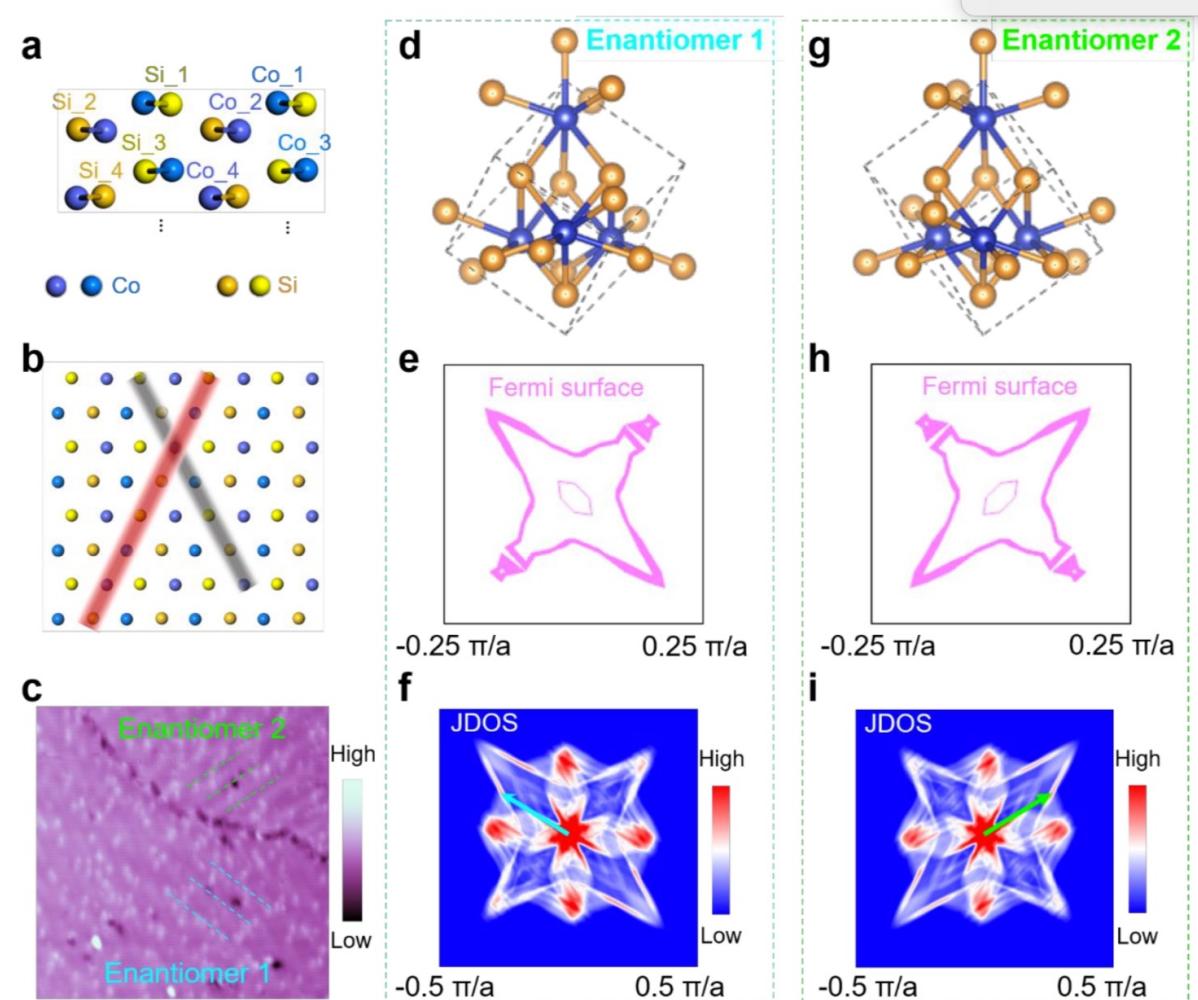
CoSi

Charge instability of topological Fermi arcs in chiral crystal CoSi

Zhicheng Rao^{1,2,†}, Quanxin Hu^{1,2,†}, Shangjie Tian^{3,†}, Shunye Gao^{1,2}, Zhenyu Yuan^{1,2}, Cenyao Tang^{1,2}, Wenhui Fan^{1,2}, Jierui Huang^{1,2}, Yaobo Huang⁴, Li Wang⁵, Lu Zhang^{1,2}, Fangsen Li⁵, Huaixin Yang^{1,2,6}, Hongming Weng^{1,2,6}, Tian Qian^{1,2,6}, Jinpeng Xu^{1,2,7*}, Kun Jiang^{1,2}, Hechang Lei^{3*}, Yu-Jie Sun^{8,1*} and Hong Ding^{1,6,7}

Chirality locking charge density waves in a chiral crystal

Geng Li^{1,2,3,4#}, Haitao Yang^{1,2#}, Peijie Jiang^{1,2#}, Cong Wang^{5#}, Qiuzhen Cheng^{1,2}, Shangjie Tian⁵, Guangyuan Han^{1,2}, Chengmin Shen^{1,2}, Xiao Lin^{1,2}, Hechang Lei^{5*}, Wei Ji^{5*}, Ziqiang Wang^{6*}, Hong-Jun Gao^{1,2,3,4*}



chiral surface states in CoSi and RhSi

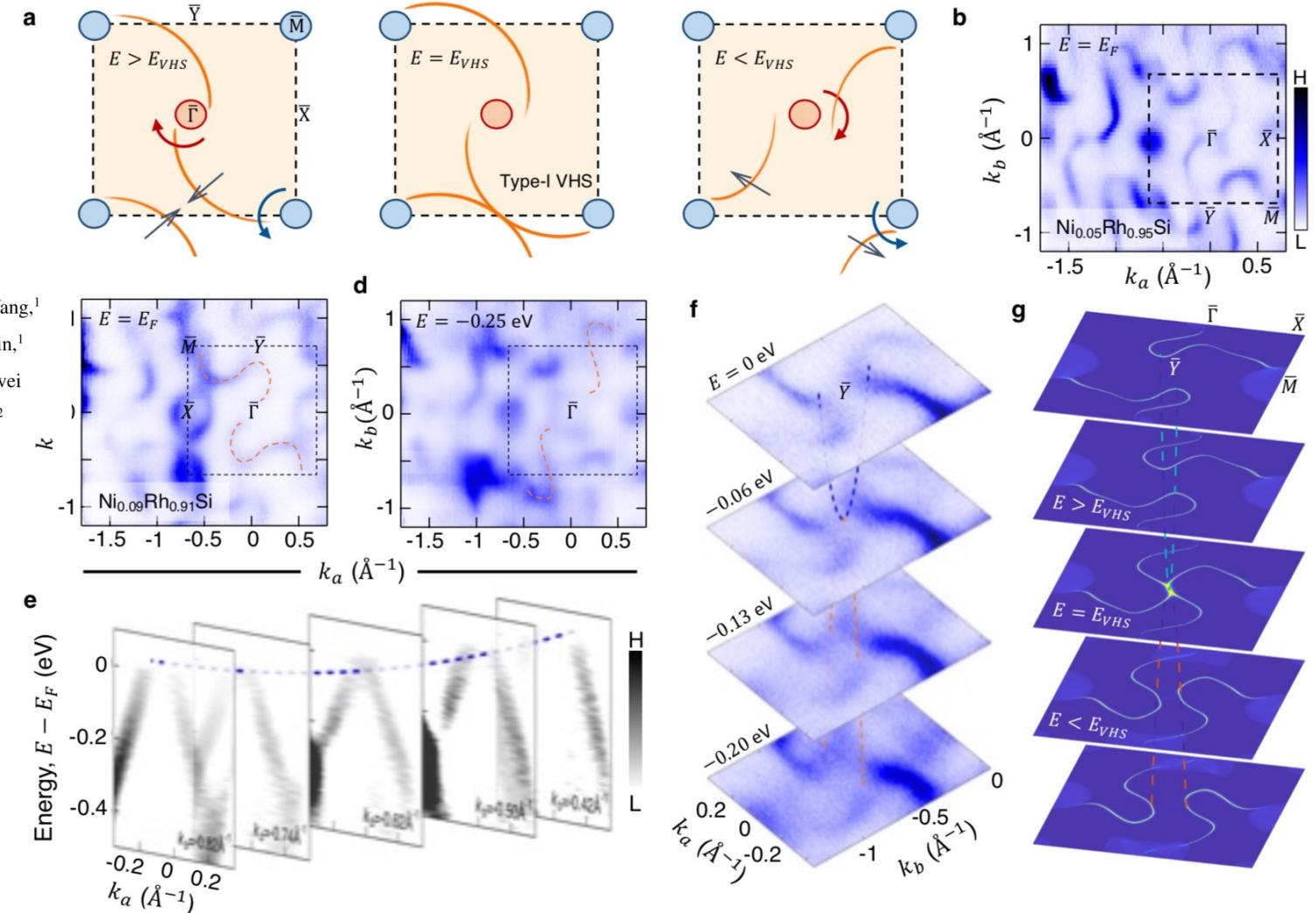
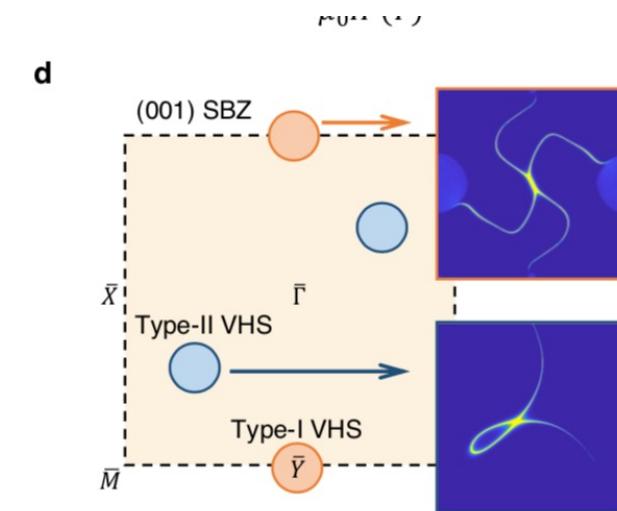
RhSi and CoSi: with multiple helicoid arc saddle points: type-I and type-II van Hove singularities.



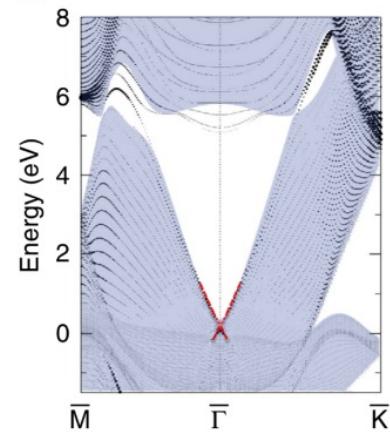
$\text{Ni}_x\text{Rh}_{1-x}\text{Si}$

Tunable topologically driven Fermi arc Van Hove singularities

Daniel S. Sanchez ^{*,1,2} Tyler A. Cochran ^{*,1} Ilya Belopolski, ^{1,3} Zi-Jia Cheng,¹ Xian P. Yang,¹ Yiyuan Liu,⁴ Tao Hou,⁵ Xitong Xu,⁴ Kaustuv Manna,^{6,7} Chandra Shekhar,⁶ Jia-Xin Yin,¹ Horst Borrmann,⁶ Alla Chikina,⁸ Jonathan D. Denlinger,⁹ Vladimir N. Strocov,⁸ Weiwei Xie,¹⁰ Claudia Felser,⁶ Shuang Jia,⁴ Guoqing Chang ^{†,5} and M. Zahid Hasan ^{†,11,12}

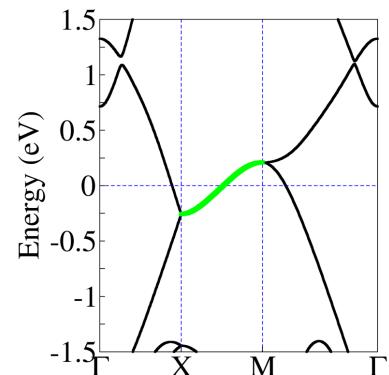


new catalyst with chiral surface states



Surface State of Platinum

B. Yan, et al. Nat. Commun. 2015, 6, 10167.

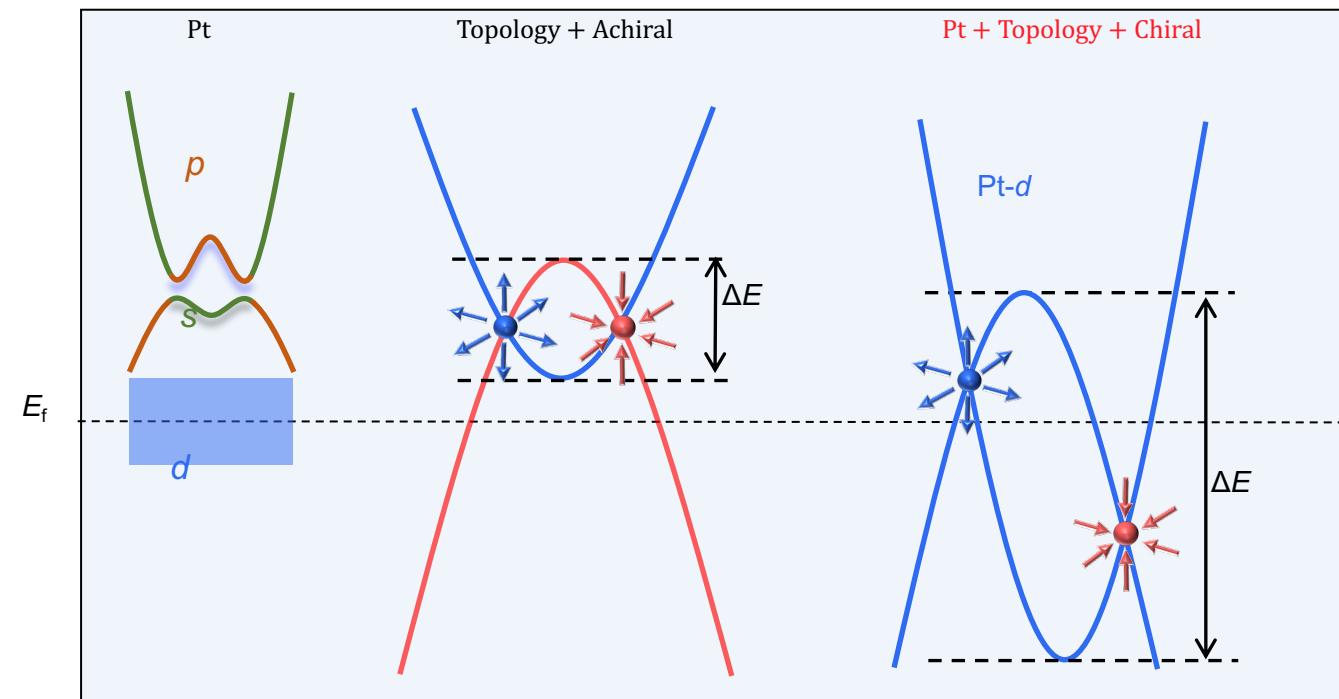
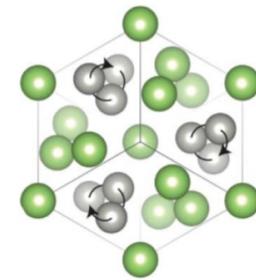


Nodal line in IrO_2

X. Xu et al. Phys. Rev. B 99 (2019) 195106

better than Pt for hydrogen evolution reaction (HER) and IrO_2 (OER, oxygen evolution reaction)

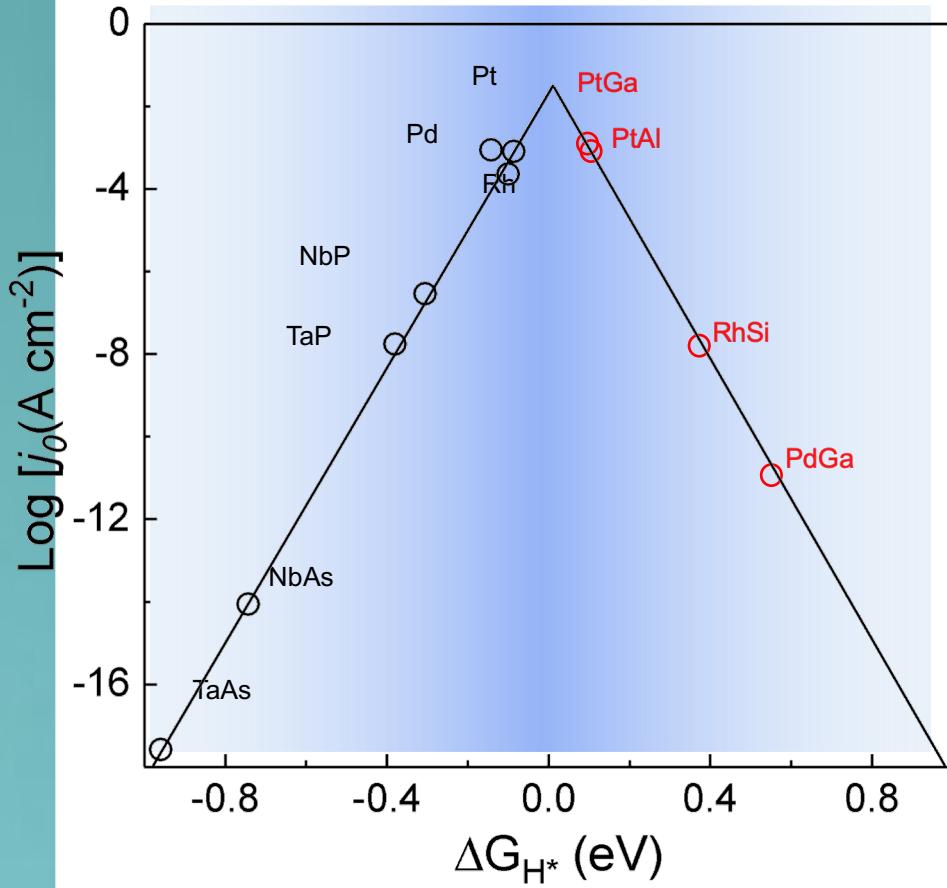
Pt and IrO_2 are topological relativistic effects and spin orbit coupling



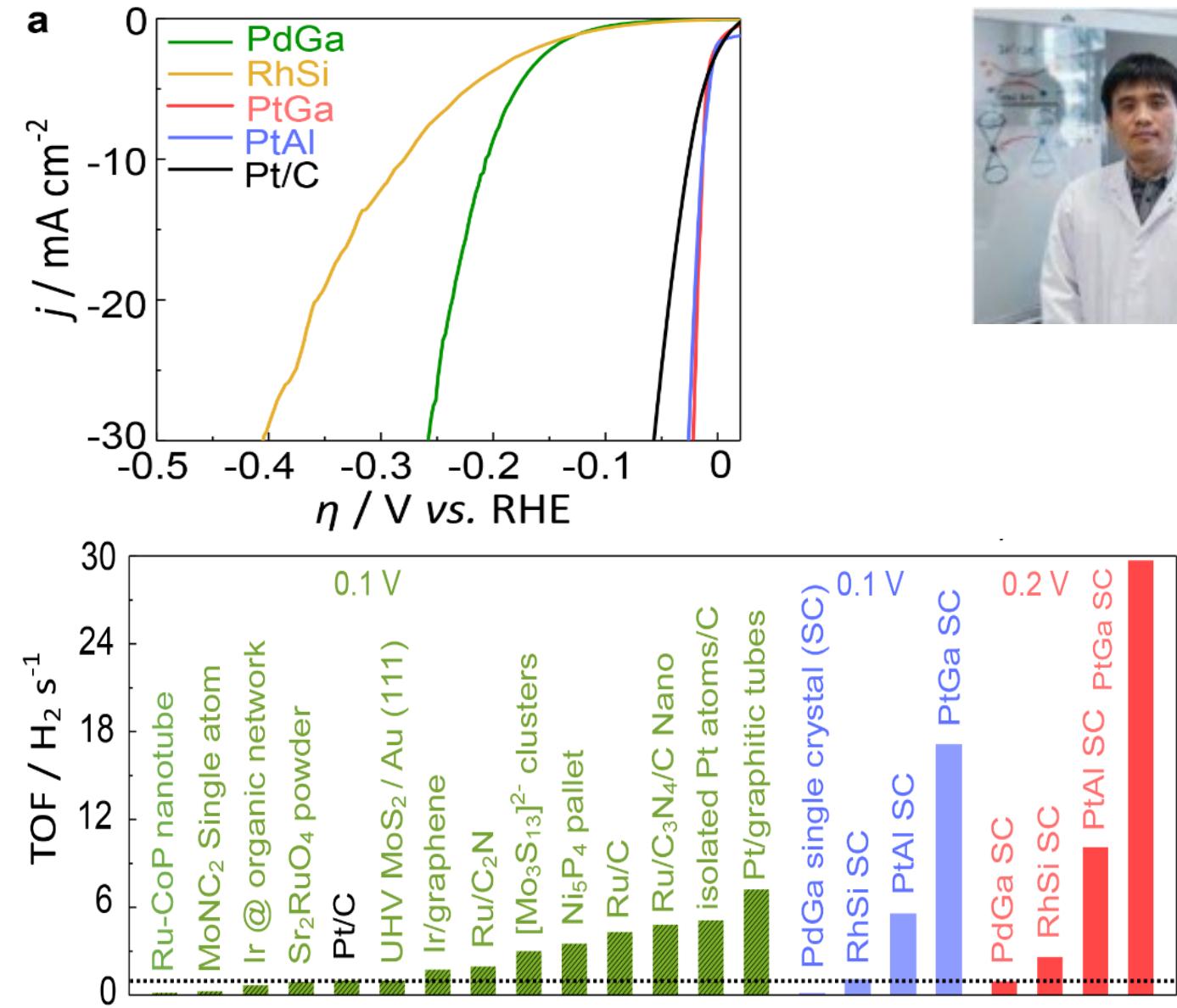
Q. Yang, et al, Adv. Mater., 2020, 32,1908518.

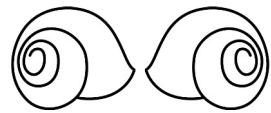
Chiral crystals for Water electrolysis and fuel cells, European patent, 19211719.0, submission 27.11.2019

new catalyst with chiral surface states

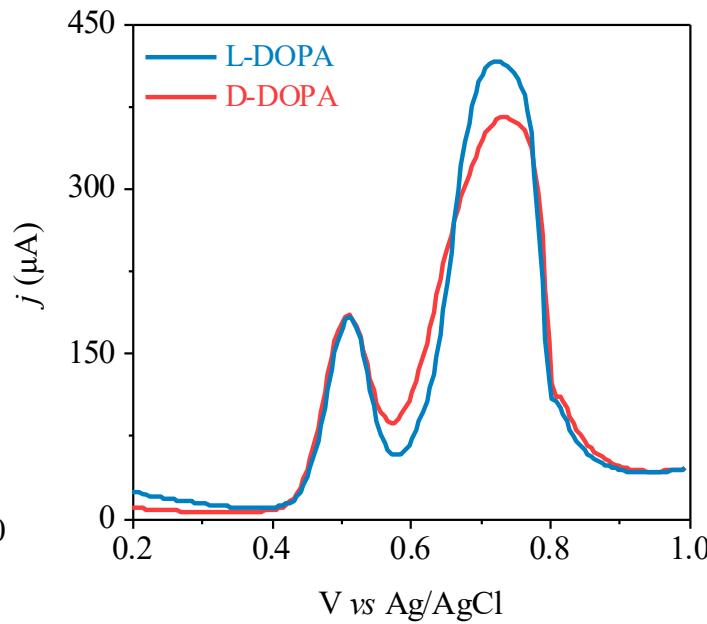
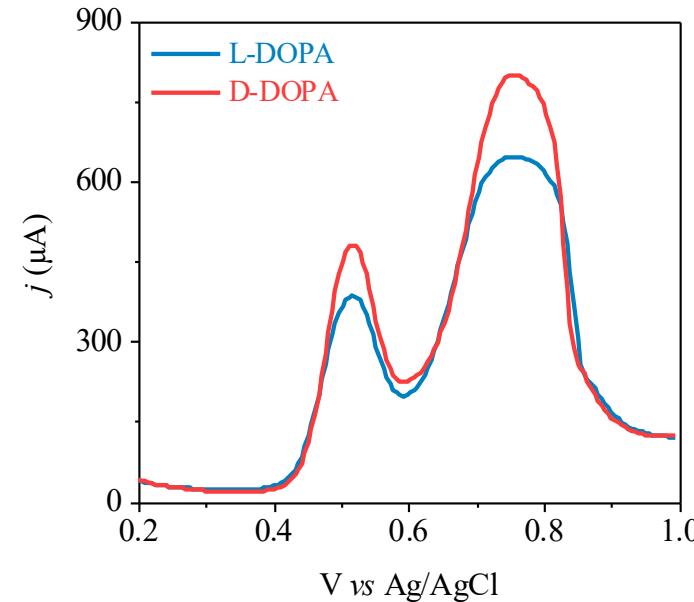
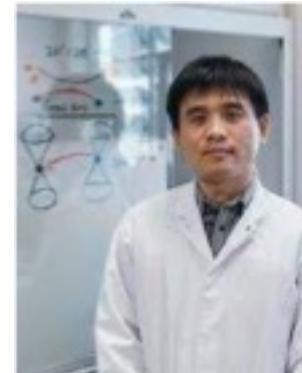
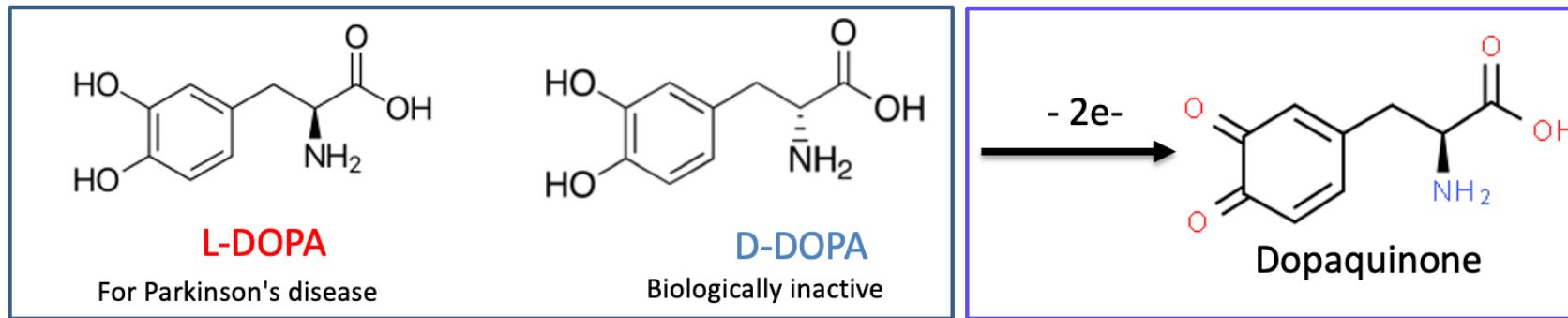


Qun Yang, et al. Advanced Materials 32 (2020) 1908518





absorption and oxidation of chiral molecules

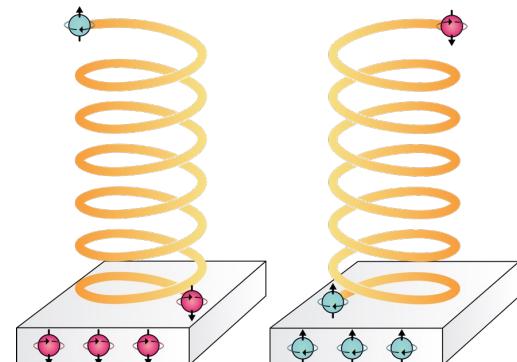
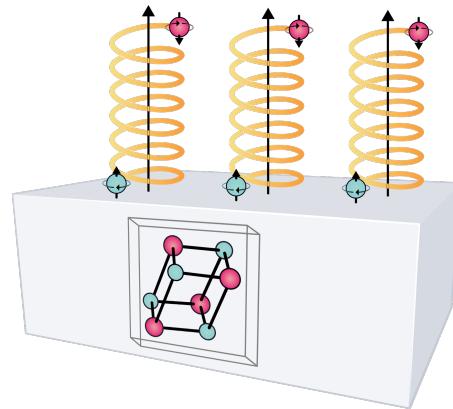
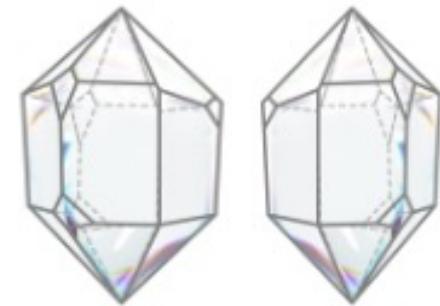
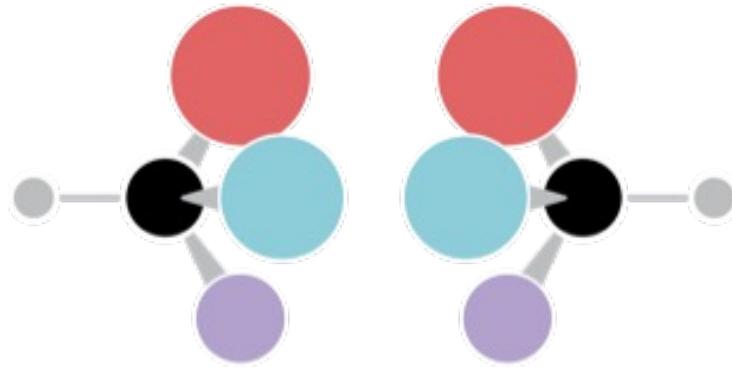
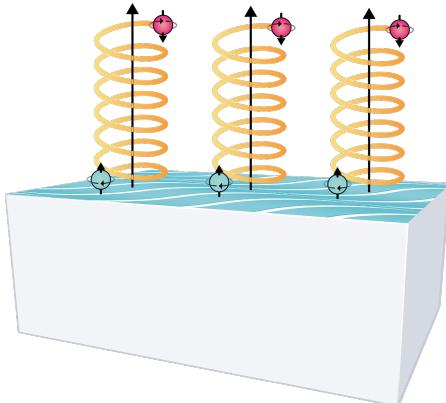
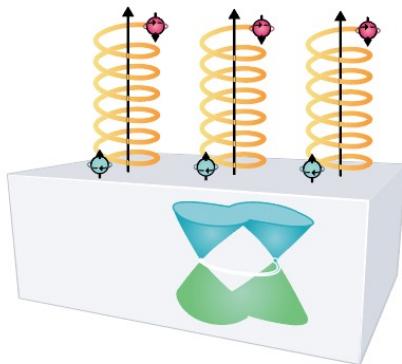
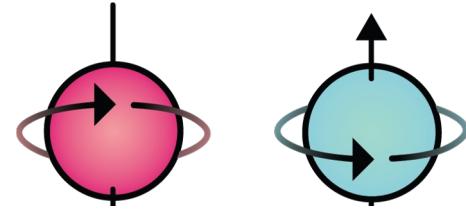


oxidation of Dopa on the surface of PdGa

all the measurement conditions are kept the same except the chirality of the PdGa crystal

D-DOPA and L-DOPA show different oxidation behaviors, depending on the chirality of PdGa

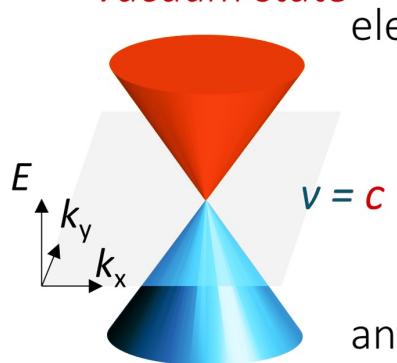
electrons – molecules, surfaces and crystals



physics meets chemistry

parity violation

Vacuum state

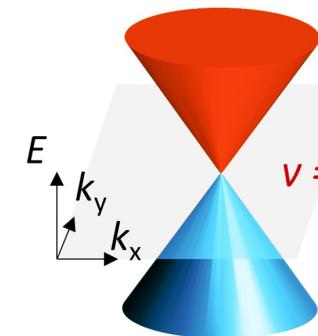


elementary particles

anti-particles

vacuum dispersion

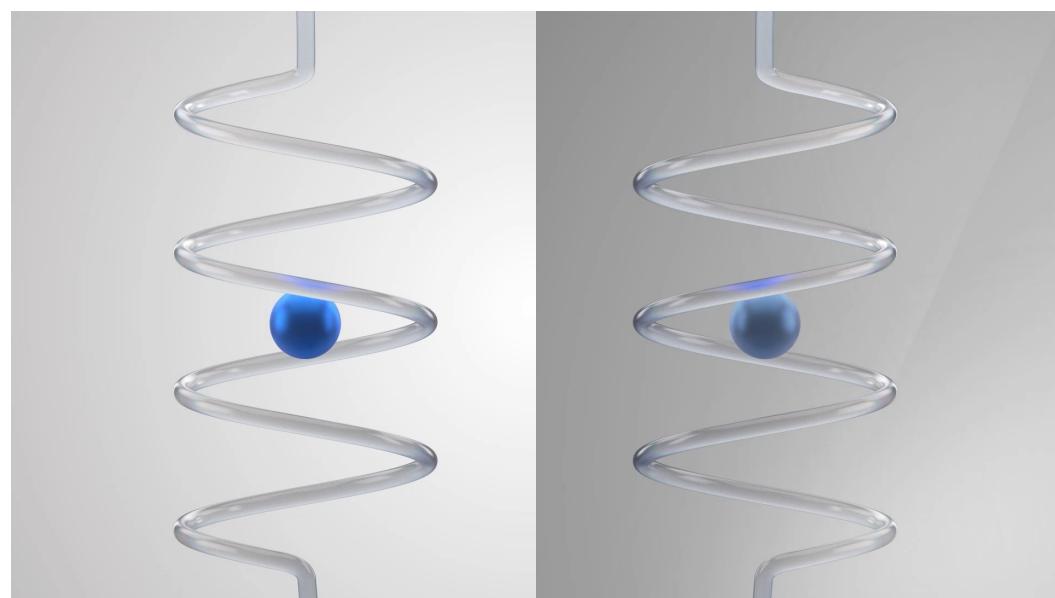
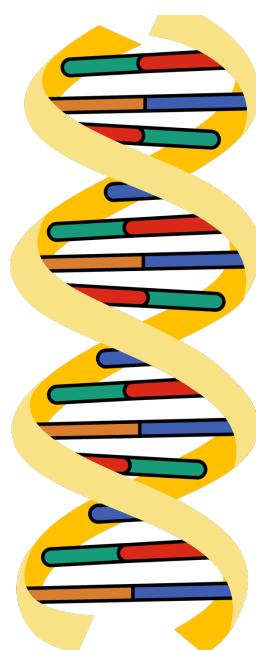
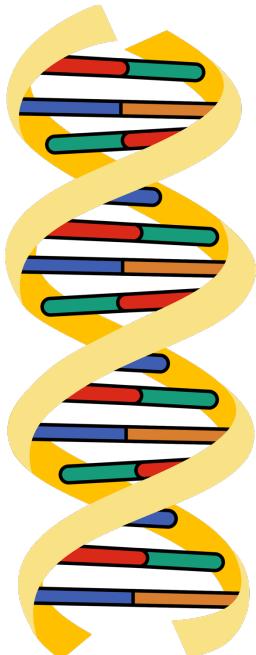
Dirac semimetal



'electrons'

'holes'

band structure



Wu experiment
β-decay

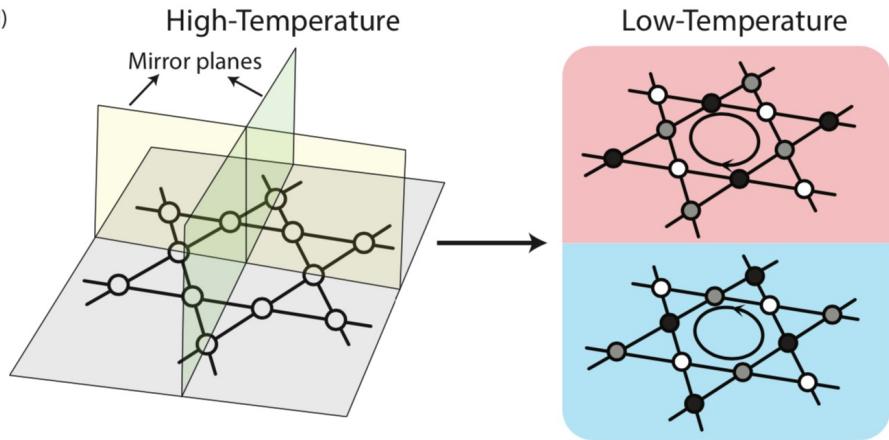


direction of
electrons

chirality in Charge density wave systems

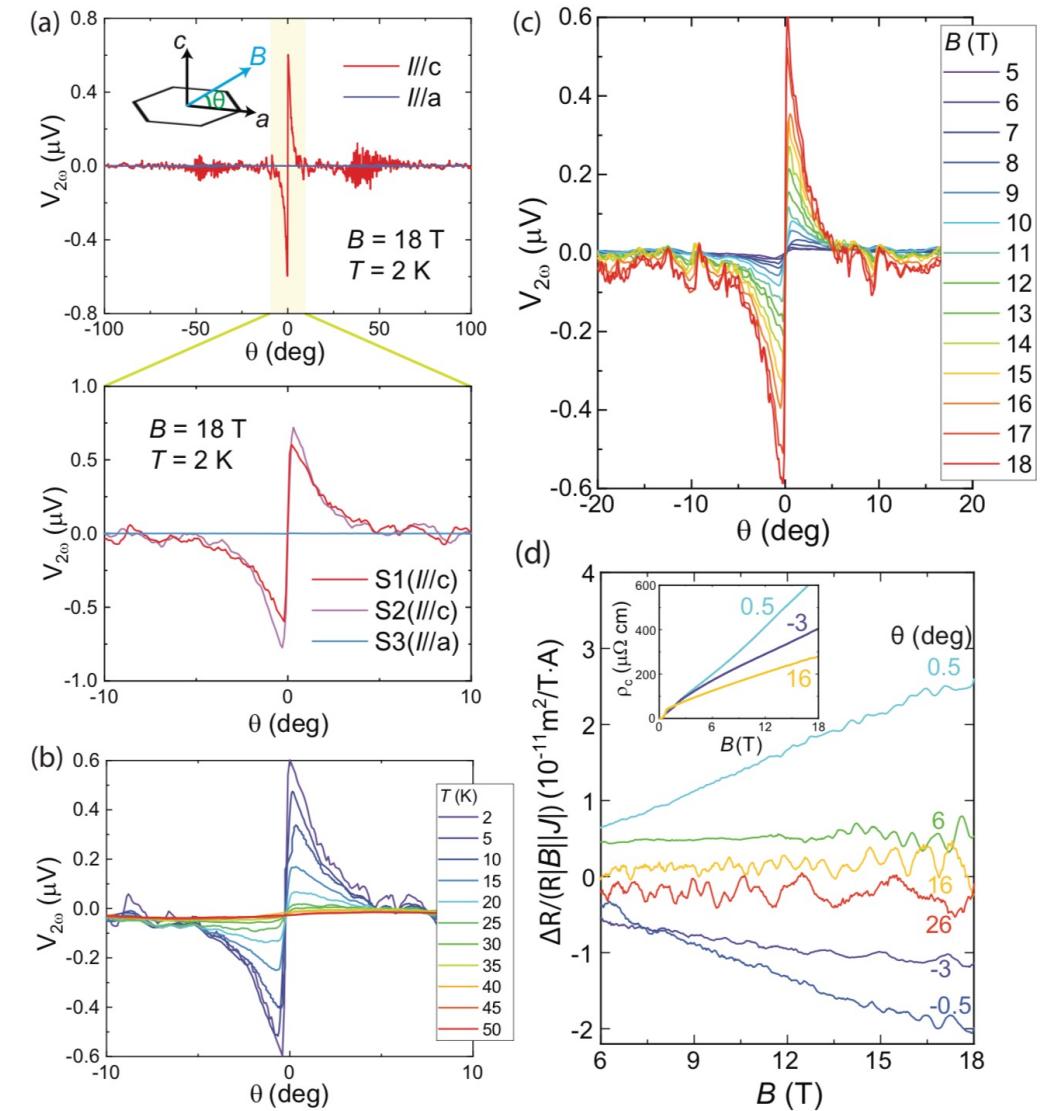


(d)

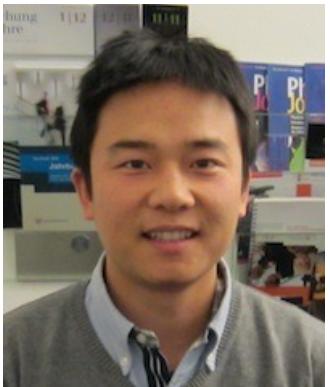


Switchable chiral transport in charge-ordered CsV_3Sb_5

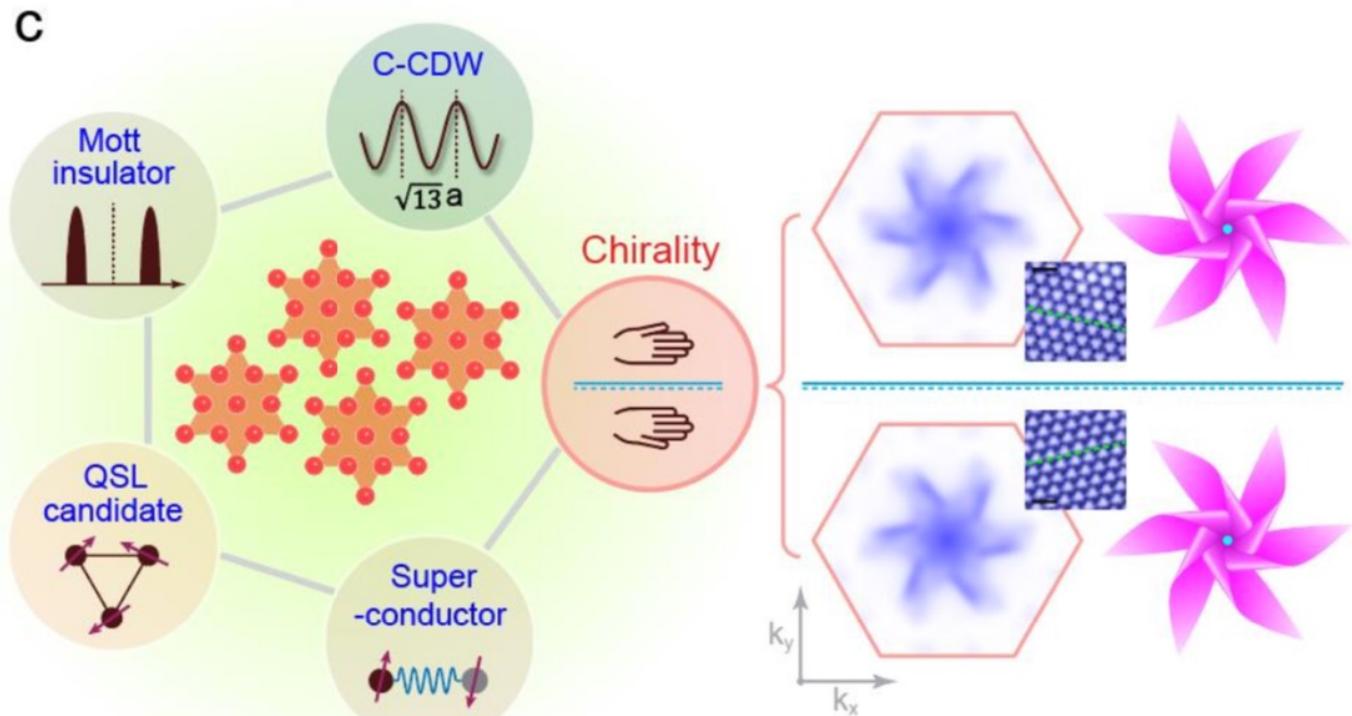
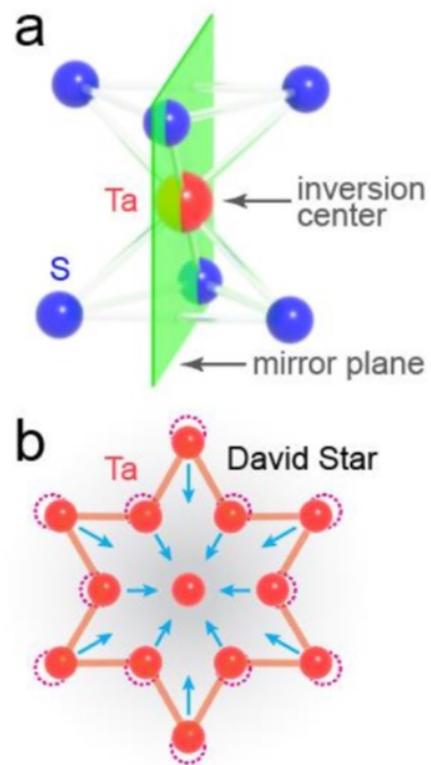
Chunyu Guo*,¹ Carsten Putzke,^{1,2} Sofia Konyzheva,¹ Xiangwei Huang,¹ Martin Gutierrez-Amigo,^{3,4} Ion Errea,^{3,5,6} Dong Chen,⁷ Maia G. Vergniory,^{5,7} Claudia Felser,⁷ Mark H. Fischer*,⁸ Titus Neupert[†],⁸ and Philip J. W. Moll^{‡,1,2}



chirality in Charge density wave systems



TaS₂

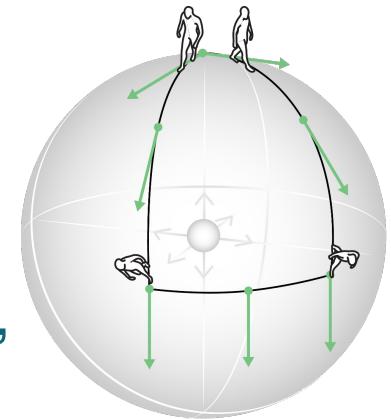


Visualization of Chiral Electronic Structure and
Anomalous Optical Response in a Material with Chiral
Charge Density Waves

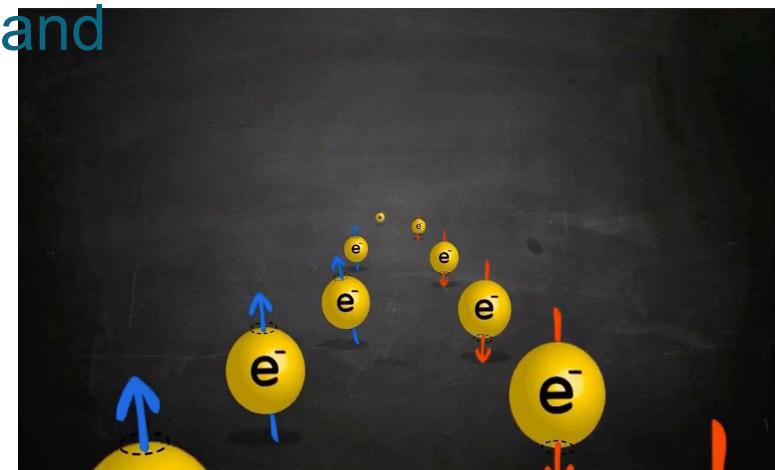
H. F. Yang^{1*}, K. Y. He^{2*}, J. Koo^{3*}, S. W. Shen¹, S. H. Zhang¹, G. Liu², Y. Z. Liu³, C. Chen^{1,4}, A. J. Liang^{1,5}, K. Huang¹, M. X. Wang^{1,5}, J. J. Gao⁶, X. Luo⁶, L. X. Yang⁷, J. P. Liu^{1,5}, Y. P. Sun^{6,8,9}, S. C. Yan^{1,5}, B. H. Yan^{3†}, Y. L. Chen^{1,4,5,7†}, X. Xi^{2,9†}, and Z. K. Liu^{1,5†}

summary

- topologically protected surfaces edge or edge states in crystals
- new quantum effects in crystals
- in semimetals one observes giant effects in response of magnetic, electric field, light etc.
- Fermi arcs
- arcs extended over the reciprocal space in chiral crystals
- high mobilities, free electron path lengths up to mm
- giant Nernst effect and magnetic Seebeck effect
- giant photovoltaic effect - quantized
- Weyl semimetals as model systems for high energy and astrophysics
- Parity violation E^*B in Weyl semimetals
- chiral anomaly
- axial gravitation anomaly
- topological catalysis



Berry Phase



summary

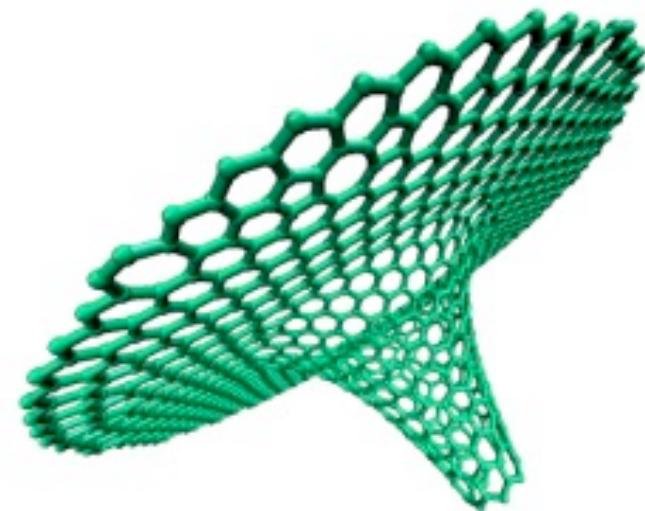
- more than **25%** of all inorganic compounds are topological
- quantum simulator for high energy and astro-physics
- paramagnetic Weyl semimetals, which break inversion: **chiral anomaly** in thermal and electrical transport
- chiral new Fermions: giant Fermi arcs, strong **Berry curvature effects**, new chiral optical effects

outlook

- chirality and magnetism
- **beyond the single particle picture** – topology in correlated materials
- experiments and understanding of **3D quantum Hall** effects
- connect chiral molecules and catalysis with Berry curvature
 - the interplay between chiral structure/surface state/orbital moment

potential applications

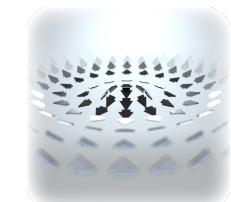
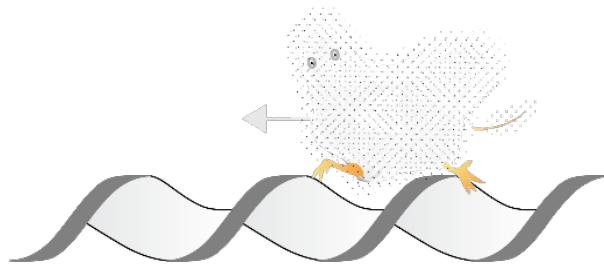
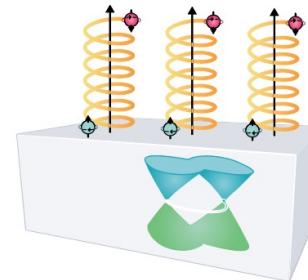
- energy conversion
 - catalysis ...
- spintronics
- quantum computing



Vision



- Crystal growth of both enantiomers of topological chiral compounds
 - interfaces, grain boundaries, chiral phonons, magnons, ...
 - Investigation of the interplay between structure, chiral surface state, orbital angular momentum, spin momentum locking ...
 - strain and magnetic field
- Chiral electrons, chiral Fermions, chiral surfaces and **catalysis**
 - enhanced light matter interaction and magnetic field
- Non local transport in chiral crystals
 - spin polarized currents over μm
- Chirality plus magnetism, correlations, superconductivity
 - Skyrmiions, Antiskyrmions ...



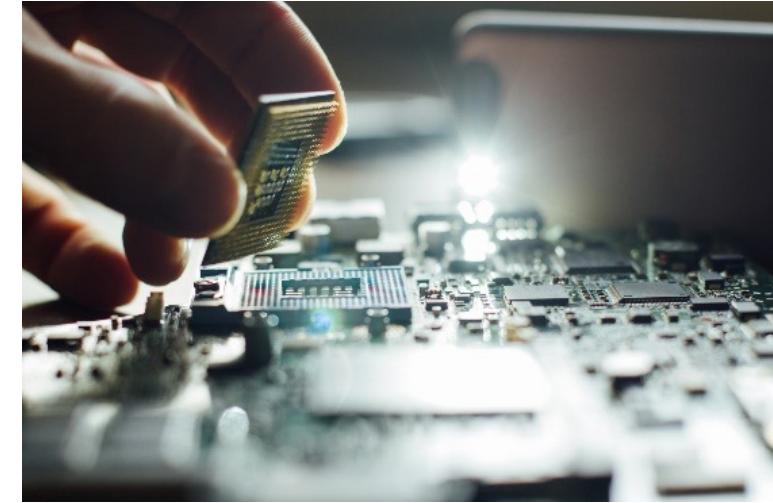
Potentielle Anwendungen



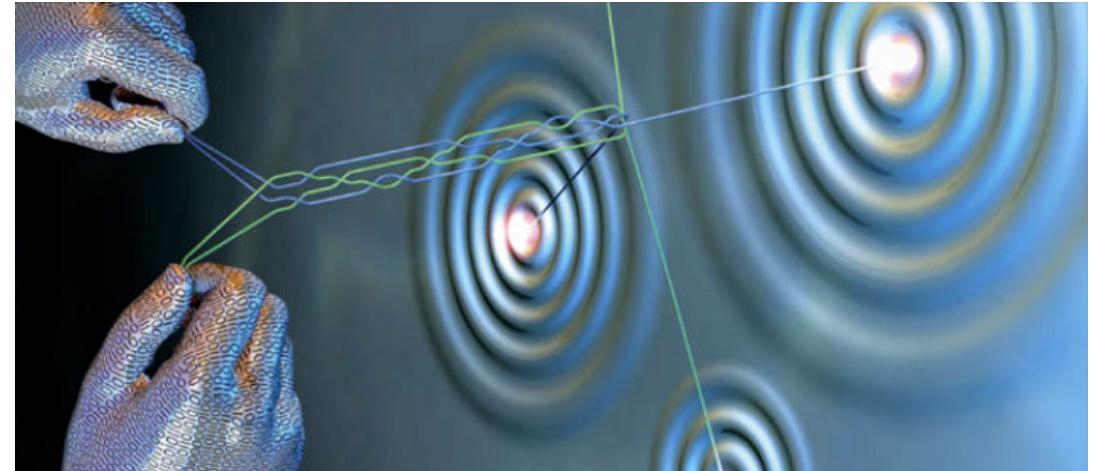
New Konzepte zur Energiekonversion



Verlustfreier Elektronentransport



Neue Elektronik – Spin



Ultraschnelle Quantencomputer

Thank you!



Johannes Gooth



Niels Schröter



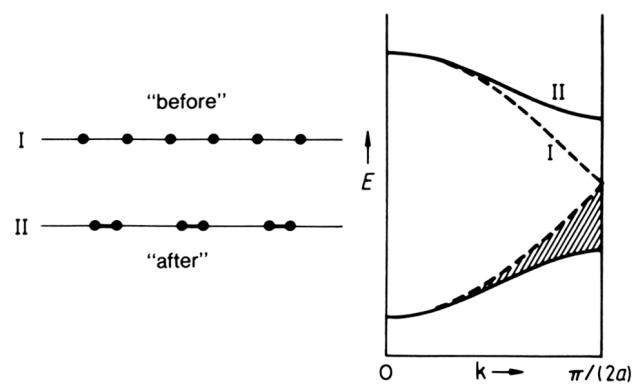
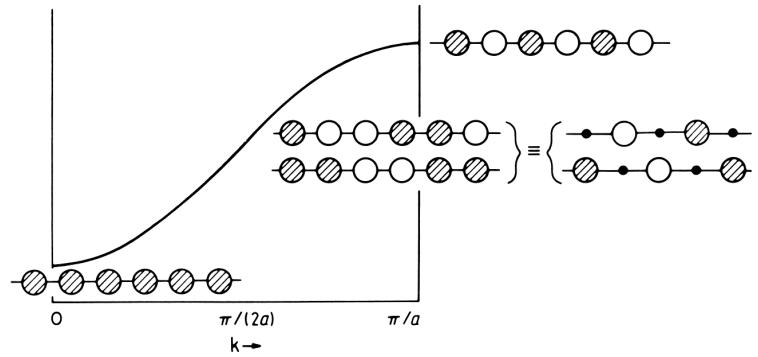
Alexander von Humboldt
Stiftung / Foundation



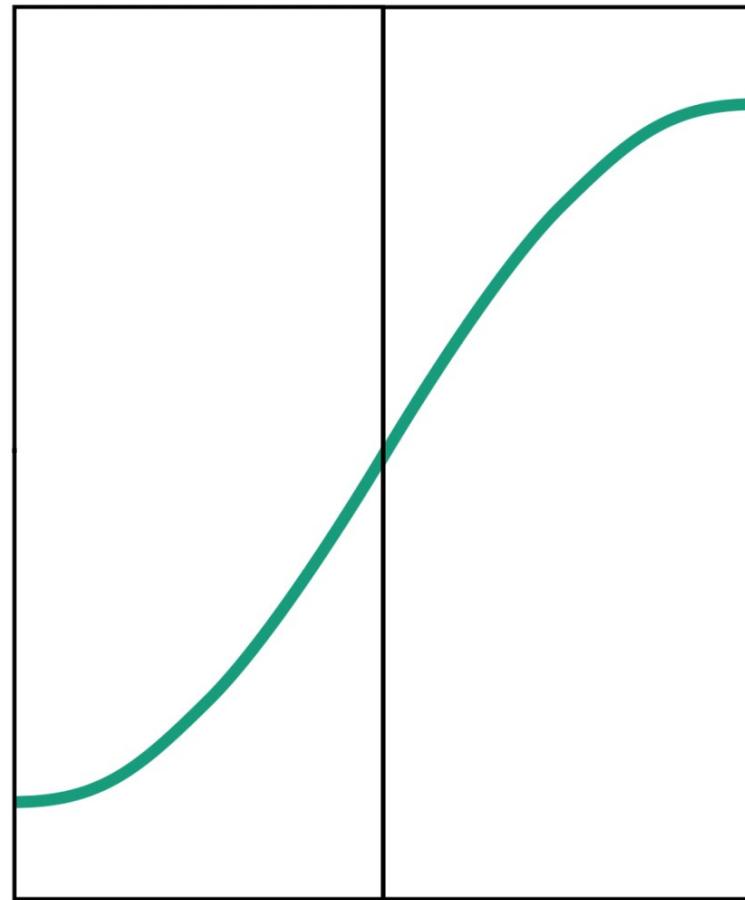
Andrei Bernevig

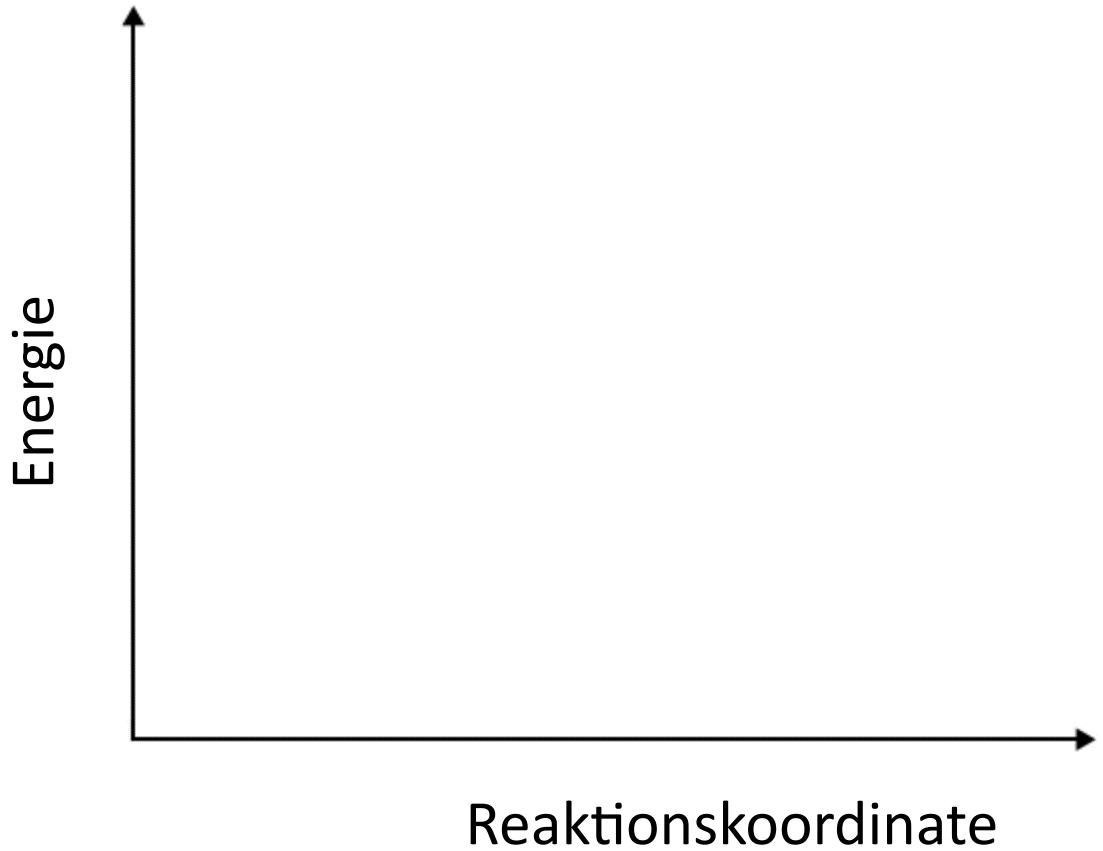


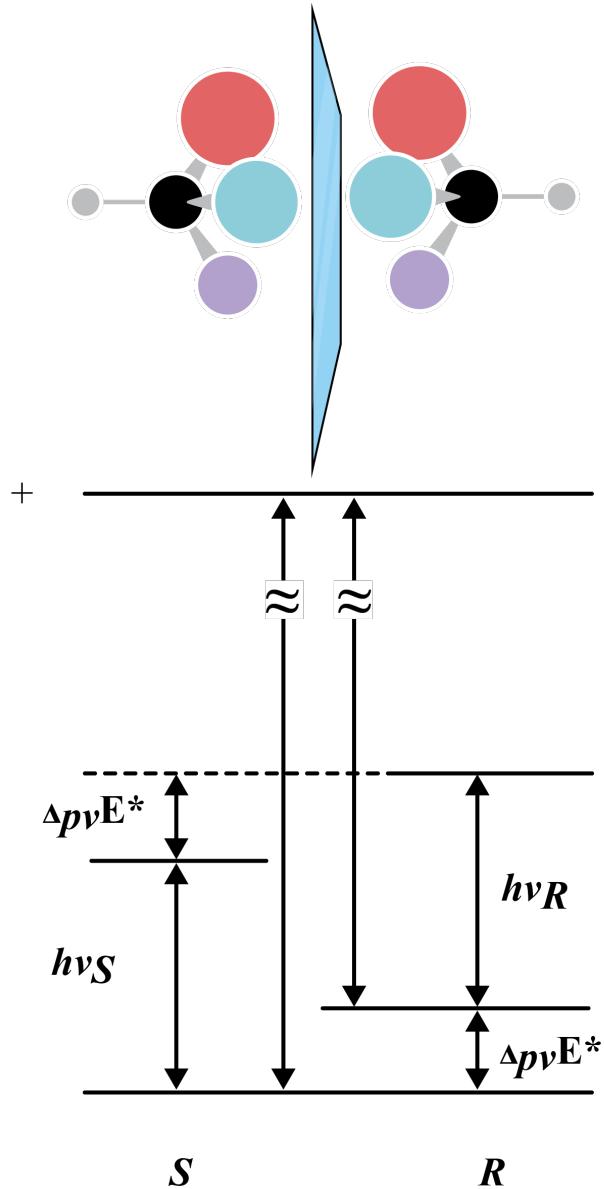
Maia Vergniory

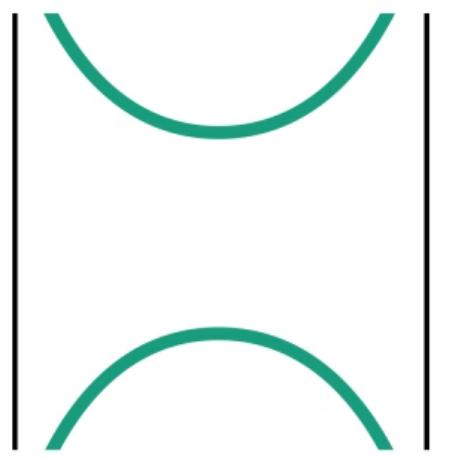


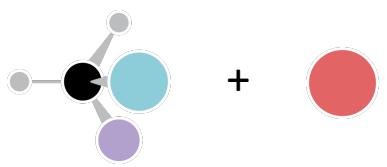
$$a' = 2a \\ k = \pi/2a$$



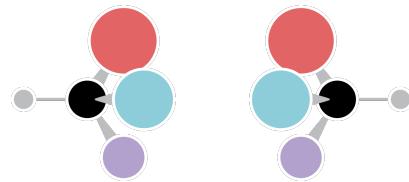








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